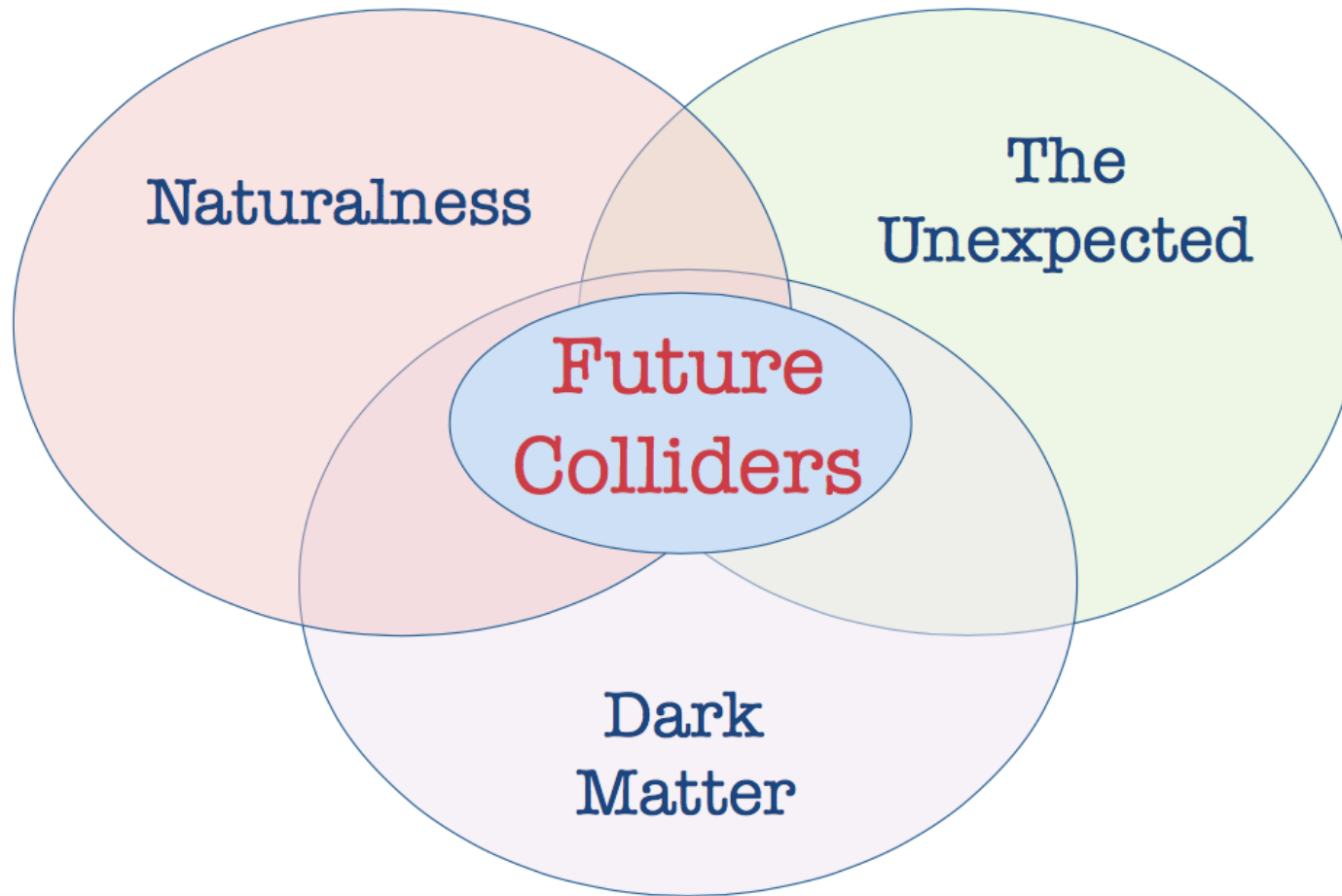


Physics and Experiments at Future pp Colliders

Ashutosh Kotwal
Duke University



BSM Workshop
T. D. Lee Institute and Shanghai Jiao Tong University
July 2018

Dawn of a New Age

- 2008 Nobel Prize in Physics

"for the discovery of the mechanism of spontaneously broken symmetry in subatomic physics"



- 2013 Nobel Prize in Physics

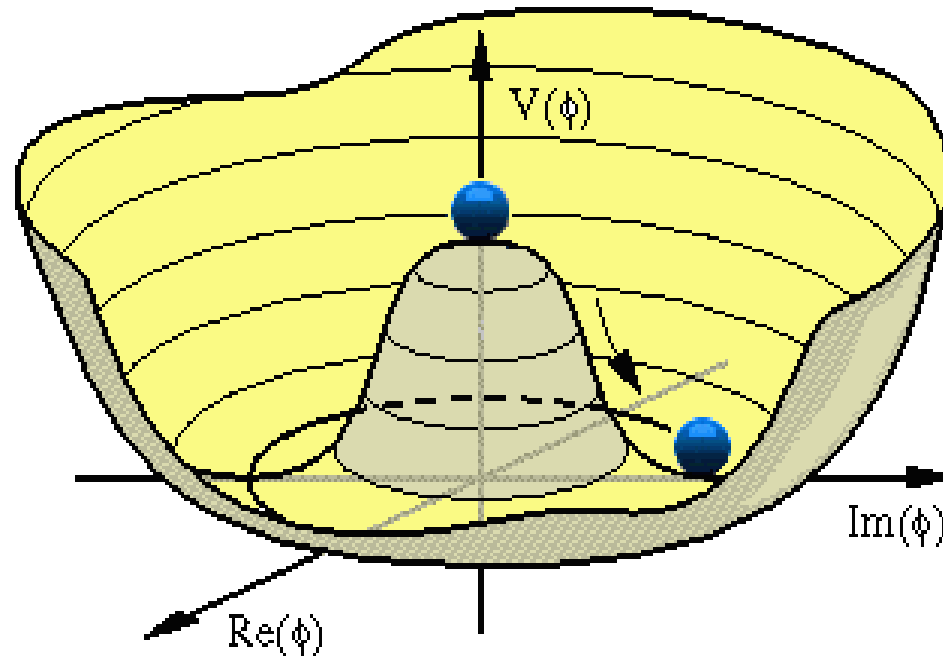
"for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"

Old and New Questions

- How to think of the vacuum as an “electroweak condensed state” ?
- How are the mysteries associated with a single, fundamental scalar field solved?
- What is the origin and nature of Dark Matter?
- What is the origin of the Baryon Asymmetry in the Universe?
- Why is Dark Energy so small but non-zero?
-
-

Spontaneous Symmetry Breaking of Gauge Symmetry

- scalar Higgs field develops a vacuum expectation value (Vev) via spontaneous symmetry breaking
 - Goldstone modes appear as the new longitudinal modes of gauge bosons



- Phase transition \rightarrow vacuum state possesses non-trivial quantum numbers
 - Dynamical origin of this phase transition is not known
 - Implies vacuum is a condensed, superconductor-like state

Fundamental vs Parametric Physics

- Fundamental principles lead to
 - Chiral fermions from irreducible representations of Lorentz group
 - fermions as spin $\frac{1}{2}$ representations of Lorentz group
 - Fermi-Dirac statistics \rightarrow Pauli Exclusion Principle
 - why matter occupies volume
 - Massless force mediators (gauge bosons) from gauge invariance
 - Massive gauge bosons and fermions from spontaneous breaking of gauge symmetry
- In comparison, the breaking of gauge symmetry by the Higgs VeV is parametrically induced
 - No dynamic or underlying principle behind it in the Standard Model

Why is Higgs Puzzling

Gauge sector

$$L = i\bar{\psi}\gamma^\mu D_\mu\psi - \frac{1}{2}F_{\mu\nu}F^{\mu\nu}$$

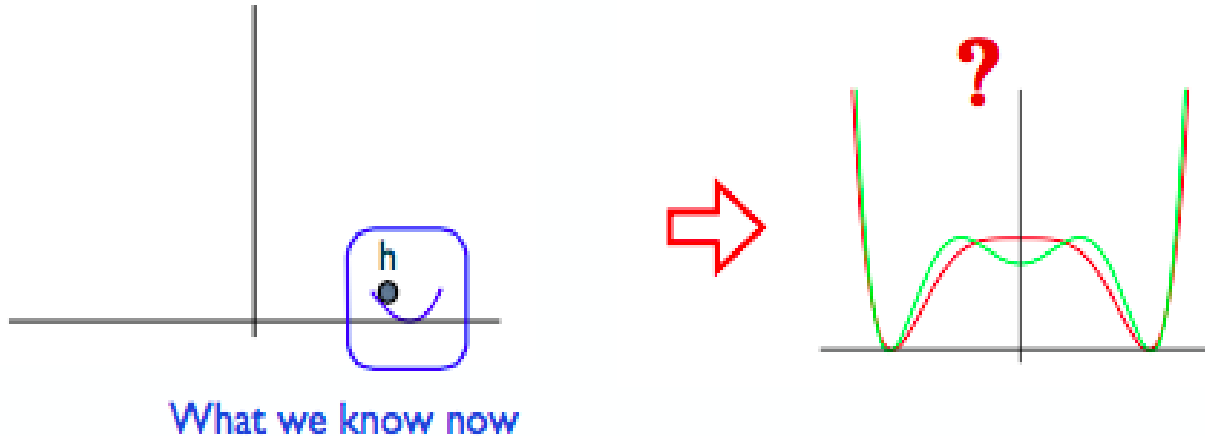
particle	spin
quark: u, d,...	1/2
lepton: e...	1/2
photon	1
W,Z	1
gluon	1
Higgs	0

h: a new kind of
elementary particle

Higgs sector

$$L = (h_{ij}\bar{\psi}_i\psi_j H + \text{h.c.}) - \lambda|H|^4 + \mu^2|H|^2 - \Lambda_{cc}^4$$

Why is Higgs Puzzling



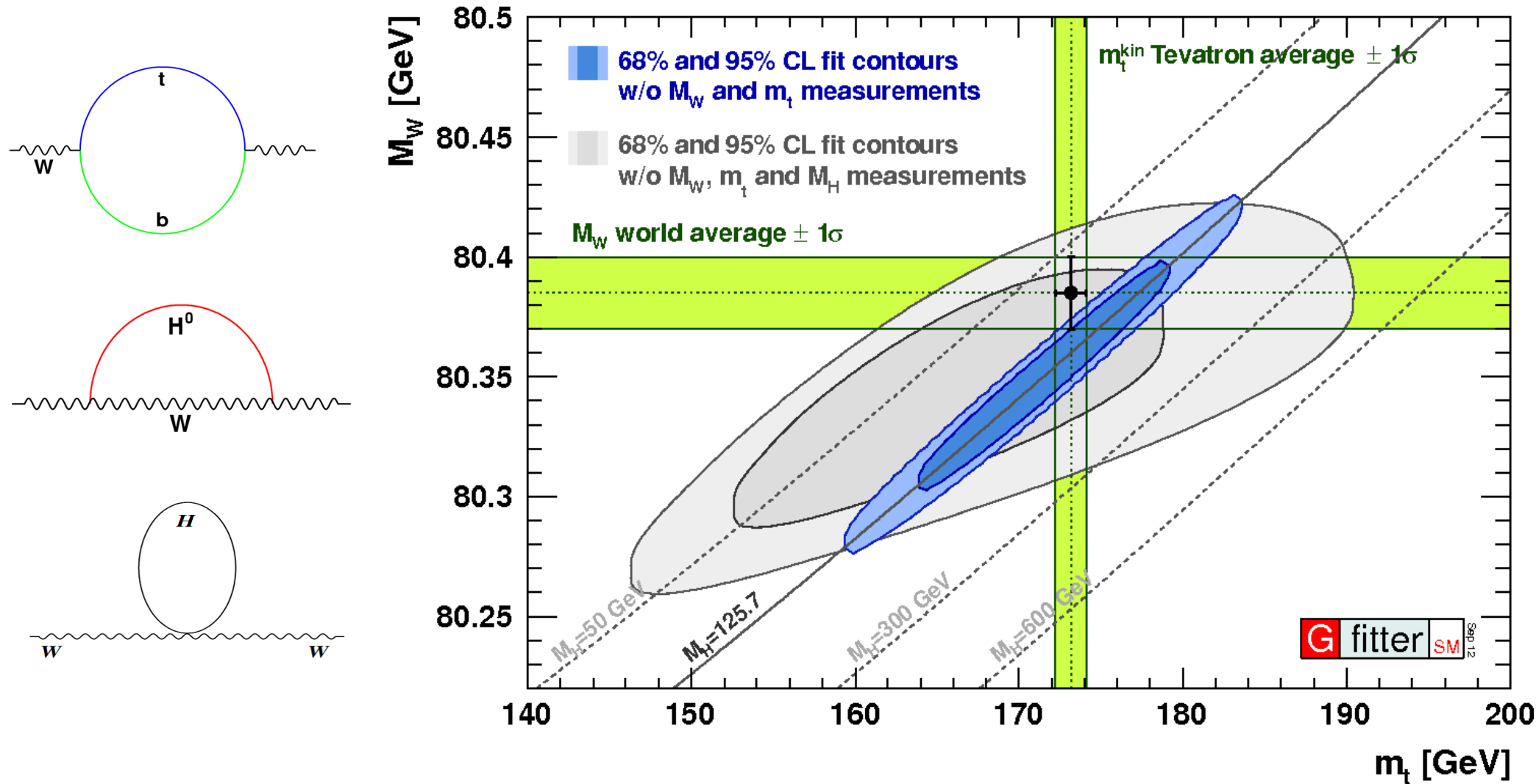
$$V(h) = \frac{1}{2}\mu^2 h^2 + \frac{\lambda}{4}h^4 \quad \text{or} \quad V(h) = \frac{1}{2}\mu^2 h^2 + \frac{\lambda}{4}h^4 + \frac{1}{\Lambda^2}h^6$$

Ad-hoc potential, similar to and motivated by Landau-Ginzburg theory of superconductivity

Higgs potential in SM can be extrapolated to Planck scale without additional parameters; but no a-priori reason for a parameterization to respect this condition

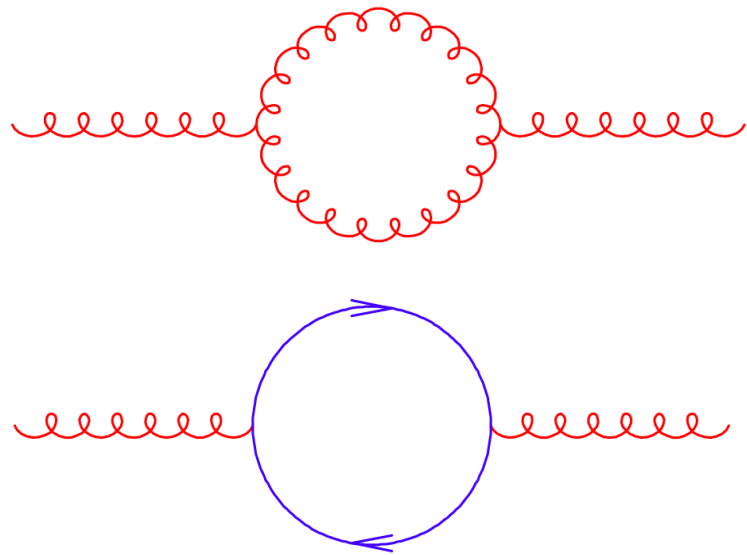
Radiative Stability of Higgs potential parameters

Test of Electroweak Quantum Loops at High Energy – Example I

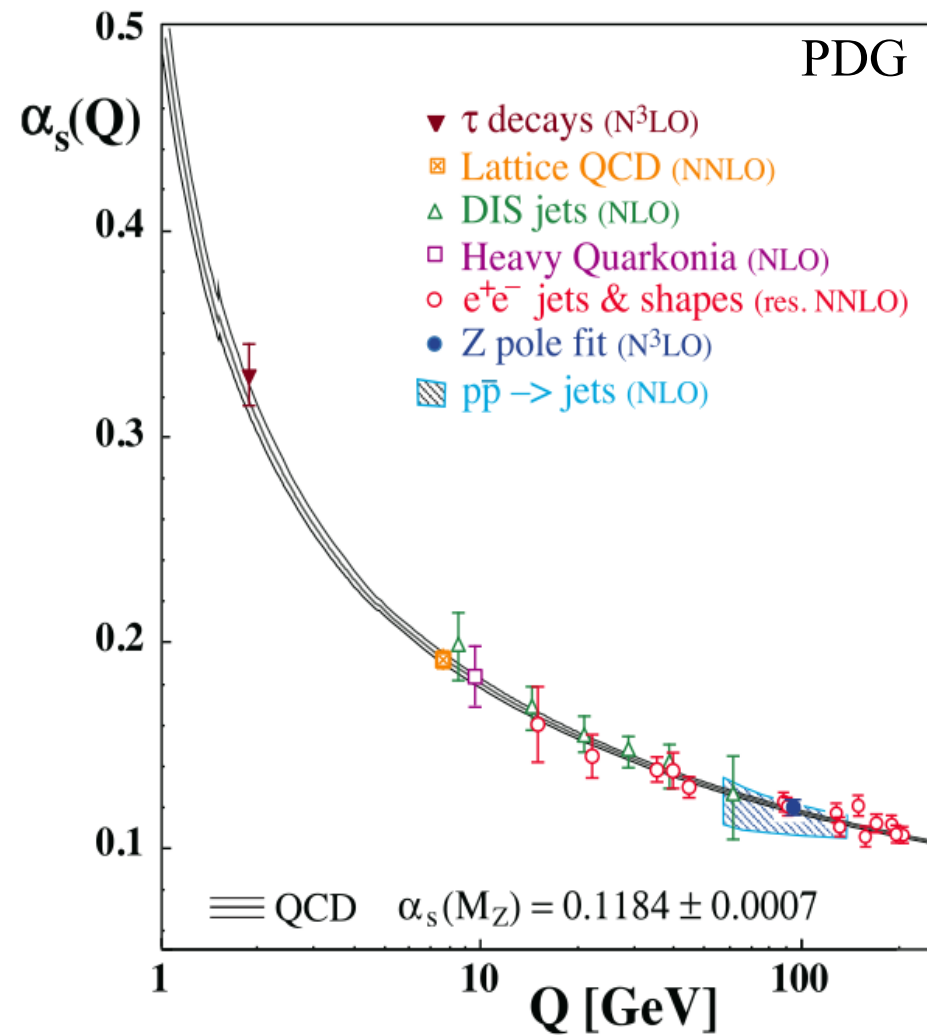


The top quark mass, the W boson mass and the mass of the Higgs boson provides a stringent test of the standard model at loop level

Example II - Test of QCD Quantum Loops at High Energy



Running of strong coupling
has been confirmed experimentally



Why is the Higgs Boson so Light?

$$m_H^2 - m_{\text{bare}}^2 = \left(\text{Higgs loop} \right) + \left(\text{top quark loop} \right) + \left(\text{W,Z loop} \right)$$

$\lambda \int^{\Lambda} d^4k (k^2 - m_H^2)^{-1} \sim \Lambda^2 \lambda$

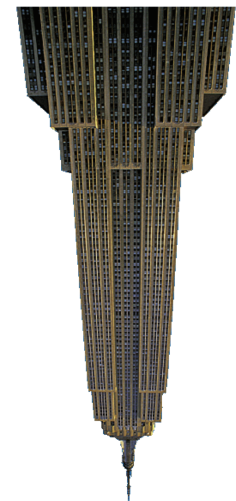
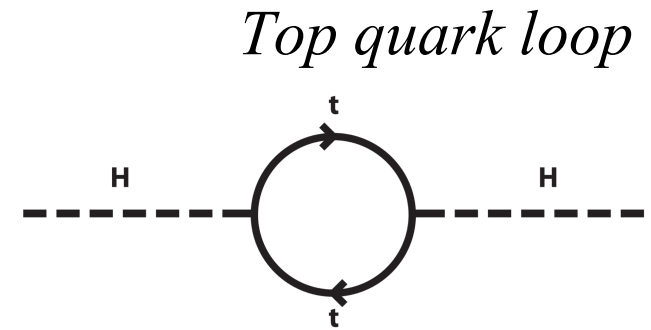
For the first time, we have additive corrections to parameters which are quadratically divergent

The Higgs boson ought to be a very heavy particle, naturally

However, observed $m_H \ll \Lambda$

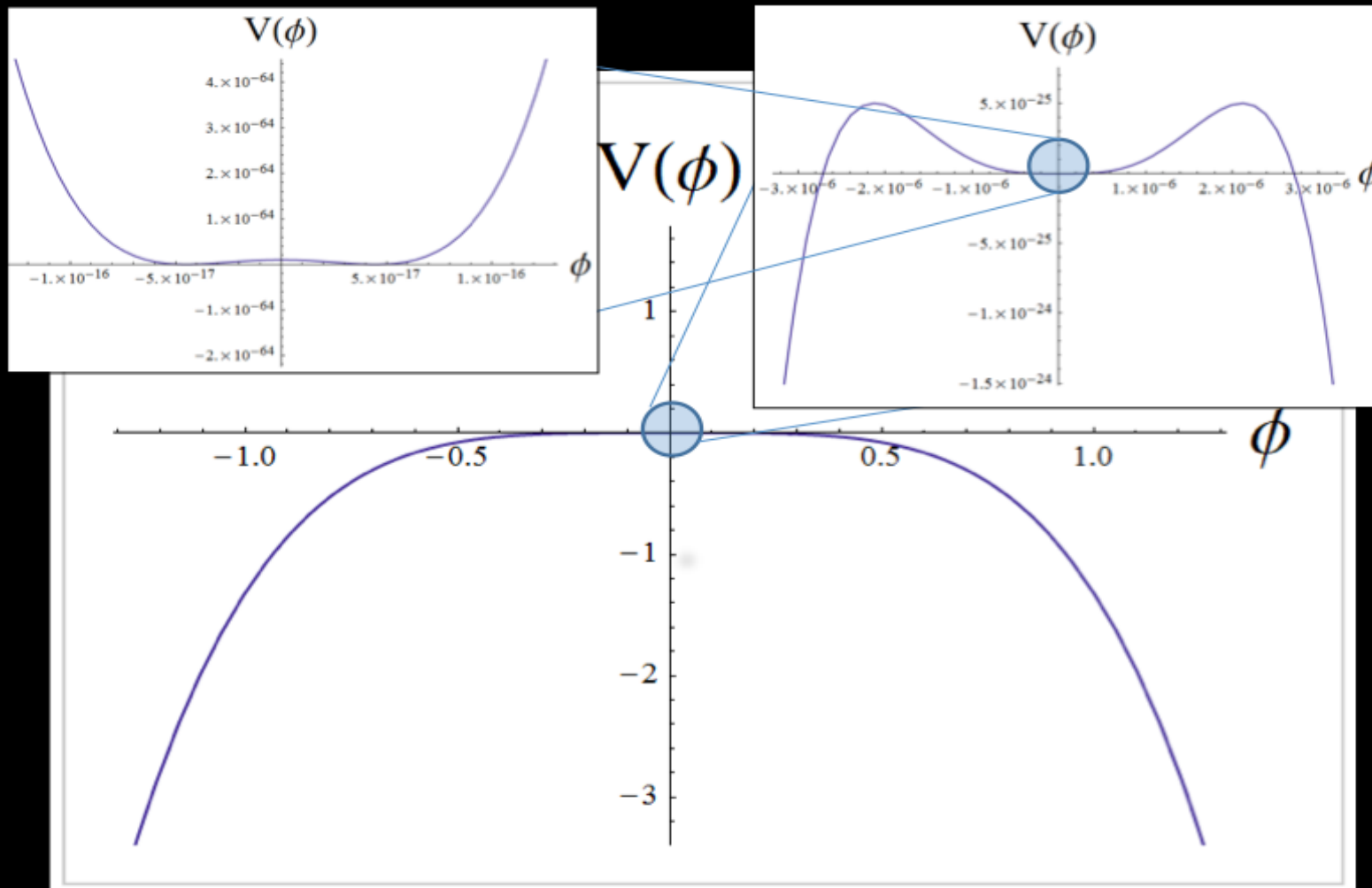
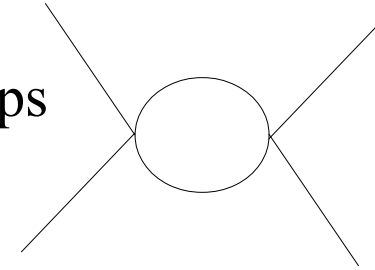
Fine-tuning Problem of Higgs Boson Mass

- The divergent integral in this quantum loop must be regulated by a high-momentum cutoff, Λ , which could be the gravitational Planck energy scale $M_{\text{planck}} \sim 10^{19} \text{ GeV}$
 - Loop calculation gives Higgs boson mass correction $\sim M_{\text{planck}}^2$
- physical Higgs boson mass $\sim 125 \text{ GeV}$
- Therefore need extreme “fine-tuning” of bare lagrangian parameters at high energy



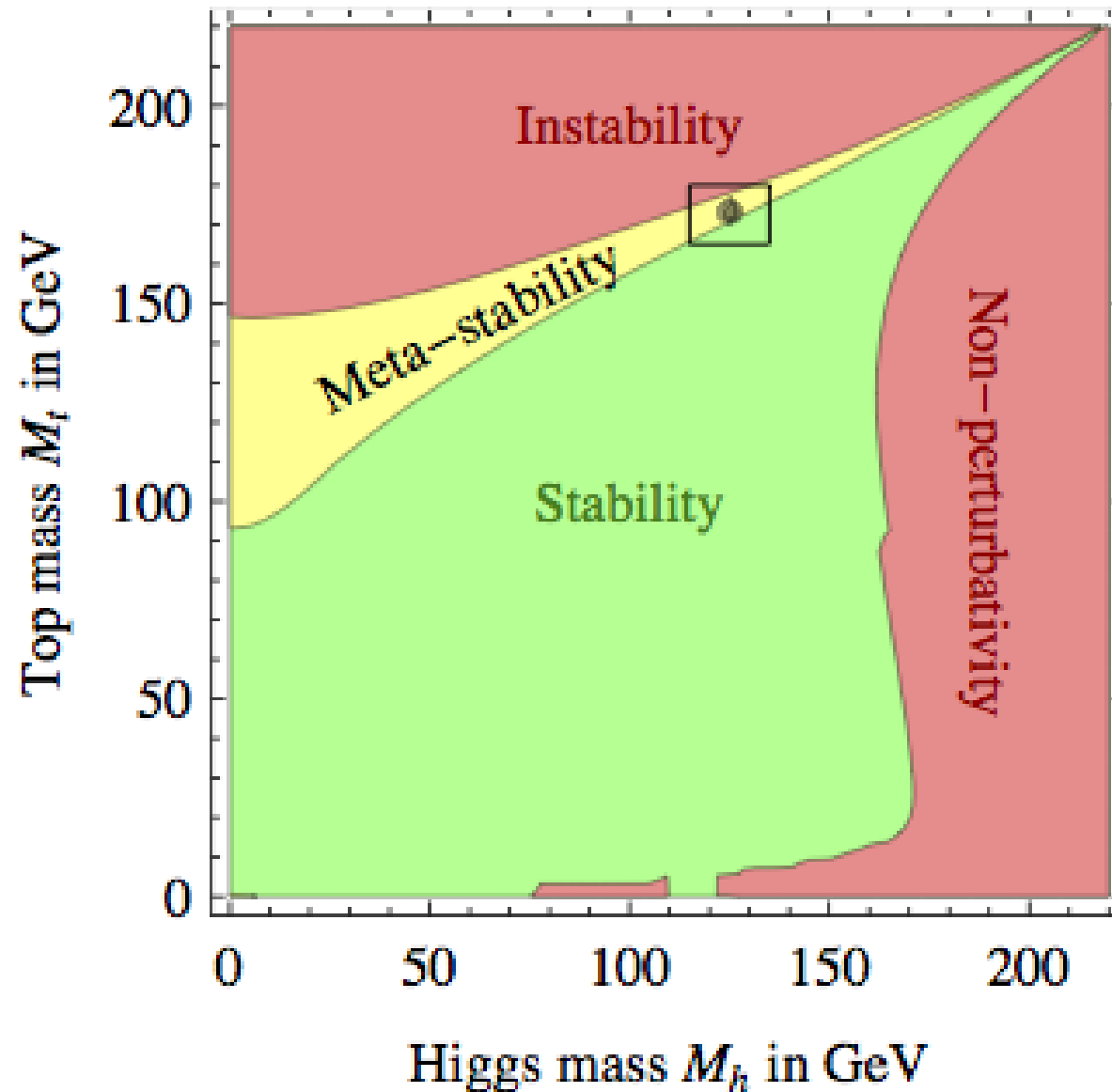
Radiative Corrections to Higgs Self-Coupling

- $\lambda|\phi|^4$ receives radiative corrections from Higgs and top-quark loops



(from Paul Steinhardt)

Stability of Electroweak Vacuum



Higgs boson puzzles

- First fundamental (?) scalar field to be discovered
- Spontaneous symmetry breaking by development of a VeV
 - But VeV is induced parametrically by ad-hoc Higgs potential, no dynamics
- Parameters of Higgs potential are not stable under radiative corrections
 - First time that the radiative correction to a particle mass is additive and quadratically divergent
 - Gauge boson masses are protected by gauge invariance
 - Fermion masses are protected by chiral symmetry of massless fermions
- Single scalar Higgs field is a strange beast, compared to fermions and gauge bosons
- Additional symmetries and/or dynamics strongly motivated by Higgs discovery

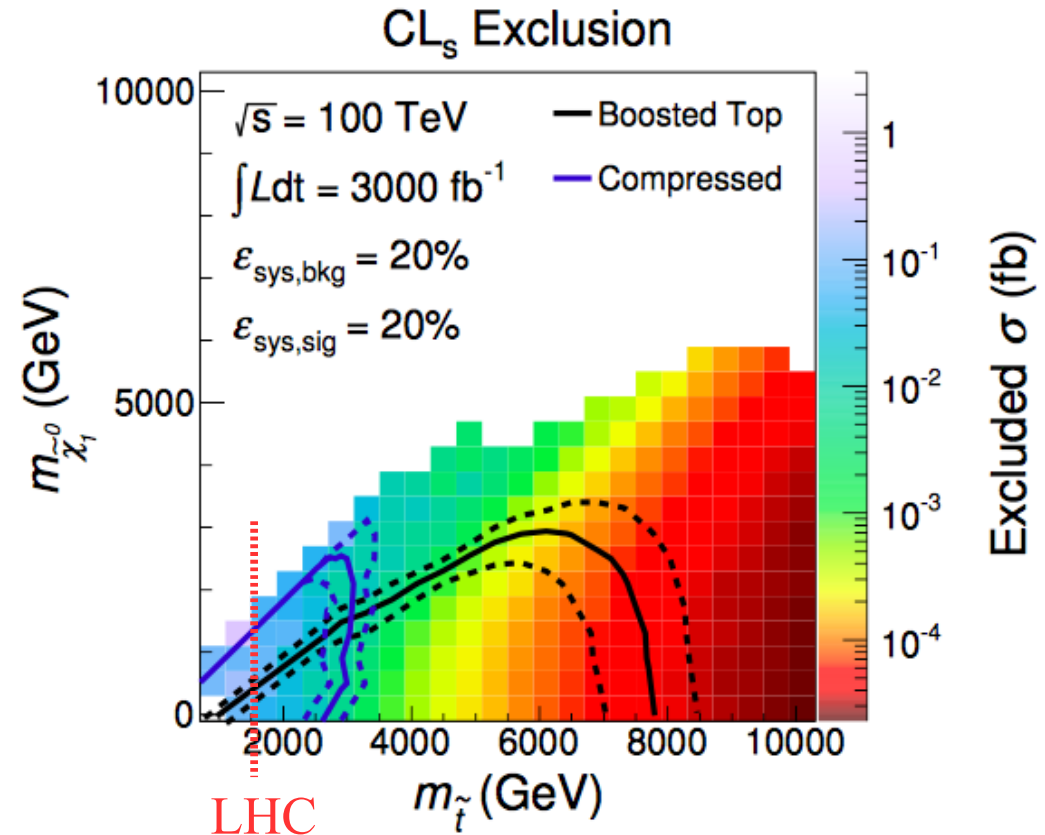
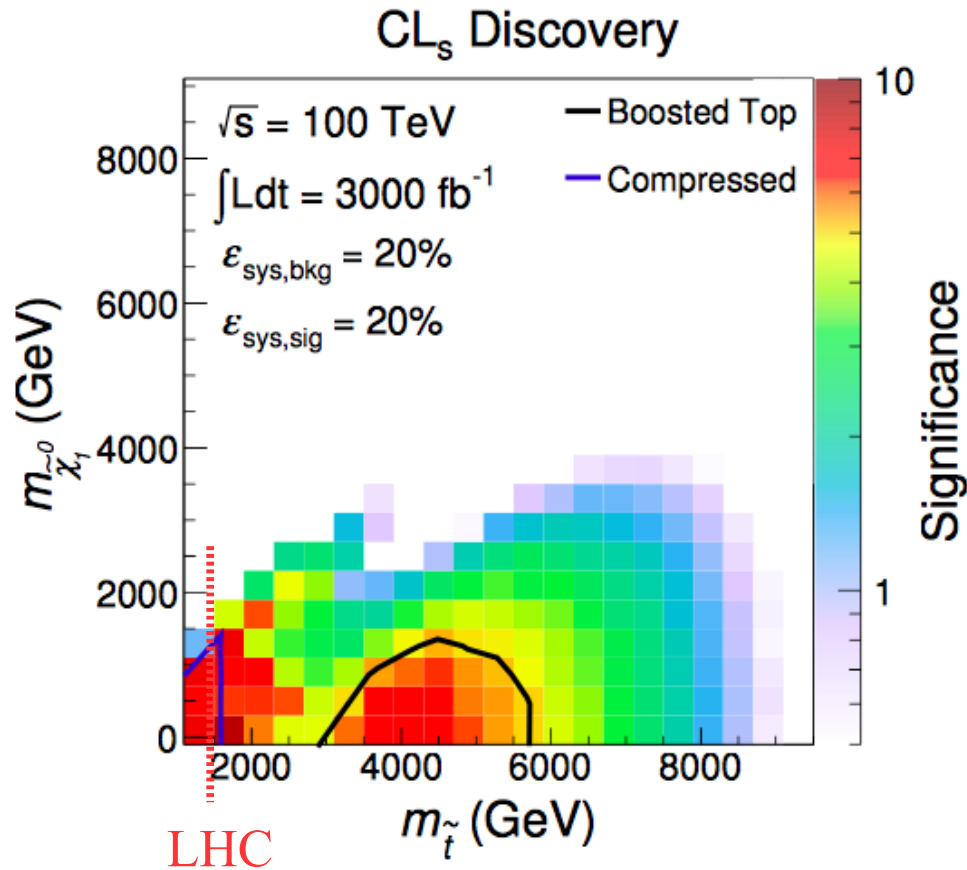
Circular pp Collider

Circular pp Collider Physics Goals

- Testable reasons why the Standard Model must be incomplete
 - Dark Matter could be
 - Weakly-interacting particles
 - Particles interacting through Higgs portal
 - Interacting with SM particles through gravity
 - Electroweak Baryogenesis
 - Can the electroweak phase transition (formation of Higgs VEV) provide the out-of-equilibrium condition needed for matter-antimatter asymmetry observed?
 - Can the parameter space of new physics be a bounded parameter space?
 - Can it be fully covered with a 100-TeV scale pp collider?
- Naturalness – the need to explain the lightness of the Higgs mass – testing Naturalness at 10^{-4}

Supersymmetric Colored Top Partner Sensitivity

(Cohen *et al*, 2014)



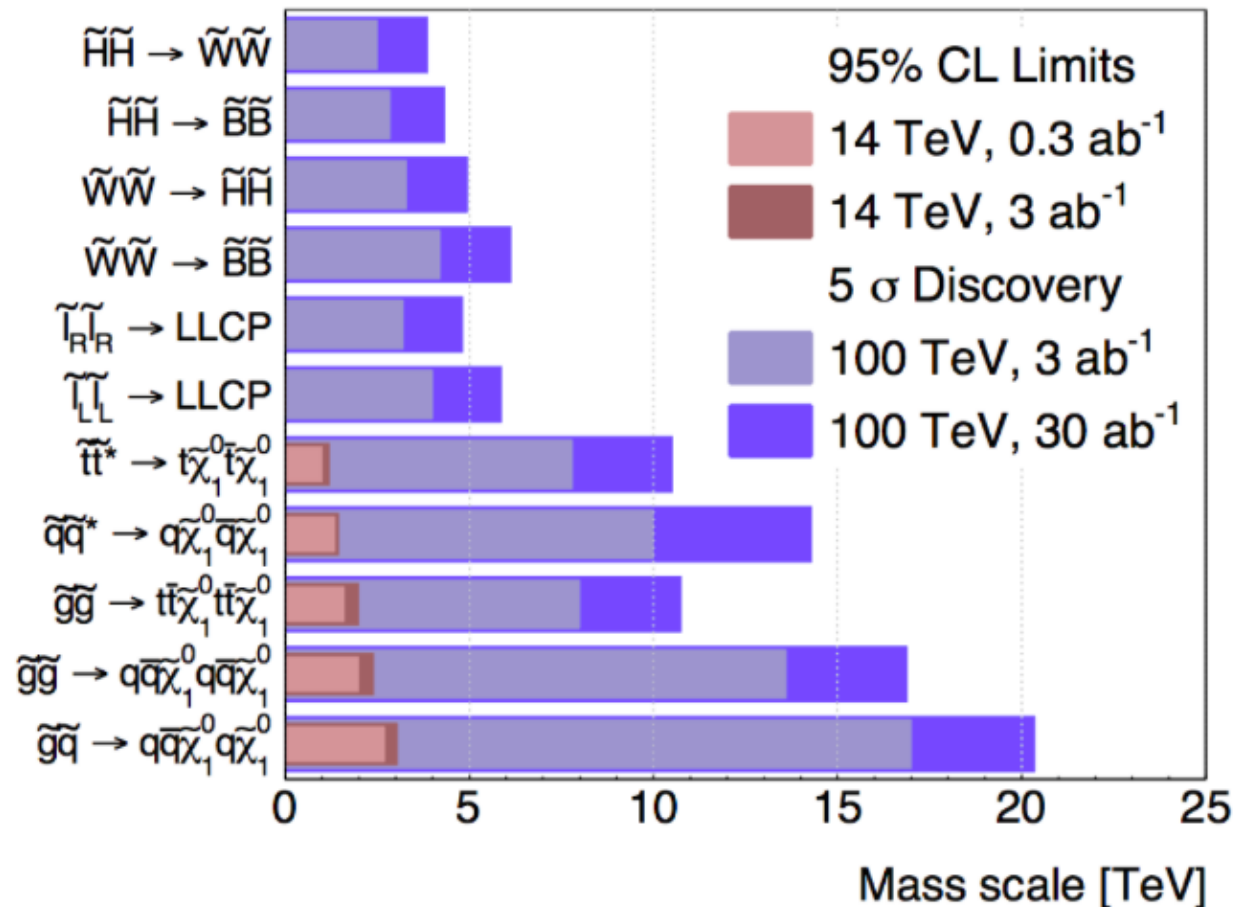
$$\text{Fine-tuning} \sim m_{\text{stop}}^2 \sim 10^{-4}$$

A big jump beyond LHC

Discovering or eliminating “natural” low-energy SUSY

Exploring New Territory – Squarks and Gluinos

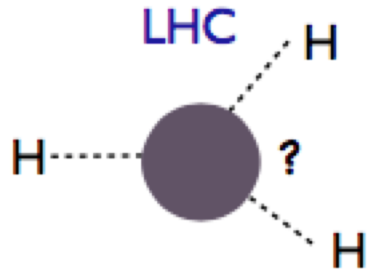
Summary from FCC Report:



Squark & gluino discovery potential up to 10-20 TeV

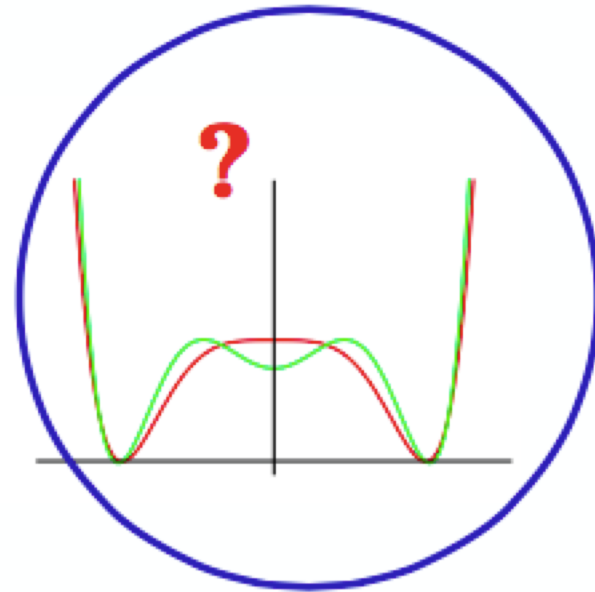
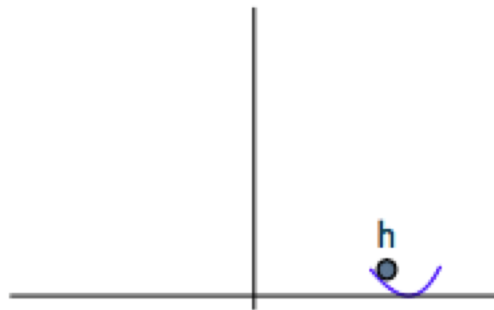
Full exploration of “low-scale” SUSY

Higgs Self-Coupling



Unique type of coupling for spin-0 scalars
Not seen before in nature!

Measuring it well is crucial to
answer this question.



Expect $O(1)$ deviations from SM in self-coupling coefficient

Measuring the Higgs Self-Coupling

- $gg \rightarrow HH$ (most promising?) , $qq \rightarrow HHqq$ (via VBF)
- Reference benchmark process: $HH \rightarrow b\bar{b} \gamma\gamma$
- Goal: 5% (or better) precision for SM selfcoupling

$HH \rightarrow b\bar{b} \gamma\gamma$	Barr,Dolan,Englert,Lima, Spannowsky JHEP 1502 (2015) 016	Contino, Azatov, Panico, Son arXiv:1502.00539	He, Ren Yao arXiv:1506.03302
FCC@100TeV 3/ab	30~40%	30%	15%
FCC@100TeV 30/ab	10%	10%	5%
S/\sqrt{B}	8.4	15.2	16.5
Details	<ul style="list-style-type: none"> ✓ λ_{HHH} modification only ✓ $c \rightarrow b$ & $j \rightarrow \gamma$ included ✓ Background systematics ○ $b\bar{b} \gamma\gamma$ not matched ✓ $m_{\gamma\gamma} = 125 \pm 1$ GeV 	<ul style="list-style-type: none"> ✓ Full EFT approach ○ No $c \rightarrow b$ & $j \rightarrow \gamma$ ✓ Marginalized ✓ $b\bar{b} \gamma\gamma$ matched ✓ $m_{\gamma\gamma} = 125 \pm 5$ GeV ✓ Jet / W_{had} veto 	<ul style="list-style-type: none"> ✓ λ_{HHH} modification only ✓ $c \rightarrow b$ & $j \rightarrow \gamma$ included ○ No marginalization ✓ $b\bar{b} \gamma\gamma$ matched ✓ $m_{\gamma\gamma} = 125 \pm 3$ GeV

**Work in progress to compare studies, harmonize performance assumptions, optimize, etc
⇒ ideal benchmarking framework**

Origin of Matter-Antimatter Asymmetry

Origin of Baryon Asymmetry

POSSIBLE EXPLANATIONS...

$$\frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-9} \text{ (from BBN)}$$

⇒ **Baryogenesis at EW Scale**  **TESTABLE!**

⇒ ...

SAKHAROV CONDITIONS (for dynamical generation of baryon asymmetry)

B Violation ✓ *Sphalerons*

V. A. Kuzmin, V. A. Rubakov, M. Shaposhnikov, Phys. Lett. B155 (1985) 36

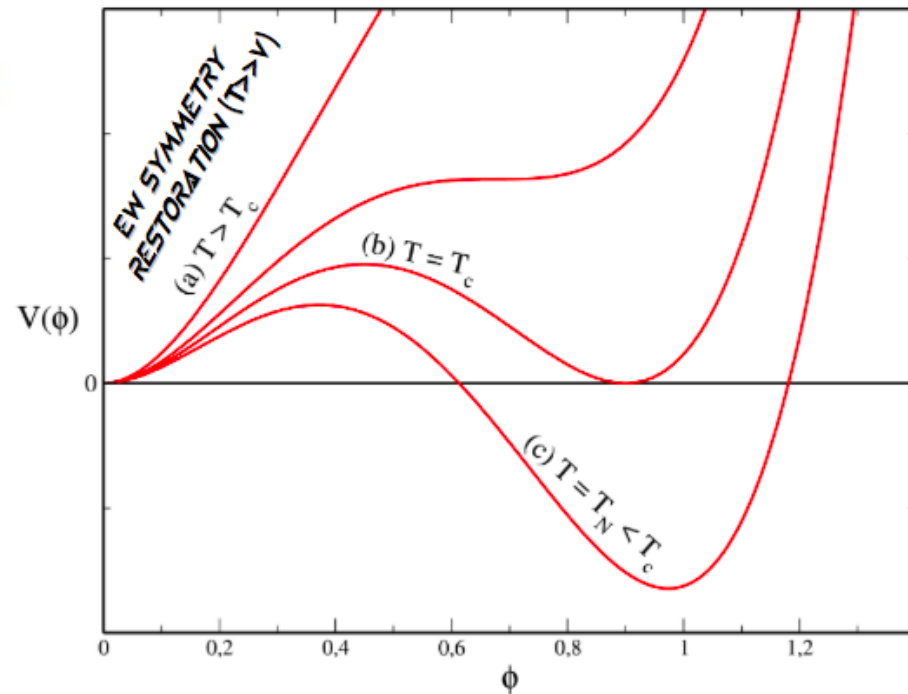
C/CP Violation ✗ *not enough*

Departure from Thermal Equilibrium ✗ *not enough*

Baryon Asymmetry and Electroweak Phase Transition

1st Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$ Discontinuous

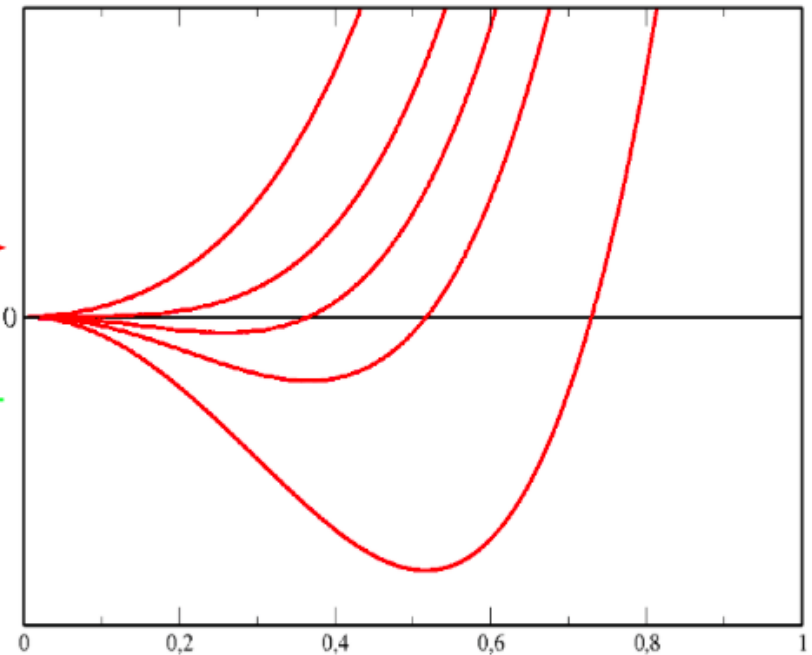


LARGER M_H

NEW BOSONS

2nd Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$ Continuous

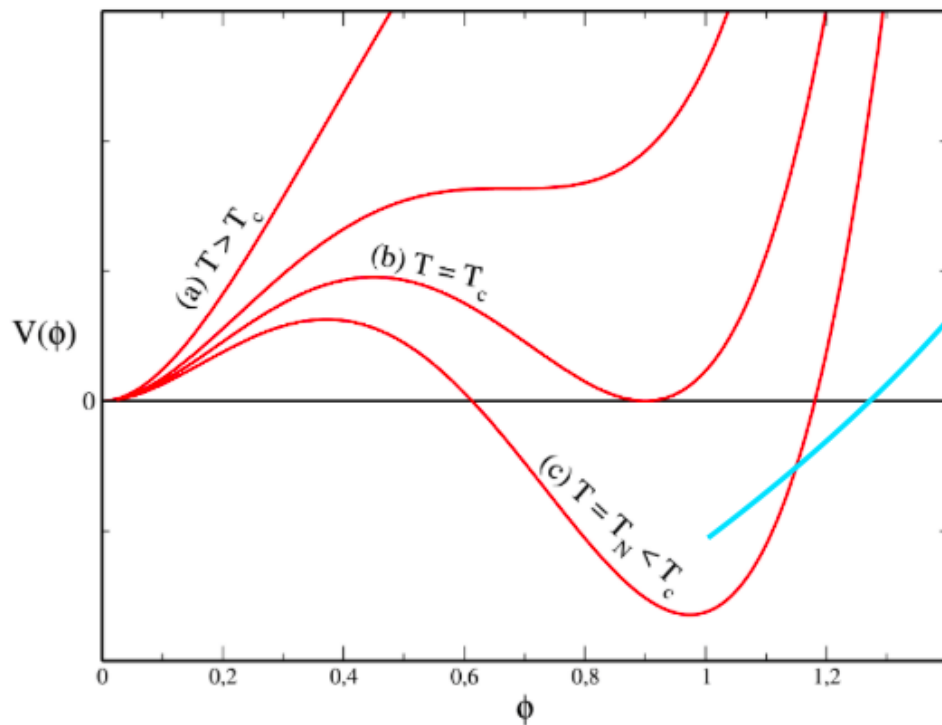


In the SM ($m_h = 125$ GeV) EW Phase Transition Smooth CrossOver
*K. Kajantie, M. Laine, K. Rummukainen, M. Shaposhnikov, Phys. Rev. Lett. **77** (1996) 2887*

Baryon Asymmetry and Electroweak Phase Transition

1st Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$ Discontinuous

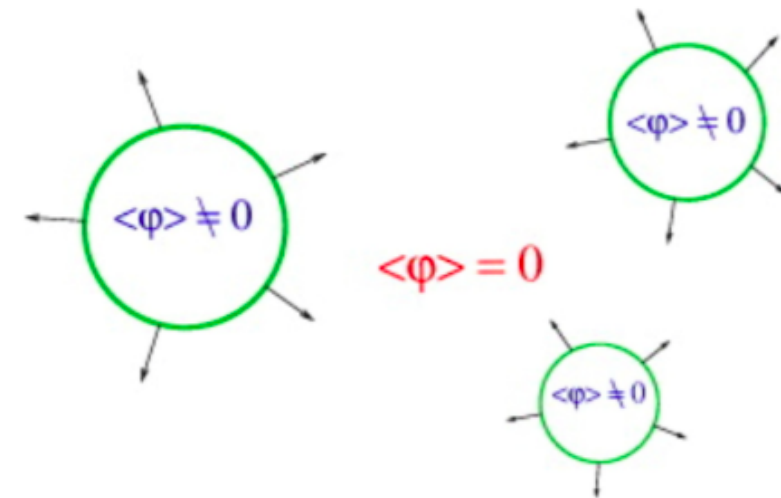


Nucleation of True Vacuum Bubbles
(in False Vacuum Sea)

*J. S. Langer, Ann. Phys. **54** (1969) 258*

*S. R. Coleman, Phys. Rev. D **15** (1977) 2929*

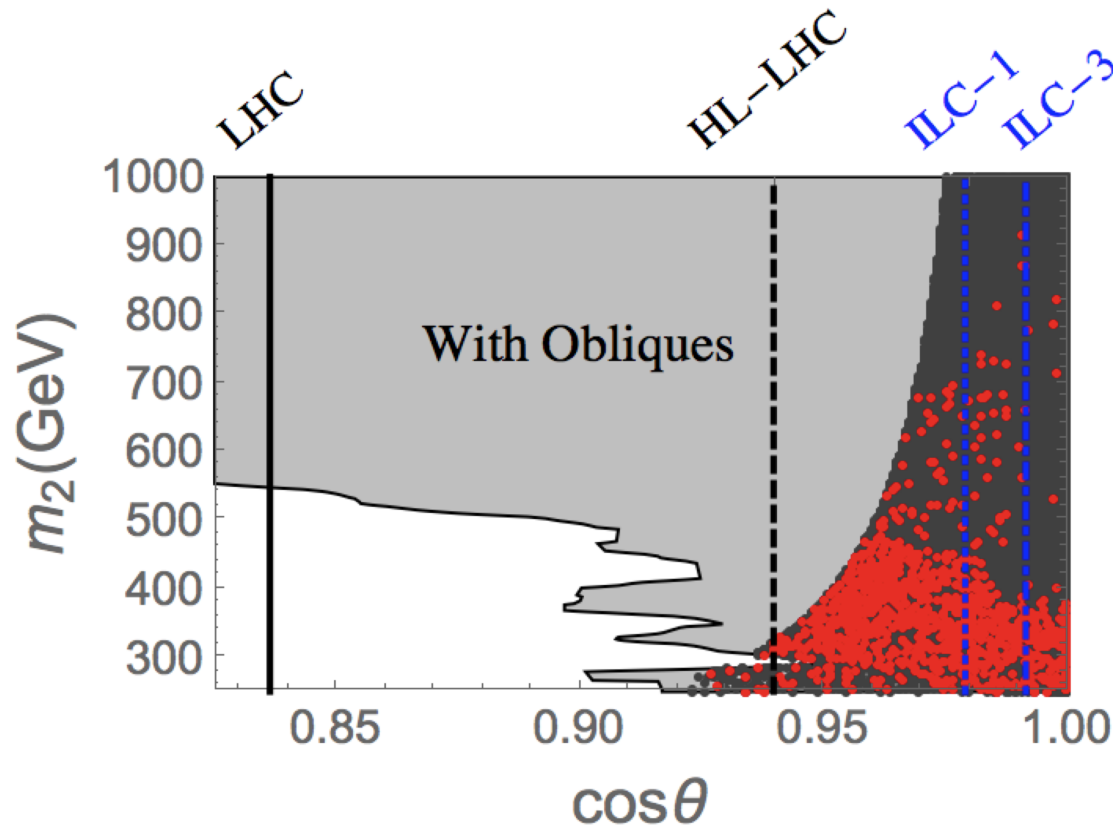
*A. D. Linde, Nucl. Phys. B **216** (1983) 421*



SUDDEN CHANGE IN HIGGS VEV

First Order Phase Transition

$$V(H, S) = -\mu^2 (H^\dagger H) + \lambda (H^\dagger H)^2 + \frac{a_1}{2} (H^\dagger H) S + \frac{a_2}{2} (H^\dagger H) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4$$



(from P. Winslow)

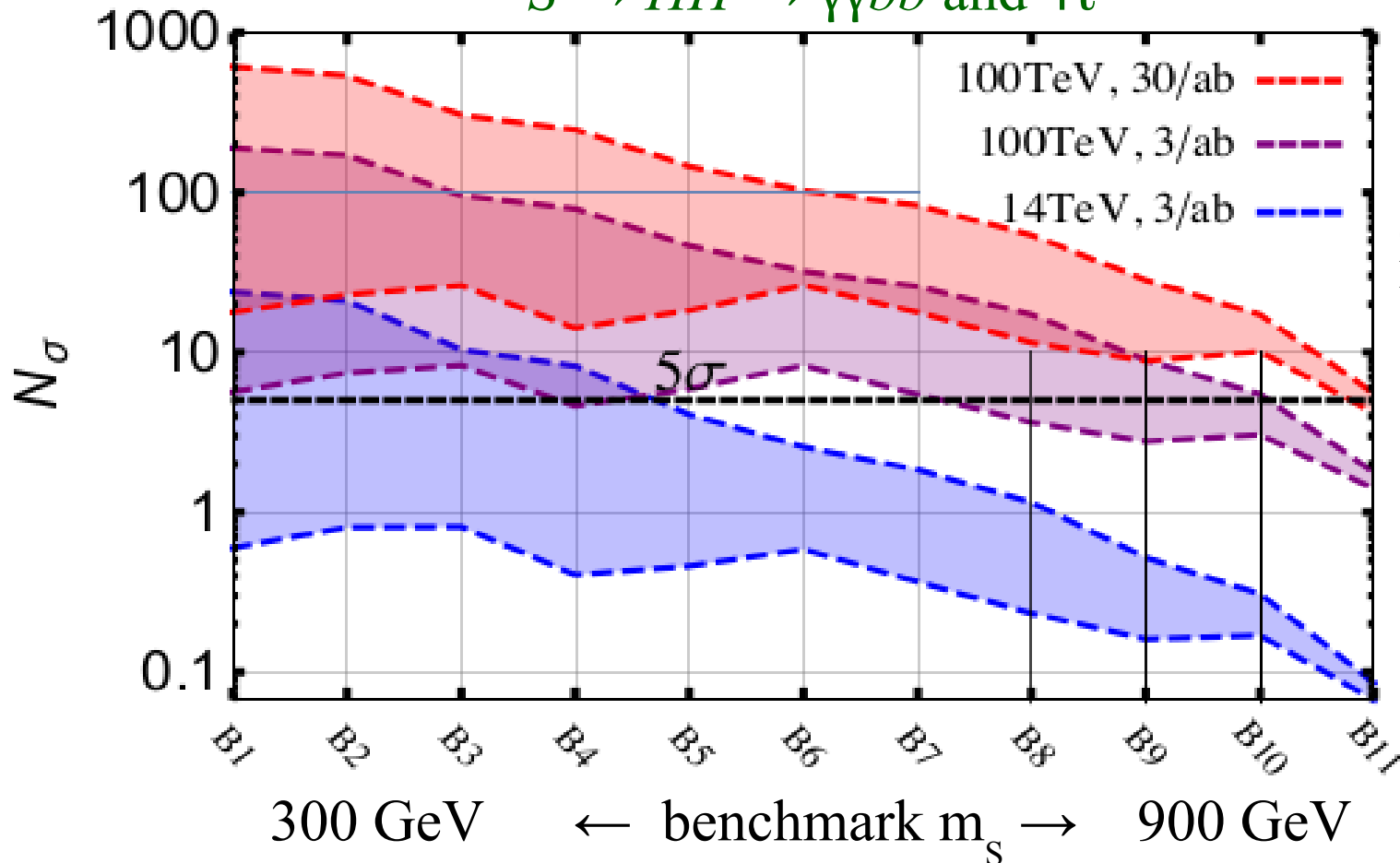
S. Profumo, M. J. Ramsey-Musolf, C. L. Wainwright and P. Winslow, arXiv:1407.5342

Can TeV-scale new physics associated with 1st order phase transition be completely covered by a pp collider?

Inducing First-Order Electroweak Phase Transition

$$V(H, S) = -\mu^2 (H^\dagger H) + \lambda (H^\dagger H)^2 + \frac{a_1}{2} (H^\dagger H) S + \frac{a_2}{2} (H^\dagger H) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4$$

$S \rightarrow HH \rightarrow \gamma\gamma bb$ and 4τ



(P. Winslow, J.M. No,
M.J. Ramsey-Musolf, AVK)

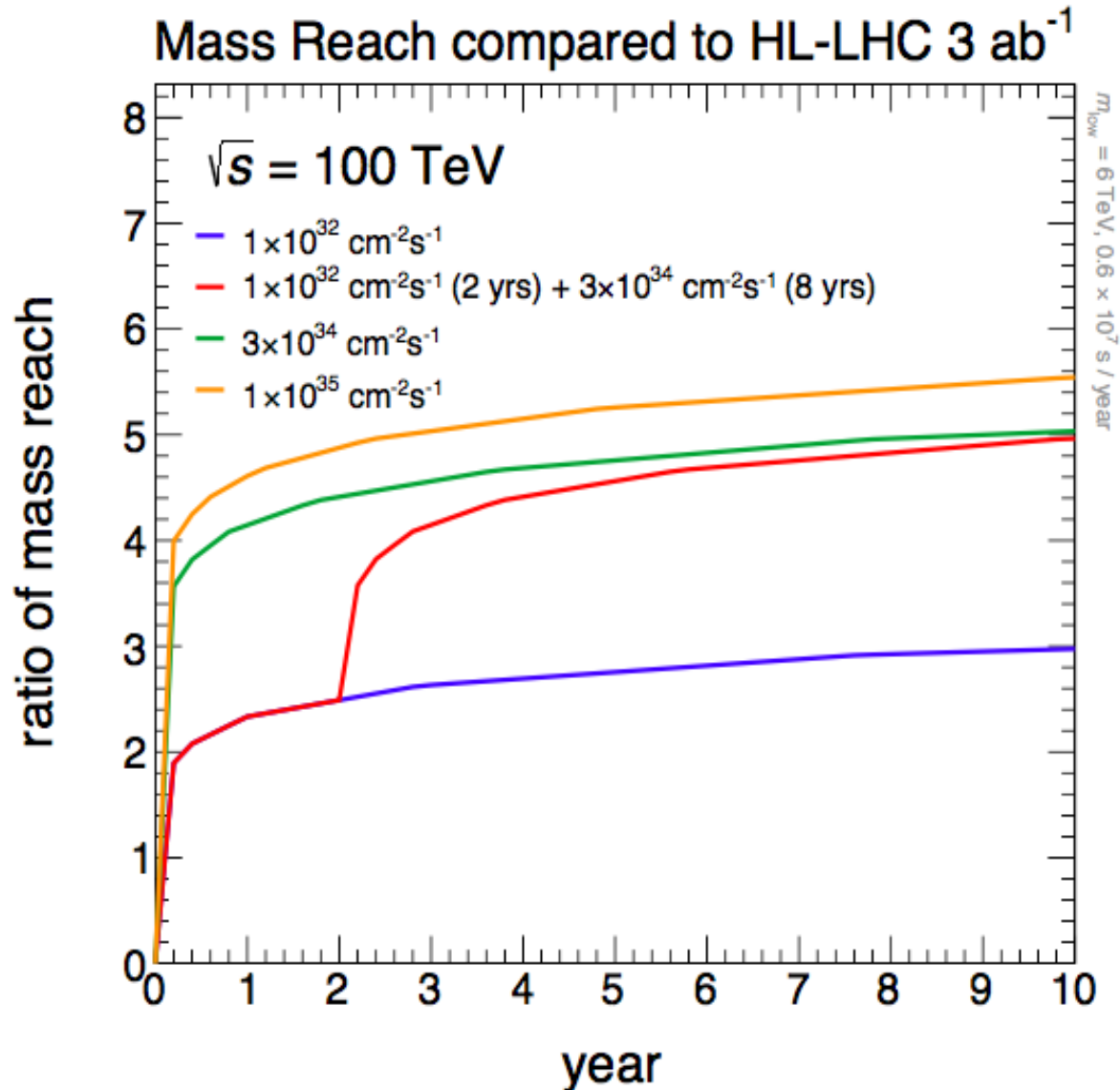
Discovery potential across entire parameter space

Guidance for Detector Design

- As long as Standard Model continues to work, “higher energy is better”
- Covering the “Naturalness-motivated” models push towards higher masses
- Dark Matter, Electroweak Baryogenesis *may* relate to physics at lower masses and smaller couplings
- Other reasons that new physics may hide at low mass with weak couplings
 - “Neutral Naturalness” (partners without QCD color charge)
 - e.g. twin Higgs, Hidden Sector
 - Higgs portal to new sector (SM interactions via Higgs only)
- Implications for detector design: larger dynamic range of p_T of objects
 - Starting at ~ 20 GeV leptons, photons and b -quarks (same as LHC, e.g. $gg \rightarrow HH$)
 - Going up to ~ 7 times the highest p_T probed at LHC
- Also large rapidity range for all objects due to higher longitudinal boost

Collider Luminosity and Energy

- Collider luminosity evolution for high-mass reach



(from L-T. Wang)

Rate comparisons at 8, 14, 100 TeV

	N_{100}	N_{100}/N_8	N_{100}/N_{14}
gg→H	16 G	4.2×10^4	110
VBF	1.6 G	5.1×10^4	120
WH	320 M	2.3×10^4	66
ZH	220 M	2.8×10^4	84
ttH	760 M	29×10^4	420
gg→HH	28 M		280

$$N_{100} = \sigma_{100\text{TeV}} \times 20 \text{ ab}^{-1}$$

$$N_8 = \sigma_{8\text{TeV}} \times 20 \text{ fb}^{-1}$$

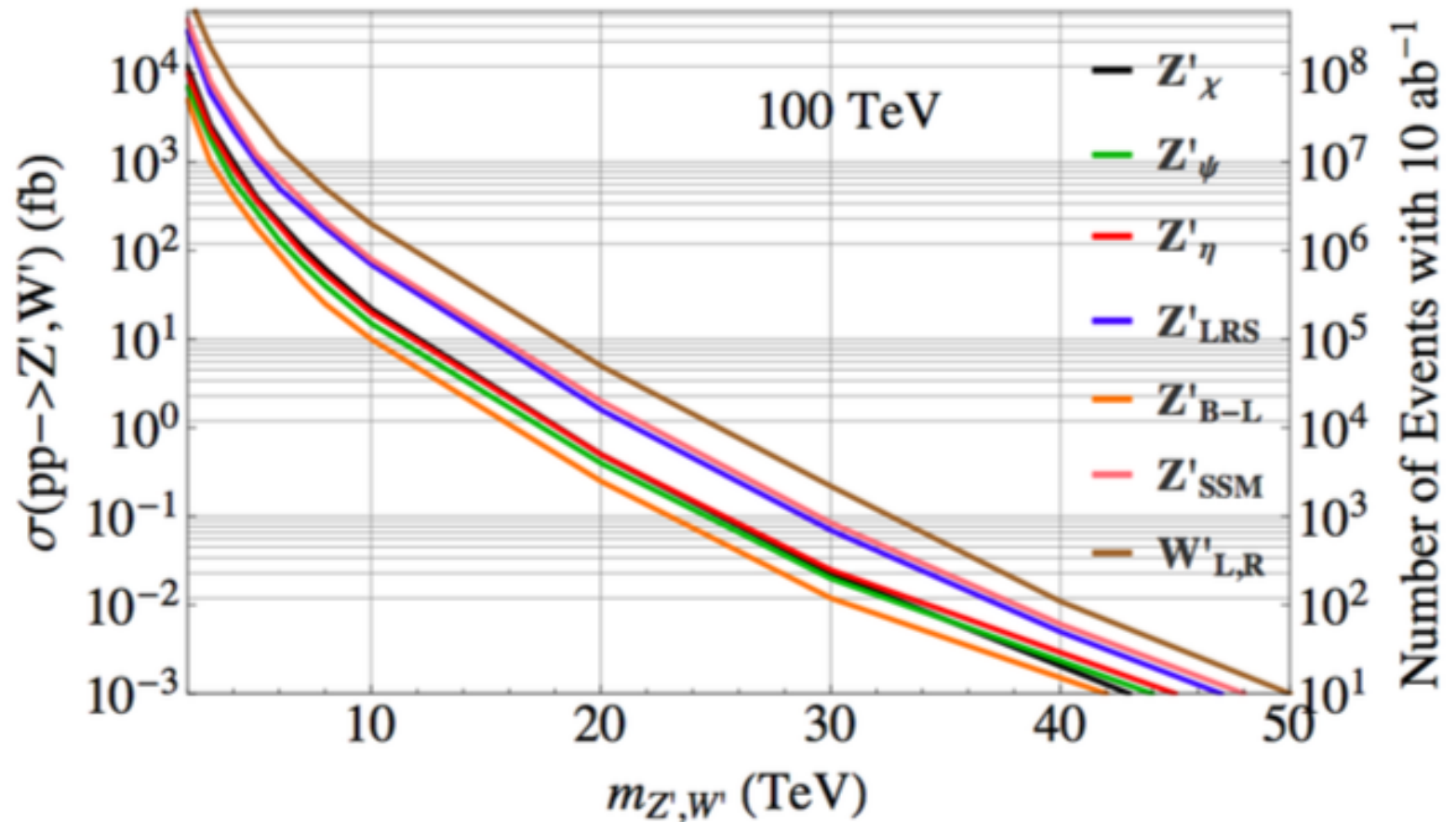
$$N_{14} = \sigma_{14\text{TeV}} \times 3 \text{ ab}^{-1}$$

Statistical precision:

- O(100 - 500) better w.r.t Run 1
- O(10 - 20) better w.r.t HL-LHC

Magnetic Tracking

Exploring New Territory - New Weak Gauge Interactions



Discovery reach

T.Rizzo, arXiv:1403.5465

Model	1 ab^{-1}	10 ab^{-1}	100 ab^{-1}
SSM	23.8	33.3	41.3
LRM	22.6	31.5	39.5
ψ	20.1	29.1	37.2
χ	22.7	30.6	38.2
η	20.3	29.8	38.0
I	22.4	29.2	36.2

10-fold increase in luminosity
 $\rightarrow \sim 7 \text{ TeV}$ increase in mass reach

Maintaining Fractional p_T Resolution

- Resolution gain with number of hits on track is slow (improves as \sqrt{N})
- Resolution improves linearly with $BL^2 \sim$ stored magnetic field energy in tracker
- Resolution improves linearly with hit resolution

Four tracker/magnet geometries being considered:

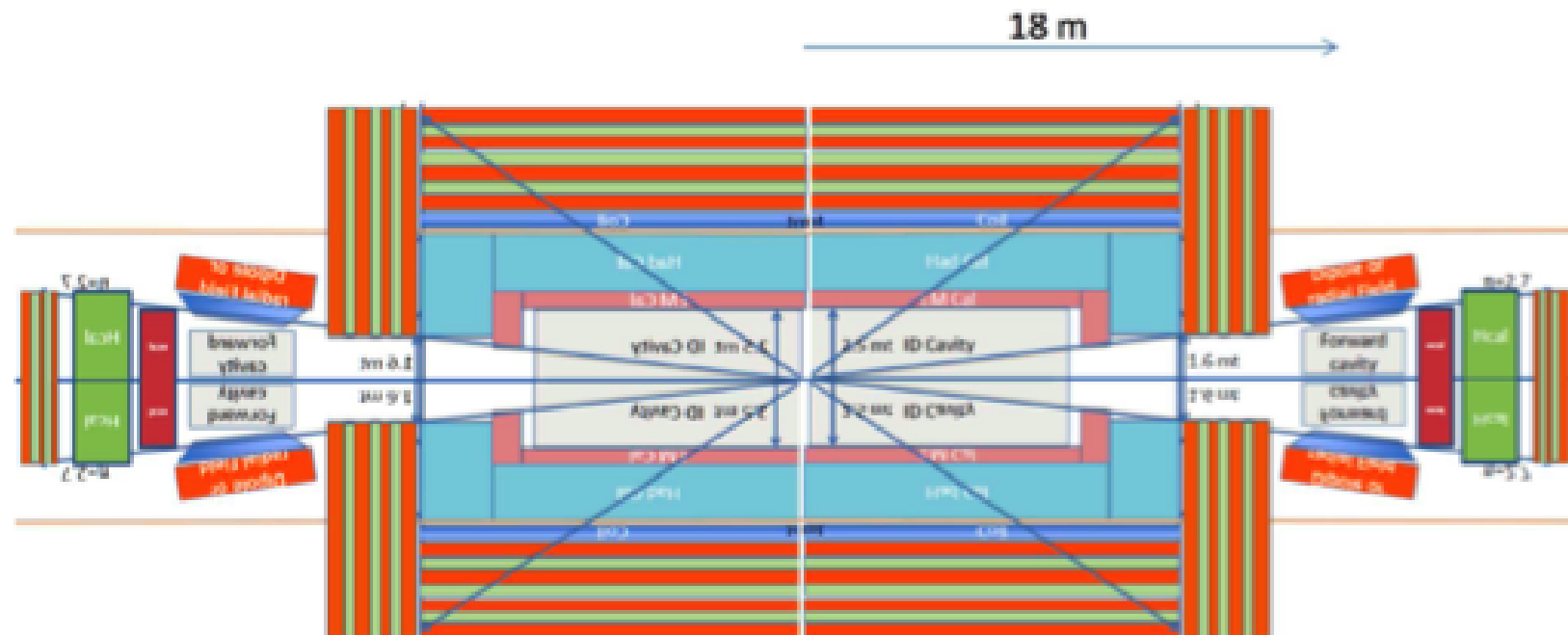
- see Dr. Marcello Mannelli's talk at Fermilab's “Next Steps in the Energy Frontier – Hadron Collider” Workshop

<https://indico.fnal.gov/conferenceOtherViews.py?view=standard&confId=7864>

Stored energy in the tracker magnetic field in the 50-100 GJ range (similar to ITER)



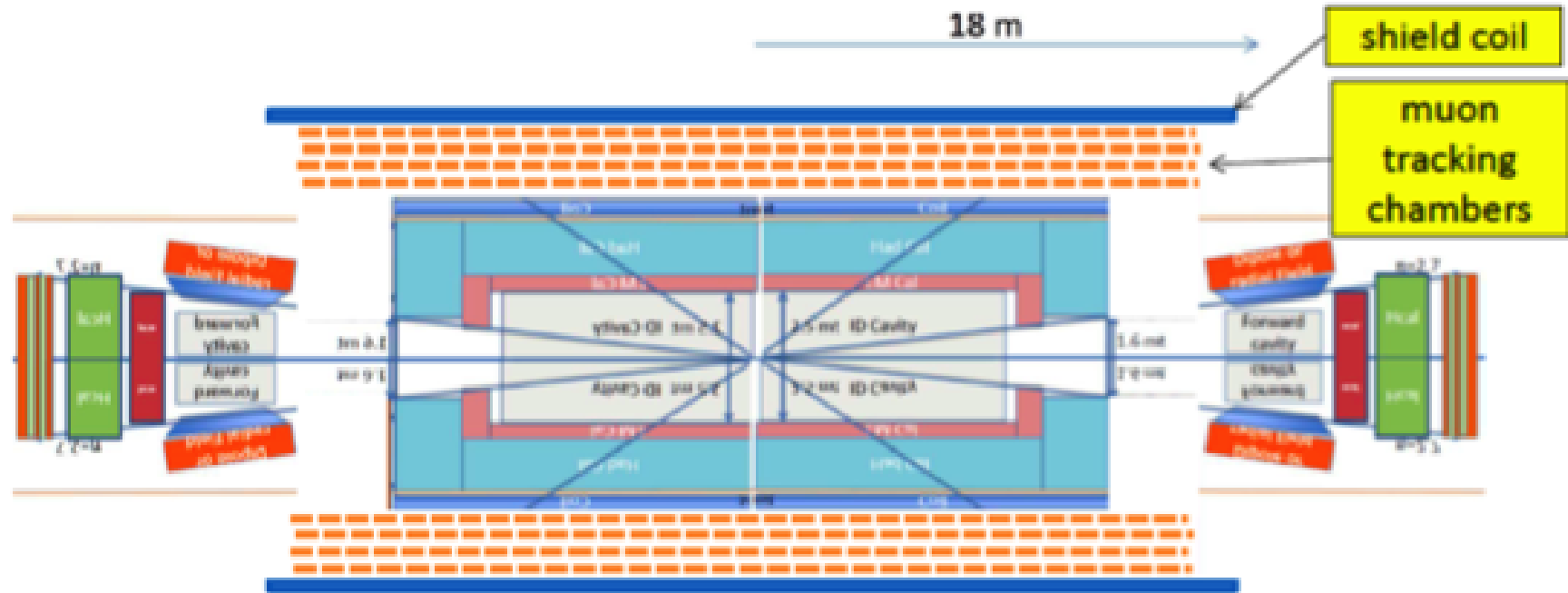
2. Option 1: Solenoid-Yoke + Dipoles (CMS inspired)



- ❖ **Solenoid:** 10-12 m diameter, 5-6 T, 23 m long
+ massive Iron yoke for flux shielding and muon tagging.
- ❖ **Dipoles:** 10 Tm with return yoke placed at $z \approx 18$ m.
Practically no coupling between dipoles and solenoid.
They can be designed independently at first.



2. Option 2: Twin Solenoid + Dipoles



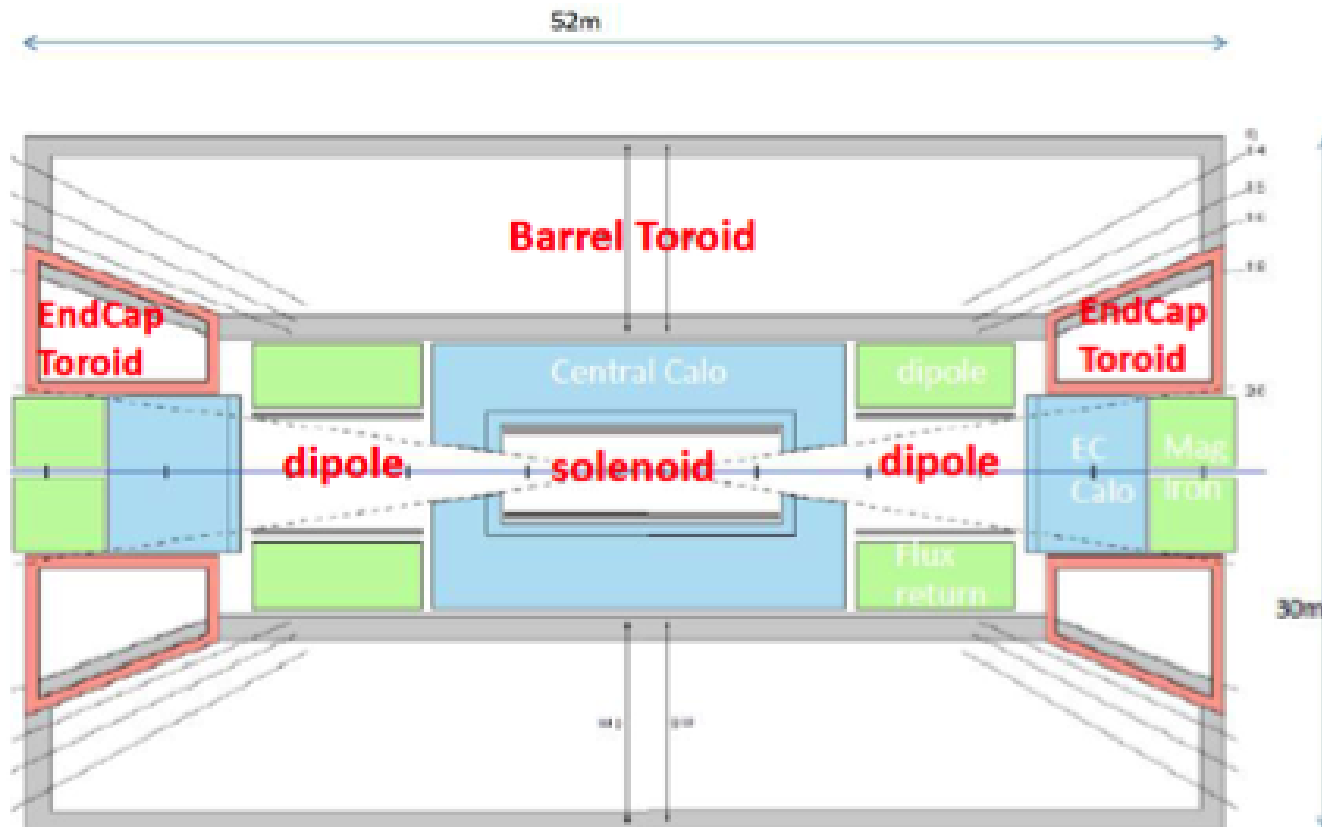
Twin Solenoid: a 6 T, 12 m dia x 23 m long main solenoid + an active shielding coil

Important advantages:

- ✓ **Nice Muon tracking space:** area with 2 to 3 T for muon tracking in 4 layers.
- ✓ **Very light:** 2 coils + structures, ≈ 5 kt, only $\approx 4\%$ of the option with iron yoke!
- ✓ **Much smaller:** system outer diameter is significantly less than with iron .



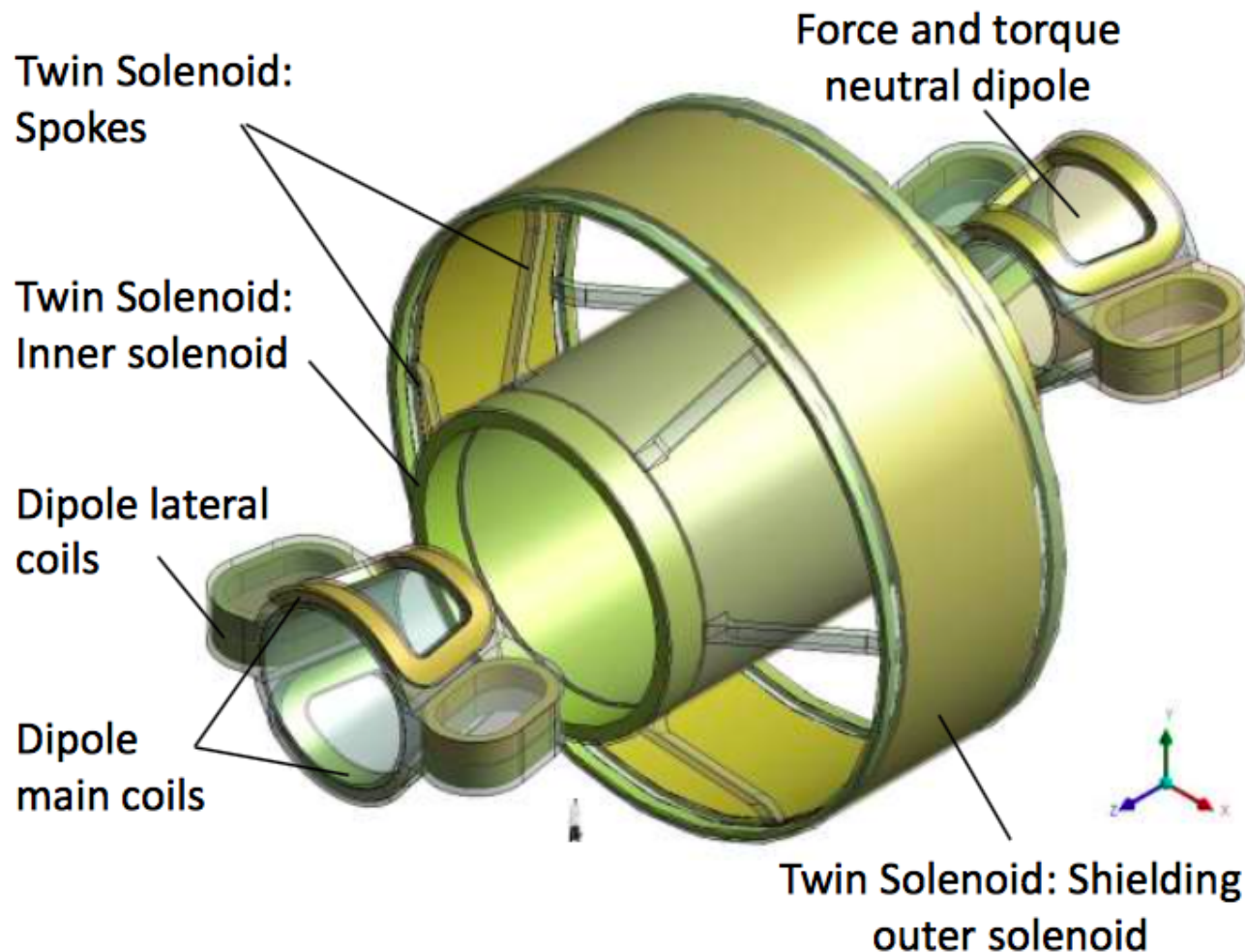
2. Option 3: Toroids + Solenoid + Dipoles (ATLAS +)



- ❖ 1 Air core Barrel Toroid with 7 x muon bending power $B_z L^2$.
- ❖ 2 End Cap Toroids to cover medium angle forward direction.
- ❖ 2 Dipoles to cover low-angle forward direction.
- ❖ Overall dimensions: 30 m diameter x 51 m length (36,000 m³).



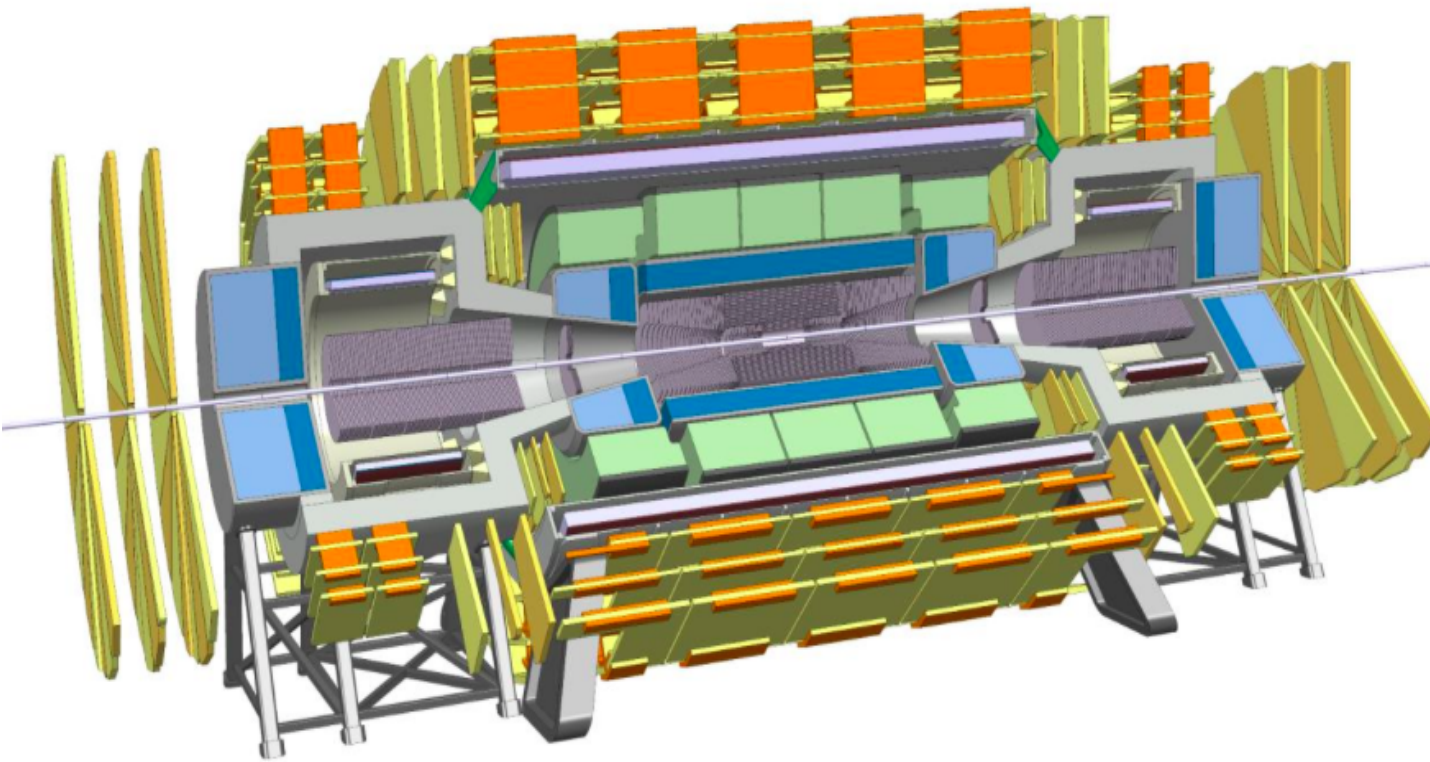
Twin Solenoid & Dipole system – bare coils



Property	Value
TS cold mass	3.2 kt
TS vacuum vessel mass	2.4 kt
TS stored energy	53 GJ
Dipoles cold mass	2x 380 t
Dipoles vac. vessel mass	To be det.
Dipoles stored energy	2x 1.5 GJ
Free bore	12 m
Outer diameter	27 m
System length	42 m
Total stored energy	56 GJ

(from Herman ten Kate)

Reference detector for the CDR



- 4T 10m solenoid
- Forward solenoids
- Silicon tracker
- Barrel ECAL LAr
- Barrel HCAL Fe/Sci
- Endcap HCAL/ECAL LAr
- Forward HCAL/ECAL LAr

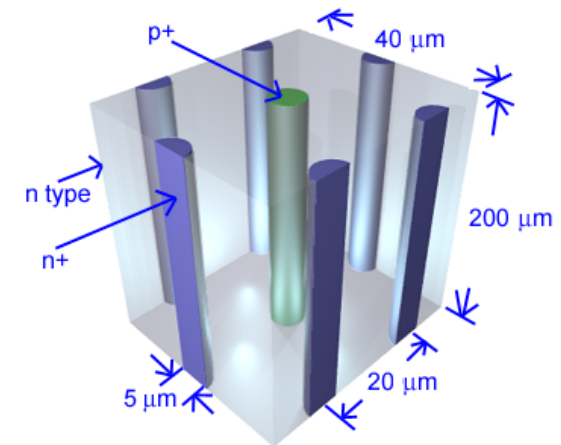
This is a reference detector that 'can do the job' and that is used to define the challenges. The question about the specific strategy for detectors at the two IPs is a different one.

Skip outer coil for baseline cost estimates...

(from Werner Riegler)

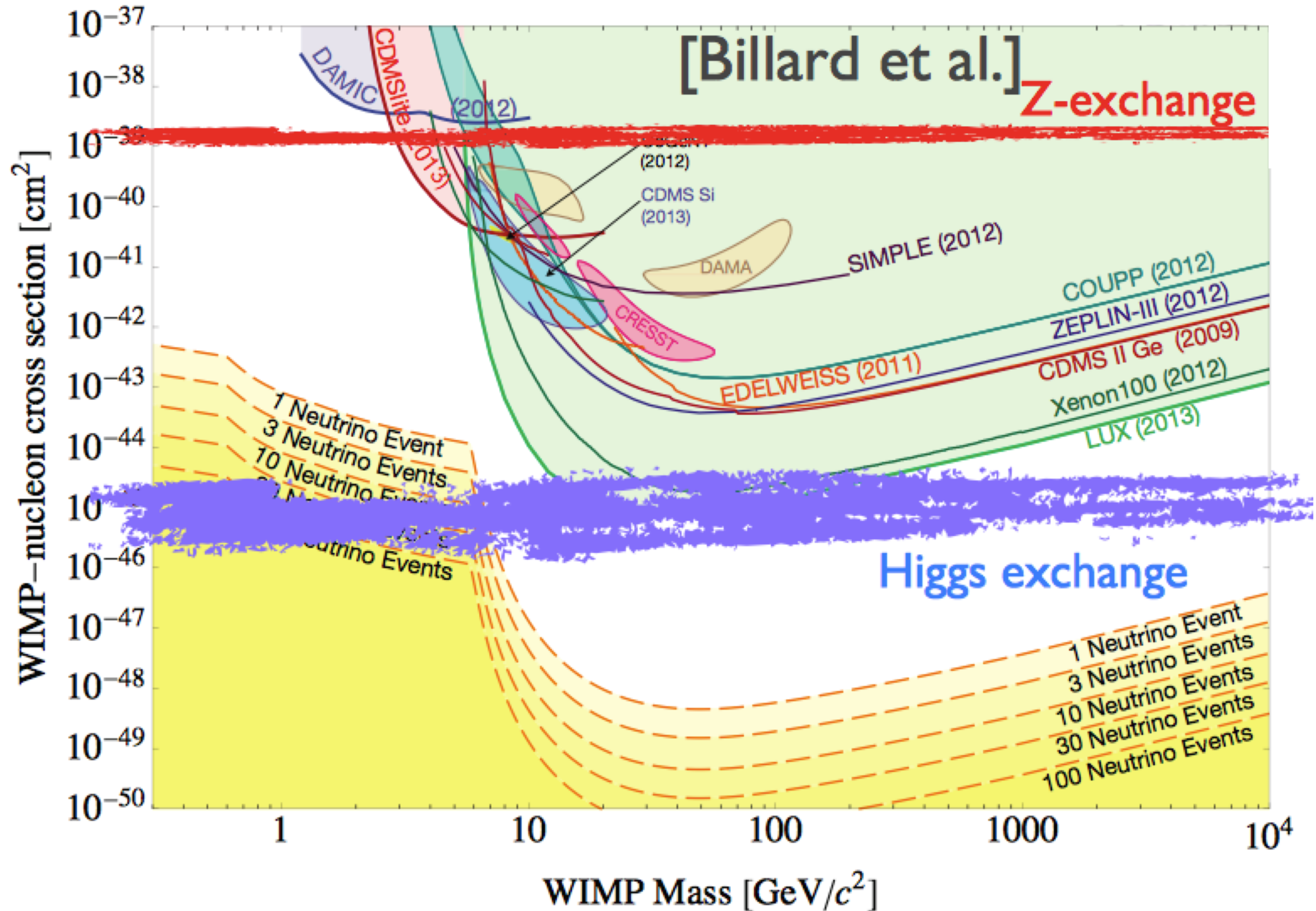
Improving Hit Resolution

- Smaller pixels with silicon sensors have multiple advantages
 - Improved hit resolution linearly improves momentum resolution at high p_T
 - Higher granularity improves two-track resolving power
 - Helps resolve close-by tracks and maintain track reconstruction efficiency in
 - high-density environment (inside boosted jets)
 - High-occupancy environment (pileup at high L)
- Issues:
 - Higher readout rate required
 - Power may be dominated by inter-pixel capacitance, which does not reduce with pixel size
 - More pixels \Rightarrow more power
- Potential solutions (3D electronics etc) under discussion



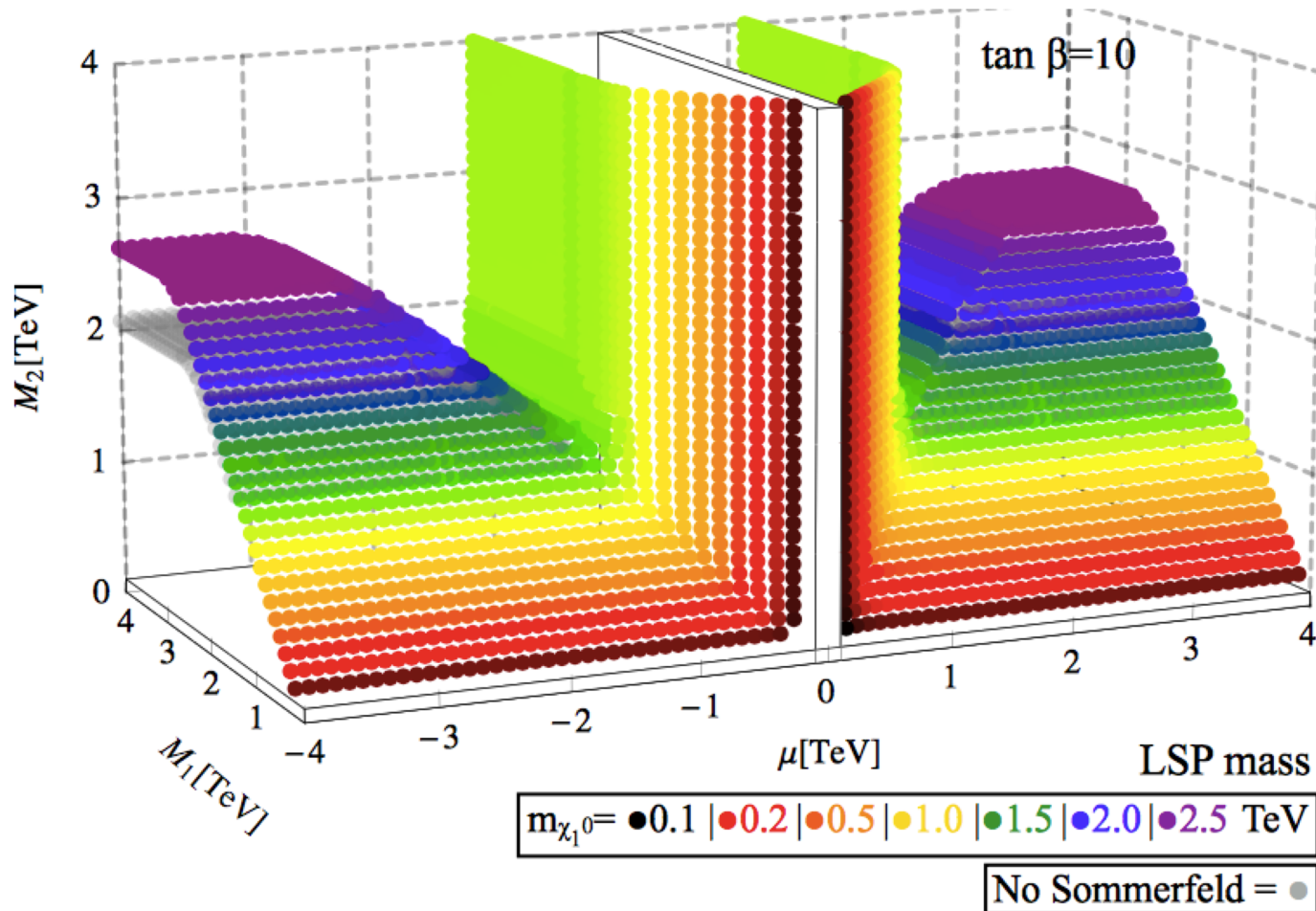
Dark Matter

Direct Searches for Dark Matter



SUSY Neutralino WIMP Relic Surface

- Supersymmetric partners of photon, Z boson or Higgs boson provide generic model of weakly interacting Dark Matter
- Combinations of Neutralino mass parameters that produce the correct relic abundance, along with Dark Matter particle (LSP) mass



Bramante *et al*,

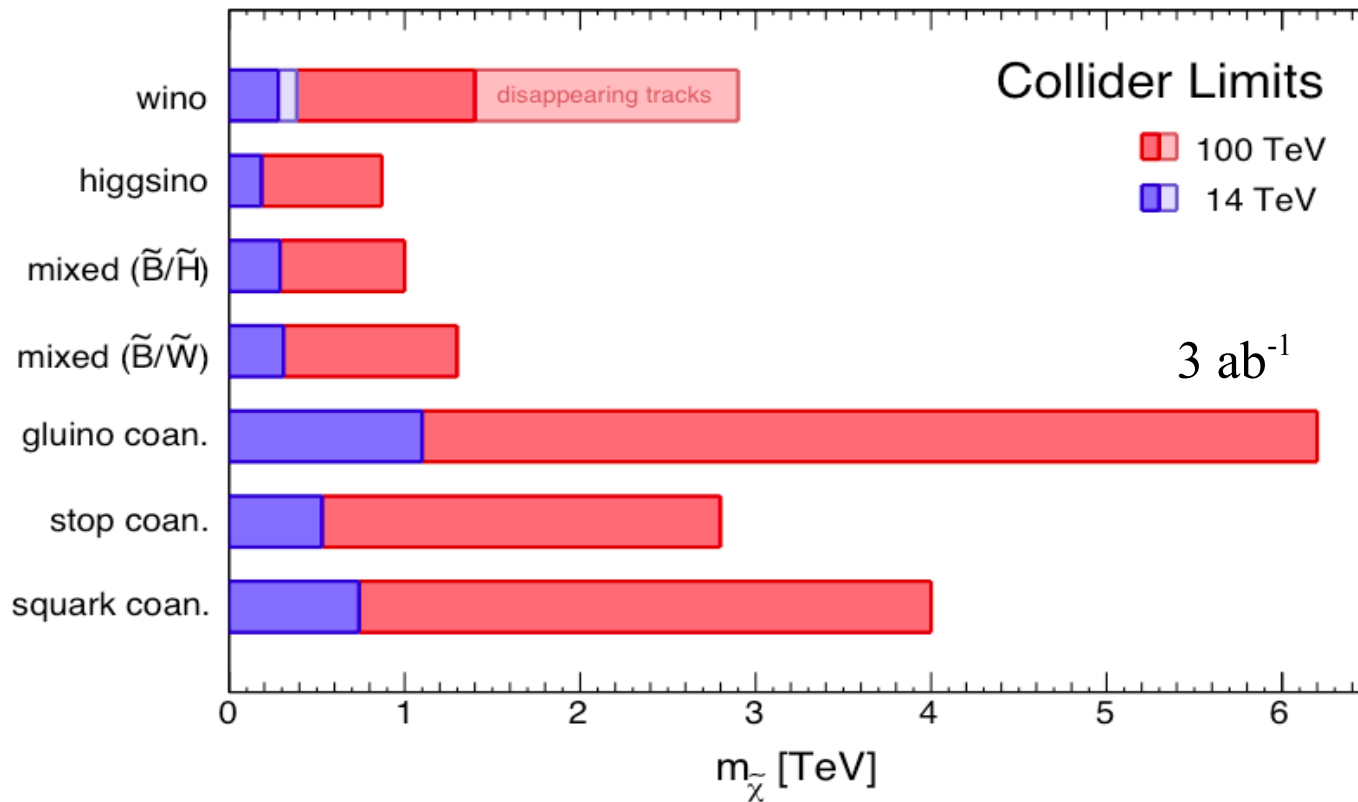
ArXiv:1510.03460

**Phys. Rev. D91 (2015)
054015**

(in the limit that other SUSY is heavy and decoupled)

Disappearing Track from Wino WIMP Decay

- $M_{\text{Dark Matter}} < 1.8 \text{ TeV } (g_{\text{DM}}^2/0.3)$ based on WIMP thermal relic hypothesis



M. Low, L-T Wang,
ArXiv:1404.0682
(mono-jet channel)

100 TeV pp collider covers most of the parameter space – 30 ab^{-1} will double the mass reach

Disappearing track: almost degenerate, long-lived $\text{Wino}^+ \rightarrow \text{Wino}^0$
requires robust tracking for reconstructing partial-length tracks

Compressed Spectrum WIMPs

$$pp \rightarrow (\tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0) (\tilde{\chi}_1^\pm \rightarrow \ell^\pm \nu_\ell \tilde{\chi}_1^0) j \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 \ell^\pm \nu_\ell \gamma j$$

Bramante *et al*, **Phys. Rev. D**93 (2016) no.6, 063525

$$\begin{array}{lll} p_{T,\ell} = [10 - 60] \text{ GeV} & |\eta_\ell| < 2.5 & \\ p_{T,\gamma} = [10 - 60] \text{ GeV} & |\eta_\gamma| < 2.5 & \Delta R_{\ell\gamma} > 0.5 \\ p_{T,j} > 0.8 \text{ TeV} & |\eta_j| < 2.5 & M_{T2}^{(\gamma,\ell)} < 10 \text{ GeV} \\ \cancel{p}_T > 1.2 \text{ TeV} . & & \end{array}$$

Soft leptons and photons are crucial for this signature

Collider vs Direct Detection Complementarity

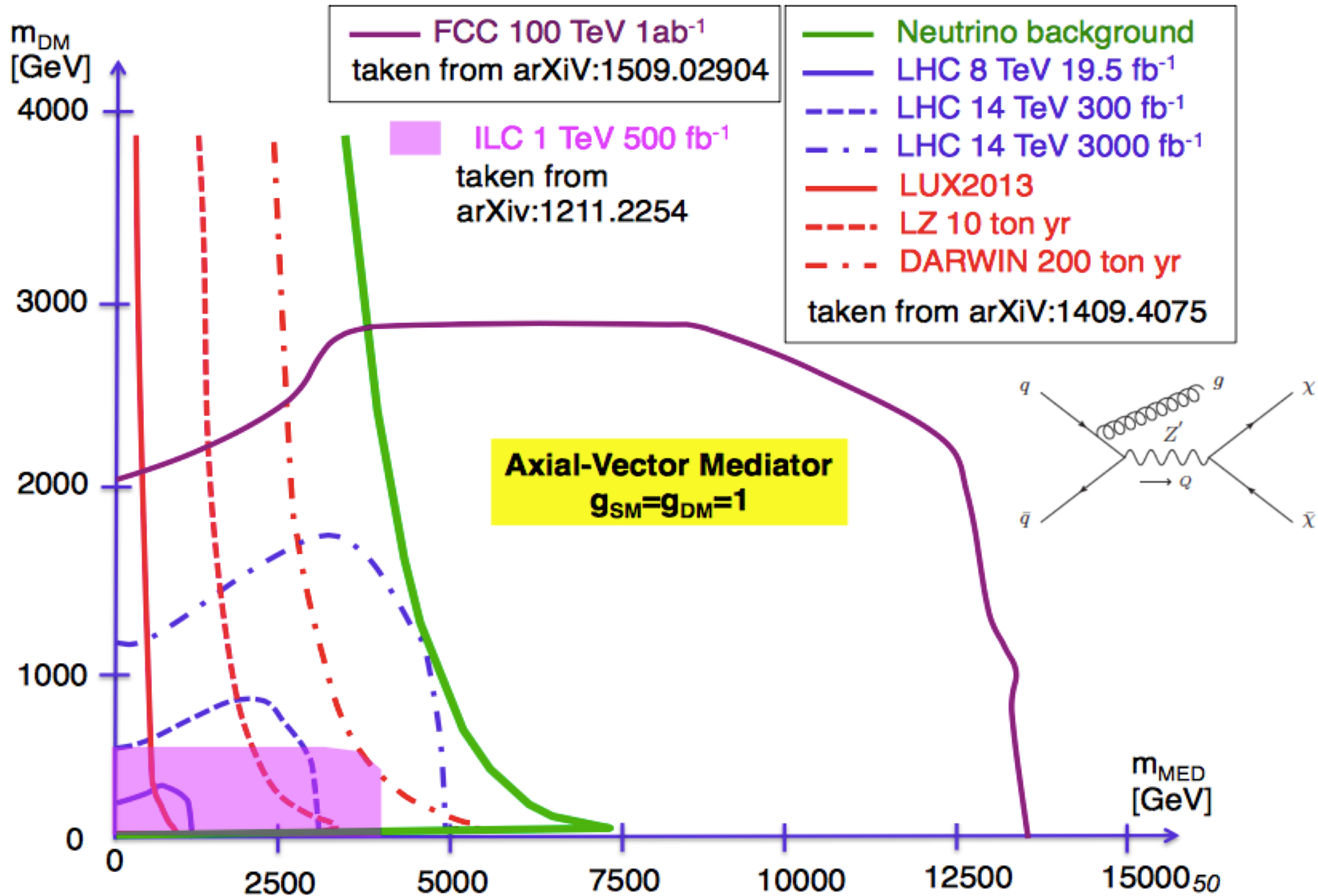
Common ground (almost)

- **Axial-Vector mediator**
DD and collider are equal in overall sensitivity but probe different regions of parameter space!
- **Scalar mediator**
DD and collider are equal in overall sensitivity but probe different regions of parameter space!

Exclusive domains (almost)

- **Vector mediator**
Besides very low DM masses DD wins clearly over collider
- **Pseudo-Scalar mediator**
No competitive limits from DD (only from indirect detection). Collider provides limits similar in sensitivity to scalar limits

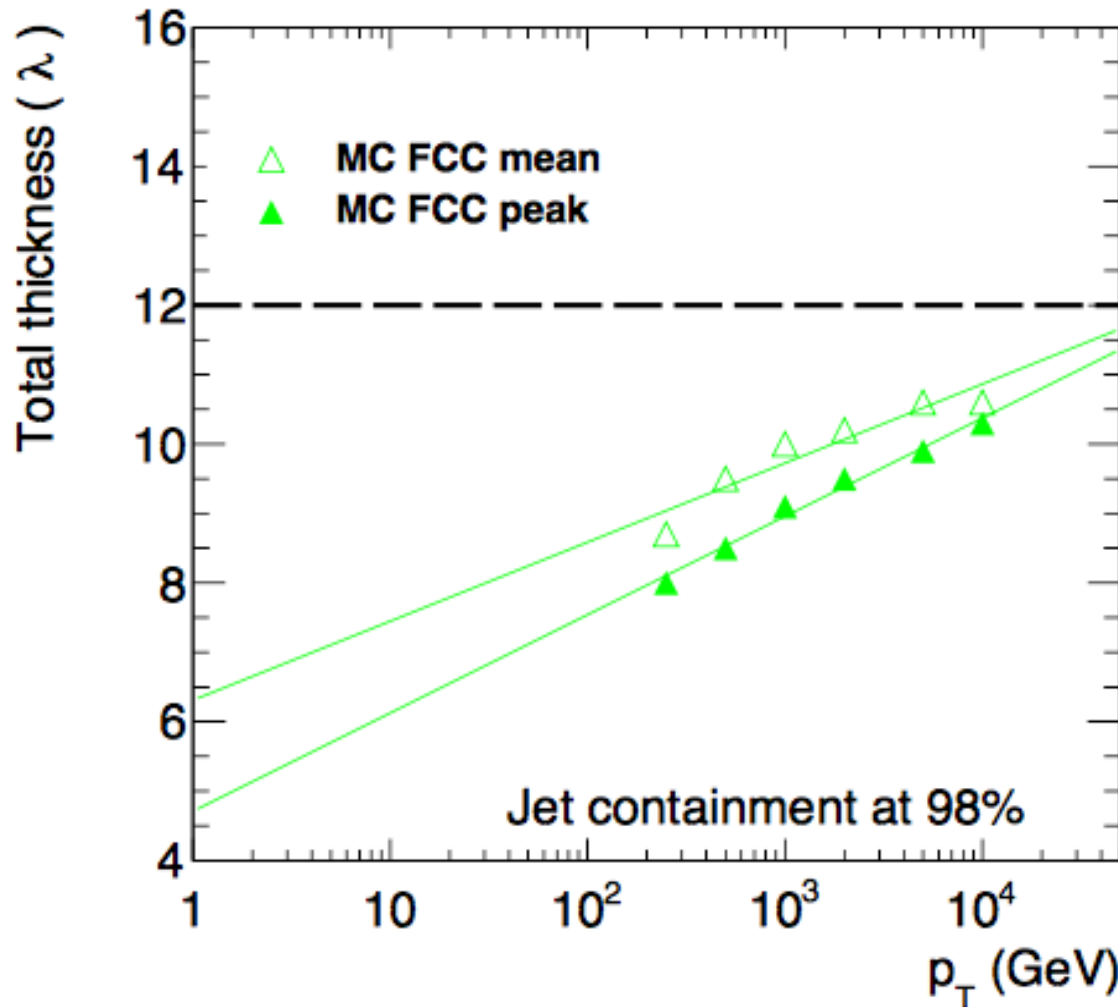
Collider Searches – Large Mediator Mass



Calorimetry

Calorimeter Geometry Issues

- Conveniences for going to higher energy:
 - Shower depth for full containment grows as $\log(E)$
 - Energy resolution improves as \sqrt{E}

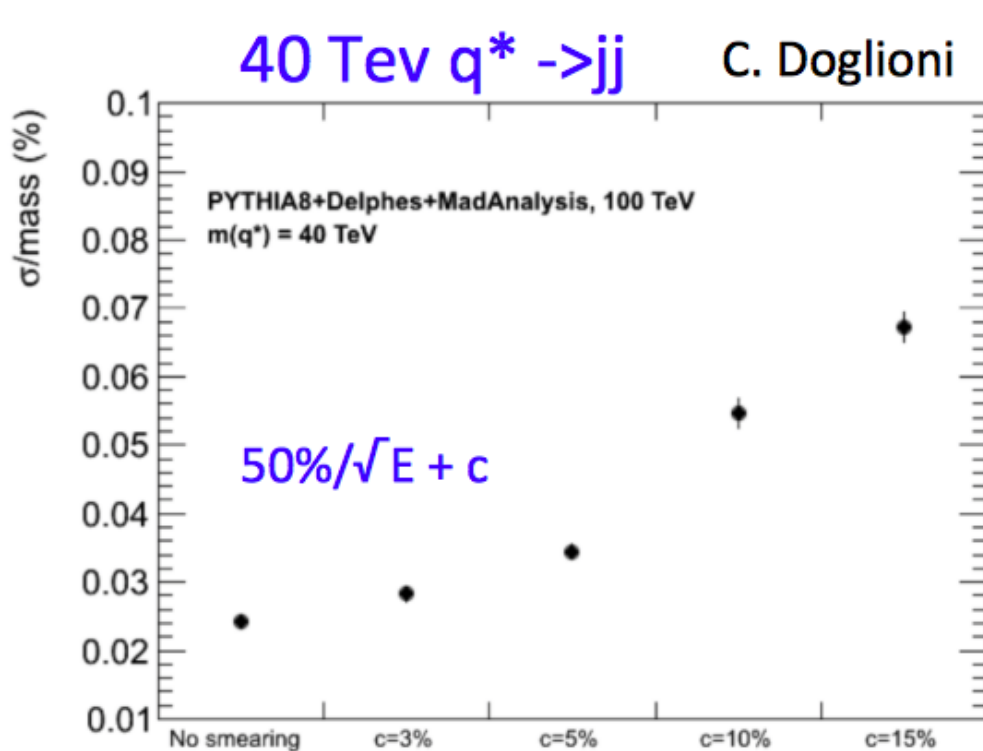


T. Carli *et al*,
arXiv:1604.01415

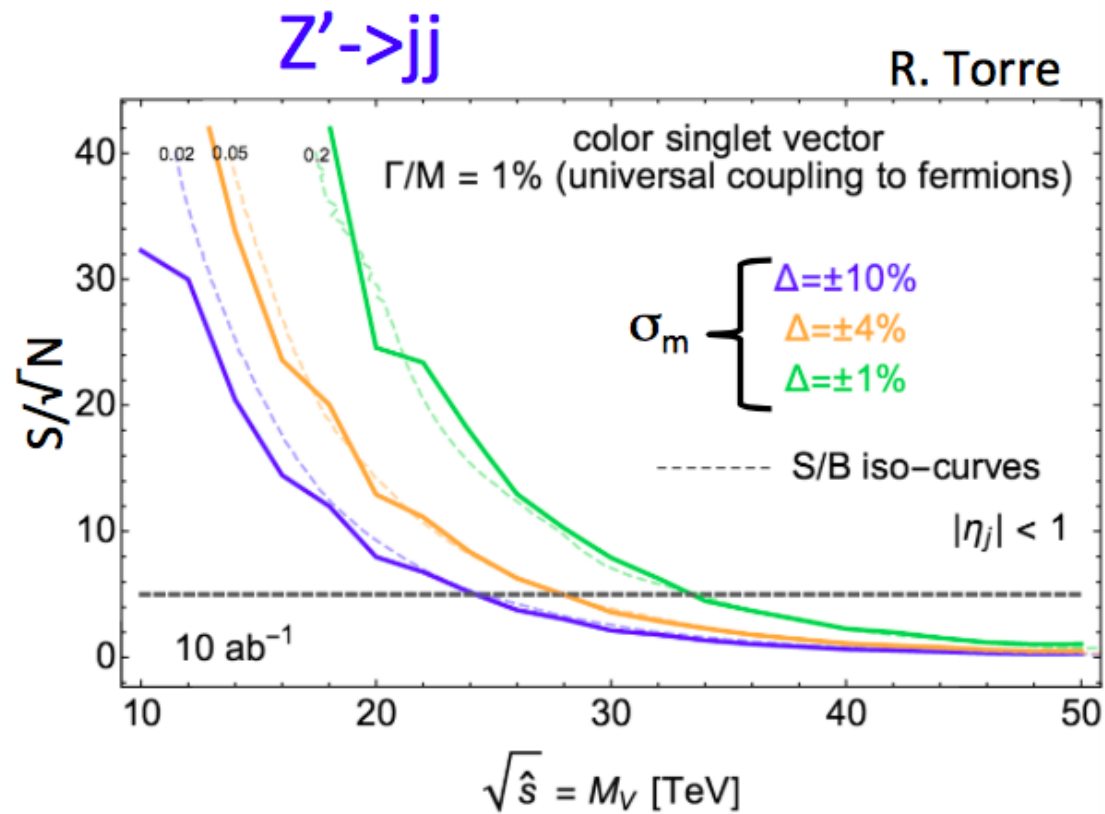
11-12 interaction lengths
needed – space constraints
(coil radius is expensive)

- Dynamic range of electronics readout required scales linearly with collider energy

Effect of HCAL Energy Resolution on Dijet Resonances



HCAL resolution constant term c



Jet resolution $\sim 2\text{-}3\%$ needed for multi TeV dijet resonances

- Extend $Z' \rightarrow jj$ discovery potential by 10TeV between $\sigma_m = 10\%$ to 1%
- Constant term will dominate at TeV energies ($\sigma/E = a/\sqrt{E} \oplus c$)
- Good shower containment is mandatory!

(from Ana Henriques)

Calorimeter Granularity

- Granularity is a KEY issue: all decay products will be boosted closer together
 - 5 TeV resonance \rightarrow HH \rightarrow 4 τ produces 1 TeV τ -lepton
 - Photons within τ -jet are separated by ~ 2 mm
 - τ -leptons from Higgs separated by ~ 10 cm
 - 20 TeV resonance \rightarrow tt , top decay products separated by ~ 3 cm
 - 10 TeV Zprime \rightarrow WW, boosted W \rightarrow jets separated by ~ 3 cm
- Tracking particles inside jets can be crucial
- Exploit particle flow algorithms to the fullest, push experience from CMS and ILC detector design effort

Geant4 simulation of a high-granular calorimeter for TeV-scale boosted particle

S. Chekanov
HEP/ANL

FCC Week. April 11-15, 2016
Rome, Italy

With contributions from:

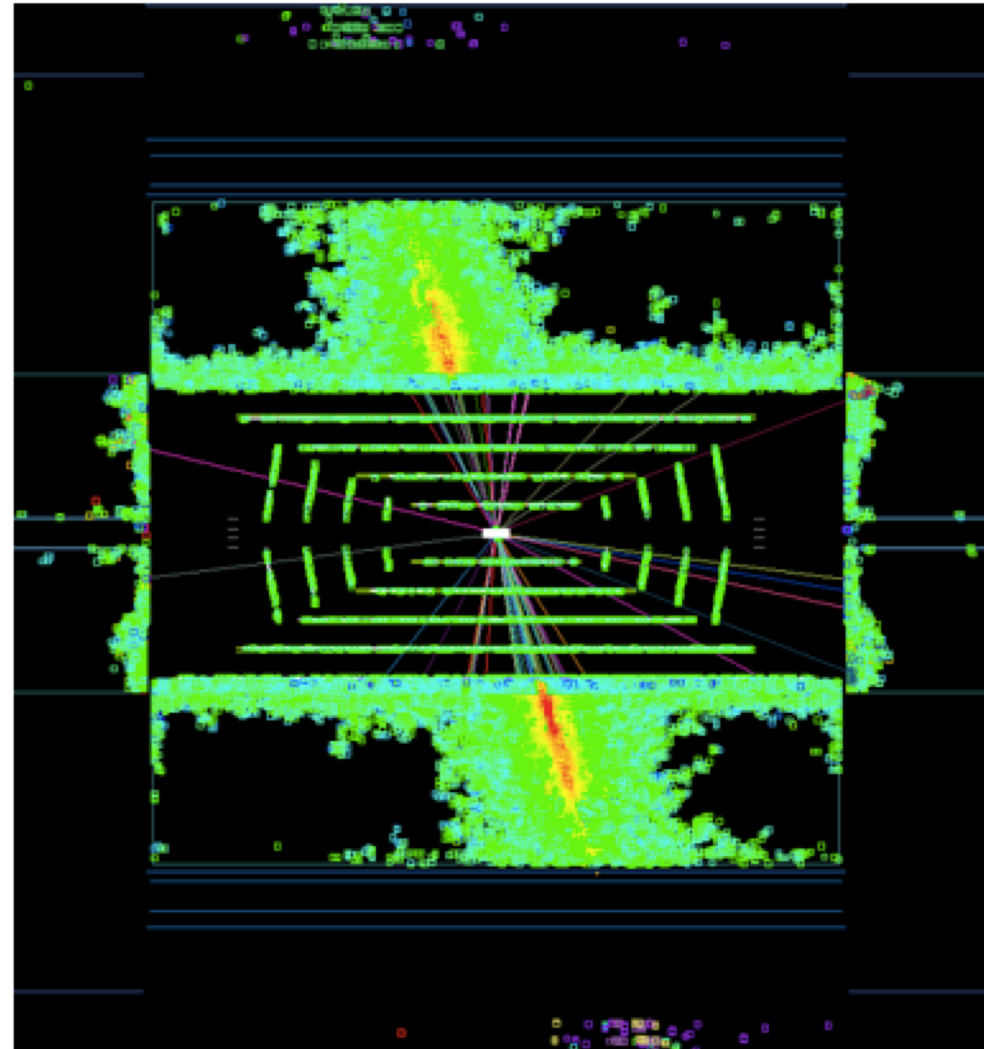
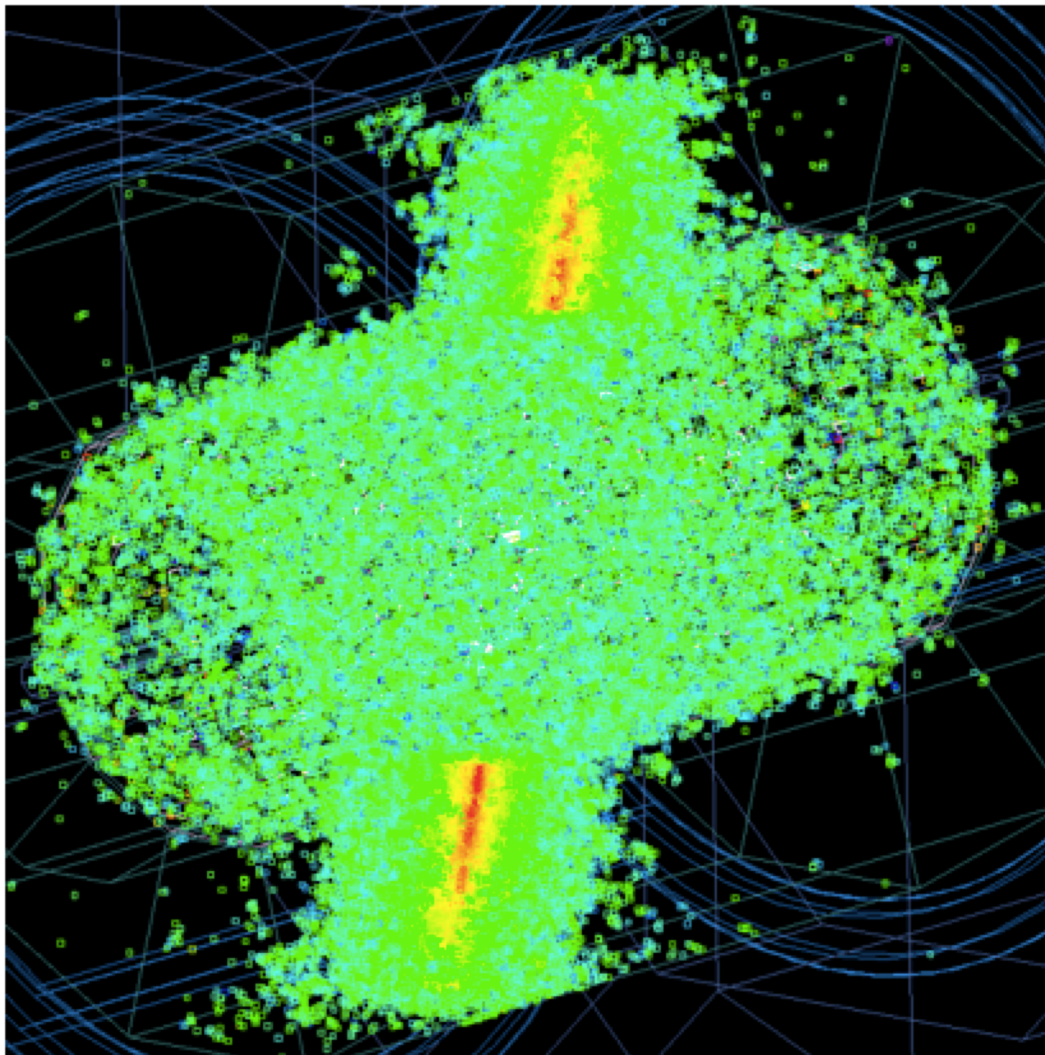
A.Kotwal (Fermilab/Duke), L.Gray (Fermilab), J.Strube (PNNL), N.Tran (Fermilab), S. Yu (NCU), S.Sen (Duke), J.Repond (ANL), J.McCormick (SLAC), J.Proudfoot (ANL), A.M.Henriques Correia (CERN), C.Solans (CERN), C.Helsens (CERN)

See Sergei Chekanov's talk in BOOST2017

GEANT Simulation of Scintillator / Iron HCAL and Silicon Tracker

5 TeV hadronic $W \rightarrow$ dijet decay with 4 cm x 4 cm scintillator readout

Background simulation in progress, will investigate different pad sizes and higher p_T



Generated on OSG by S. Chekanov

GEANT Simulation of Silicon/Tungsten EM Calorimeter

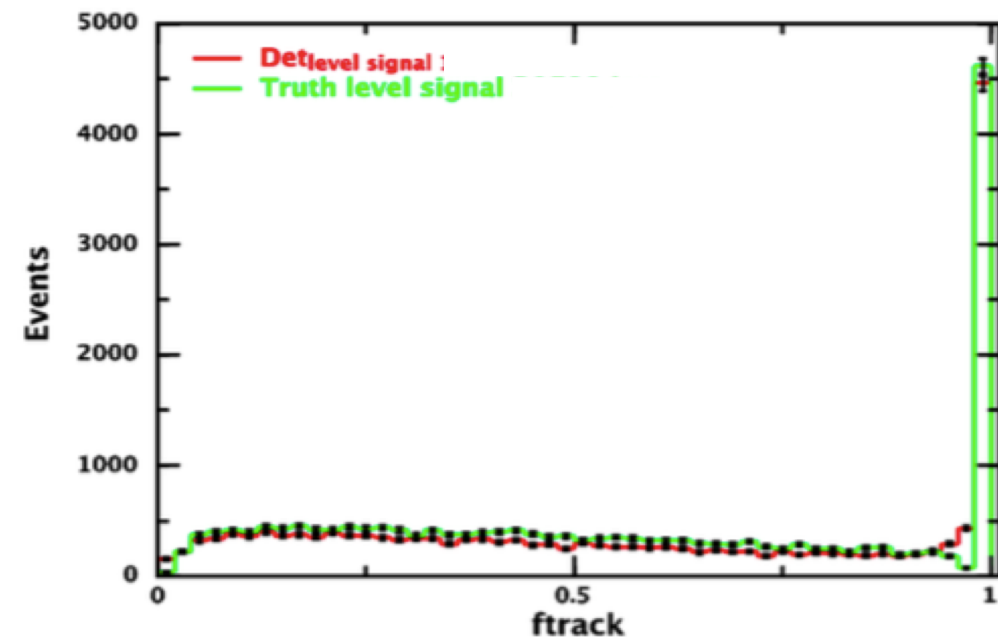
500 GeV hadronic τ -lepton decays with 4mm x 4mm silicon pads

Background simulation in progress, will investigate larger pad sizes and higher p_T

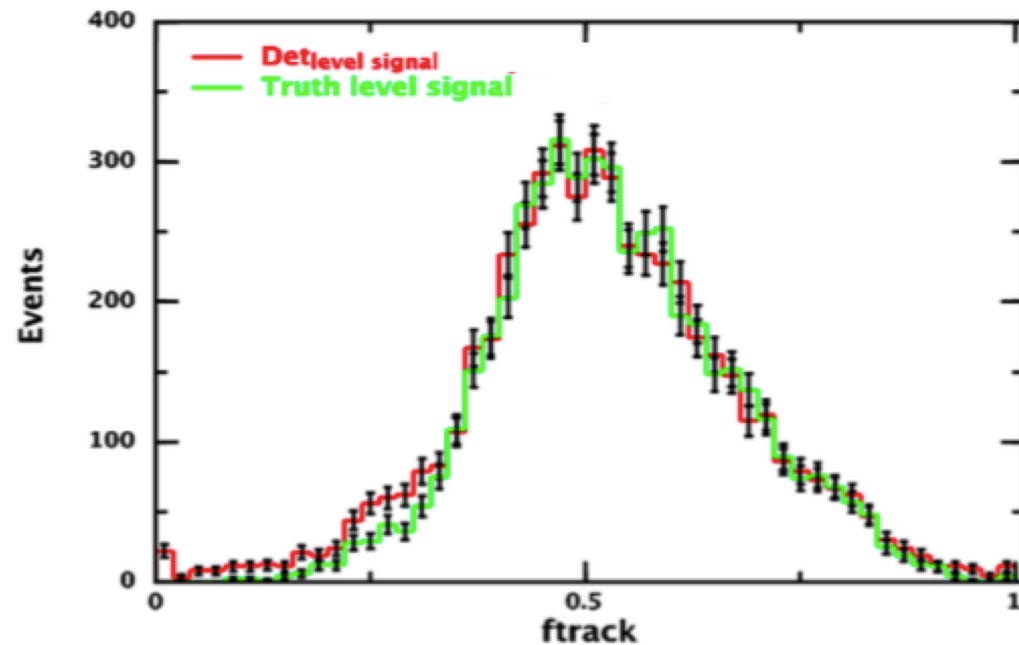
f_{track} (leading track momentum fraction)

$= (\text{pT of highest pT track in core region } (\Delta R < \text{core})) / (\text{Total } E_T \text{ deposited in } \Delta R < \text{core})$

core = 0.1



1 prong



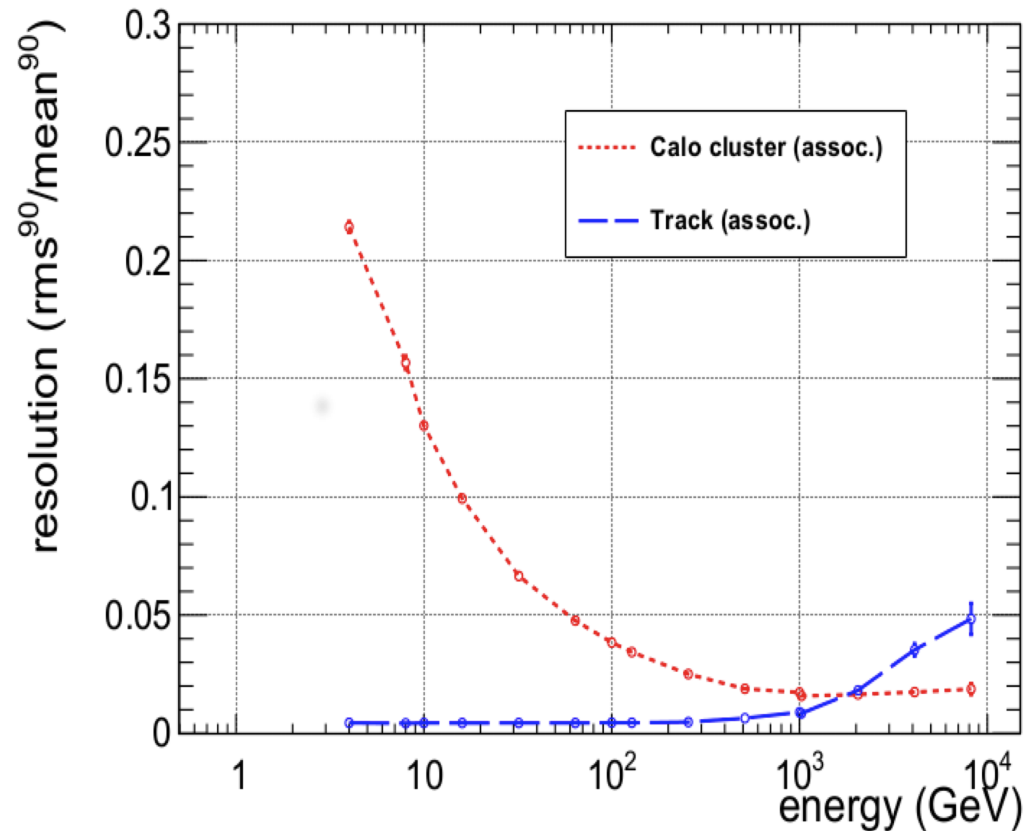
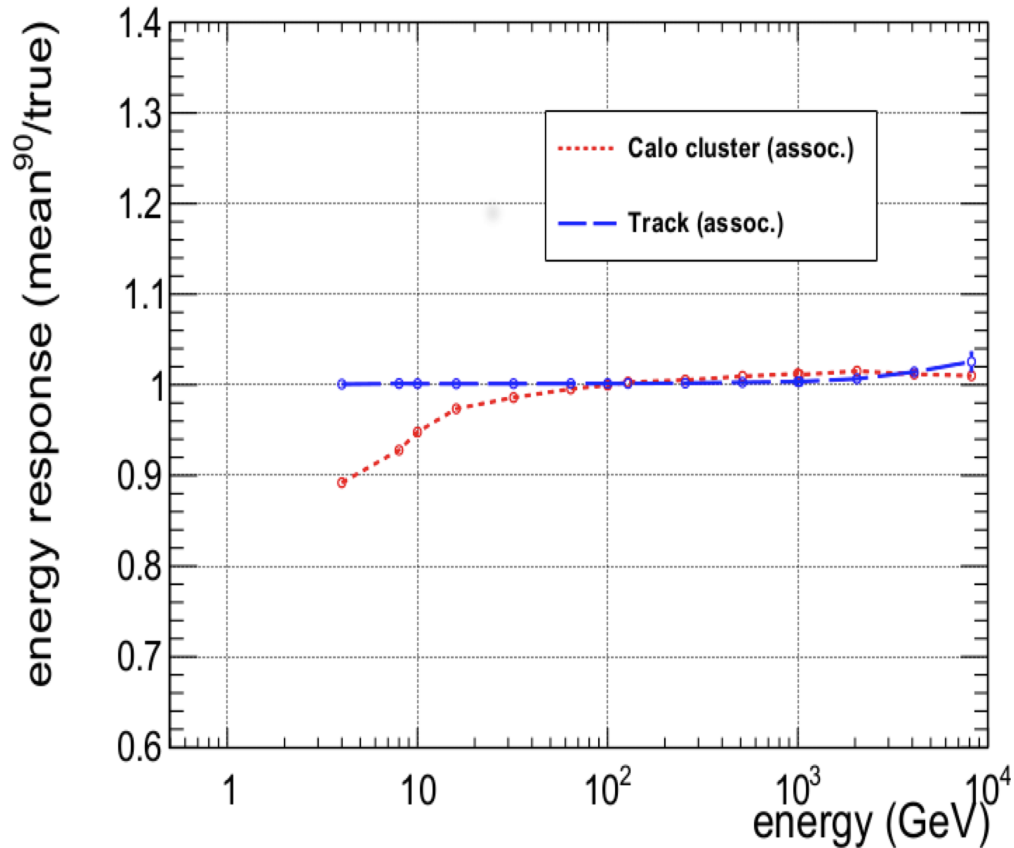
3 prong

Analysis by Sourav Sen (Duke graduate student)

Higgs $\rightarrow \tau\tau$ is an important channel to complement $\gamma\gamma$ and $b\bar{b}$

GEANT Simulation: Si/W ECAL & Scintillator/Iron HCAL

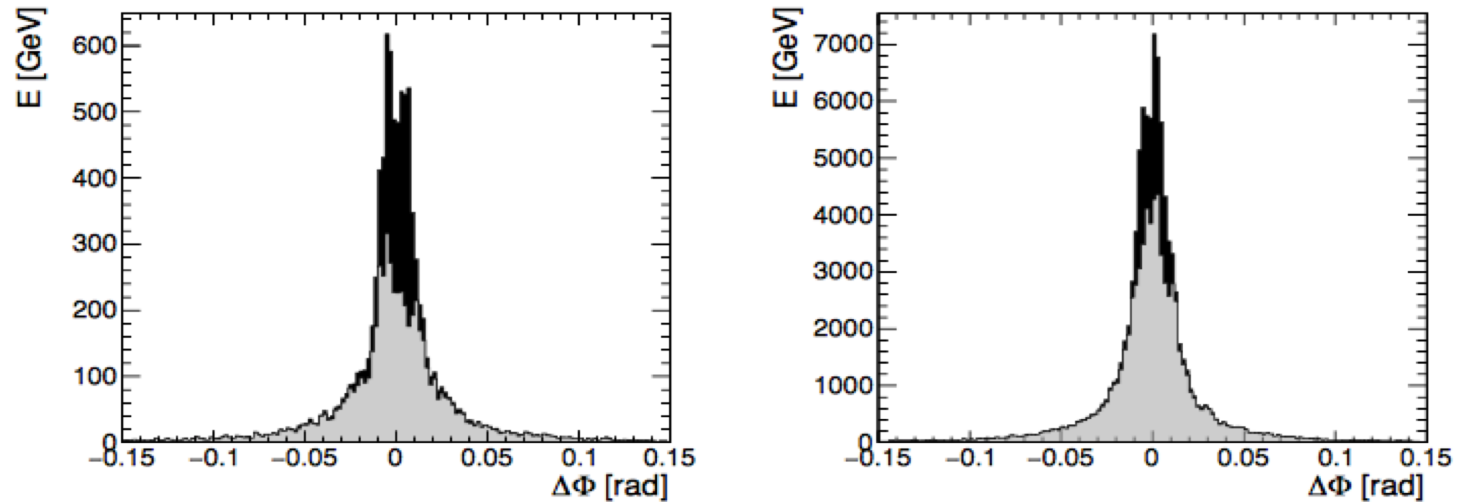
Single pion response and resolution



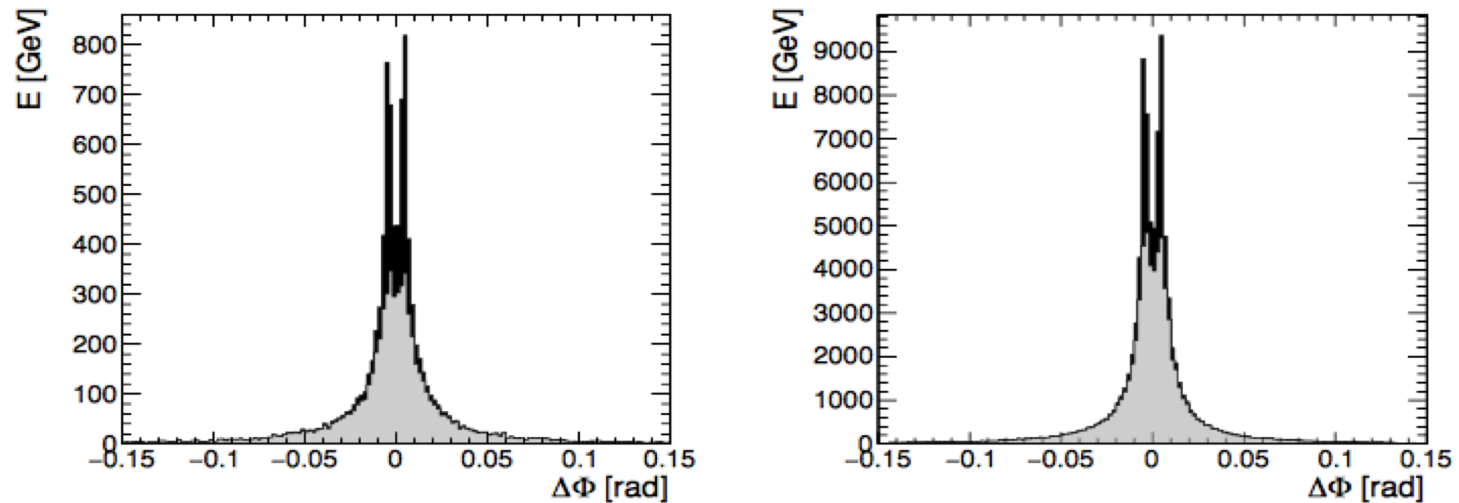
- Analysis by S. Yu, N. Tran and S. Chekanov
- First look at boosted object discriminating variables
- Published in **JINST 12 (2017) no.06, P06009**

GEANT Simulation: Silicon/Tungsten EMCAL & Iron/Scintillator HCAL

Dual K_L^0 spatial separation (generated $\Delta\phi = 10$ mrad)



(b) 5×5 cm HCAL cells and 2×2 cm ECAL cells



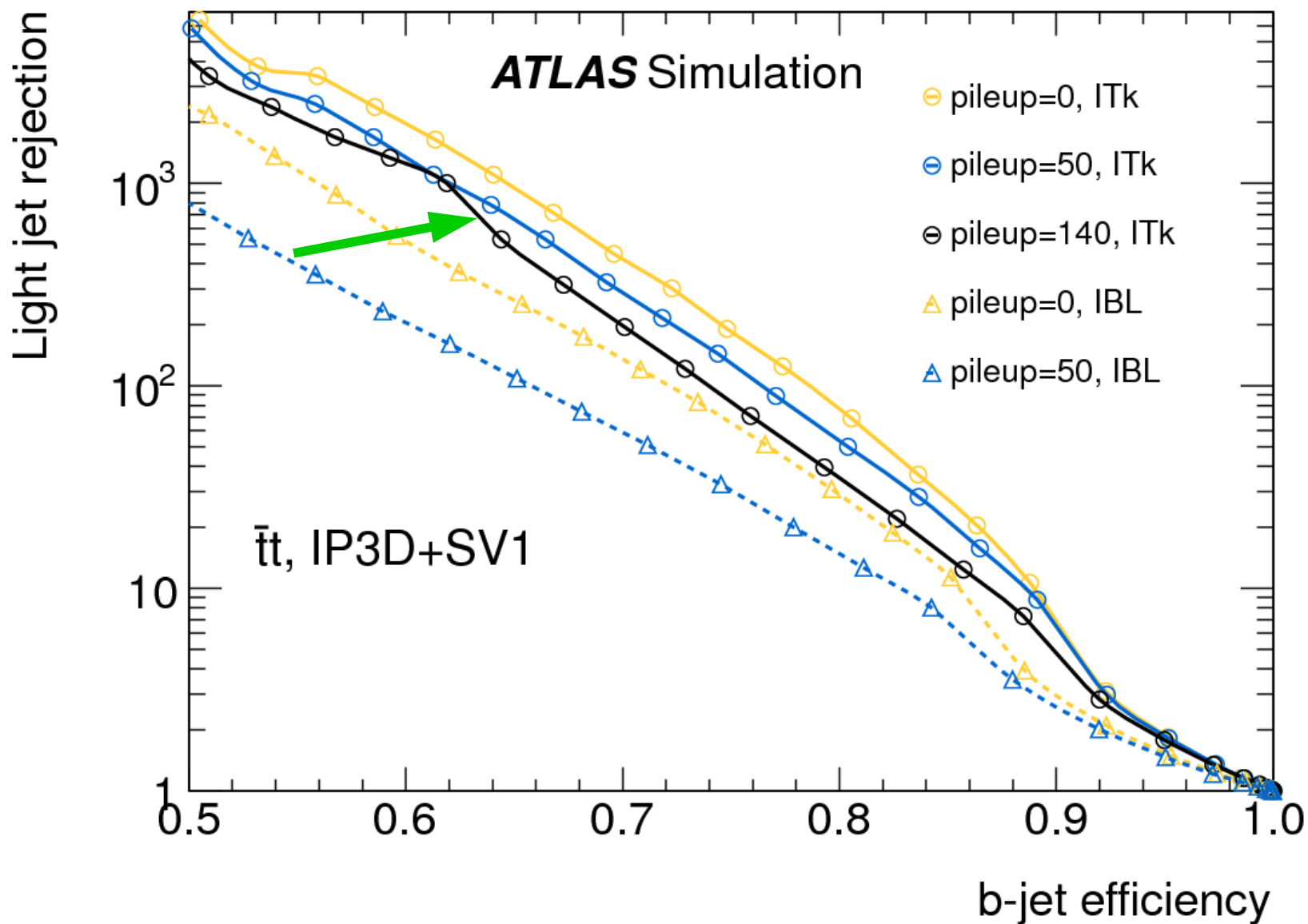
(c) 1×1 cm HCAL cells and 3×3 mm ECAL cells

- Analysis by Nhan Tran
- Published in **JINST 12 (2017) no.06, P06009**

Figure 14: Azimuthal distribution of energy deposition for pair of incident K_L^0 particles at 100 GeV (left) and 1000 GeV (right), with the angular separation of $\Delta\phi^K = 0.009$ rad. Electromagnetic calorimeter cells are indicated in black while hadronic calorimeter cells are indicated in gray.

b-tagging

b-tagging Design Performance for HL-LHC



IBL = current, ITk = HL-LHC design (3 \rightarrow 4 pixel layers, smaller pixels)

b-tagging

- FCC stage 1 plans to deliver $\sim 3 \text{ ab}^{-1}$
 - Similar conditions as HL-LHC, pileup ~ 200 at 25 ns bunch crossing
- FCC stage 2 plans to deliver $\sim 15 \text{ ab}^{-1}$
 - Pileup ~ 1000
 - or 5 ns bunch crossing? If very fast detectors have no out-of-time pileup
- Need to achieve same *b*-tagging performance in higher-density environments
 - Highly boosted top quarks and Higgs bosons from heavy resonance decays
 - Width of b-jet ~ 300 microns at 2 cm radius
 - Need to resolve tracks with factor x5 higher local density than LHC

Forward rapidity coverage

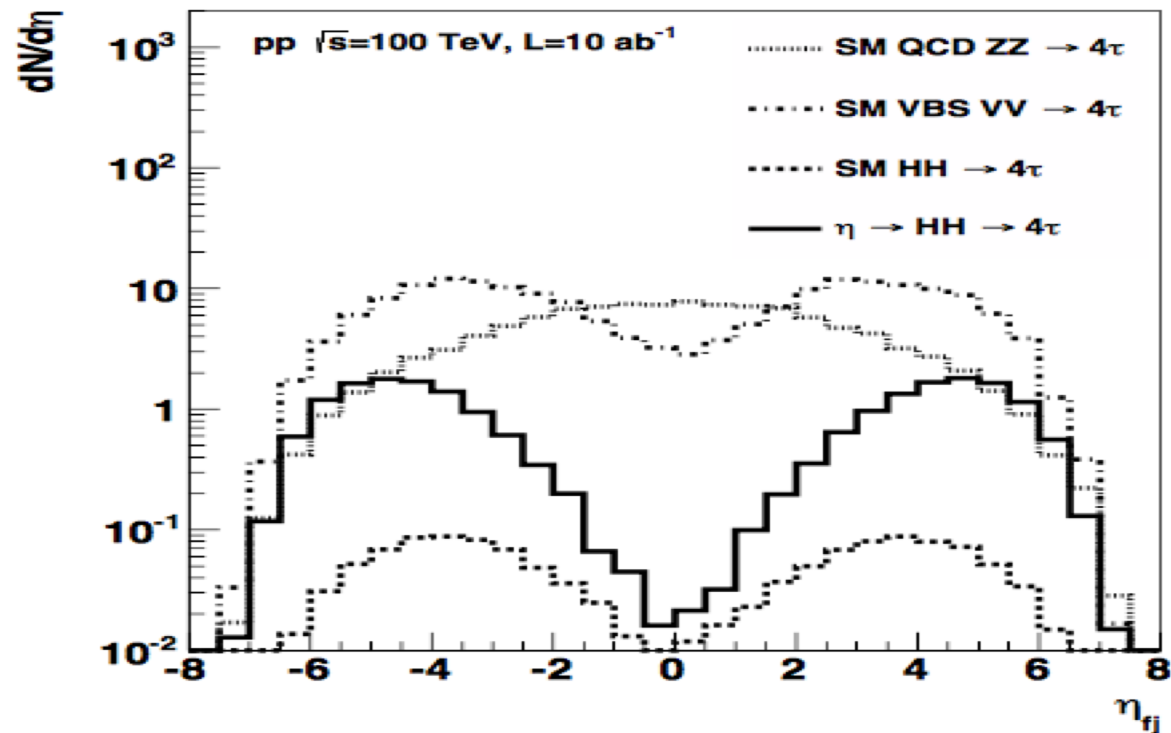
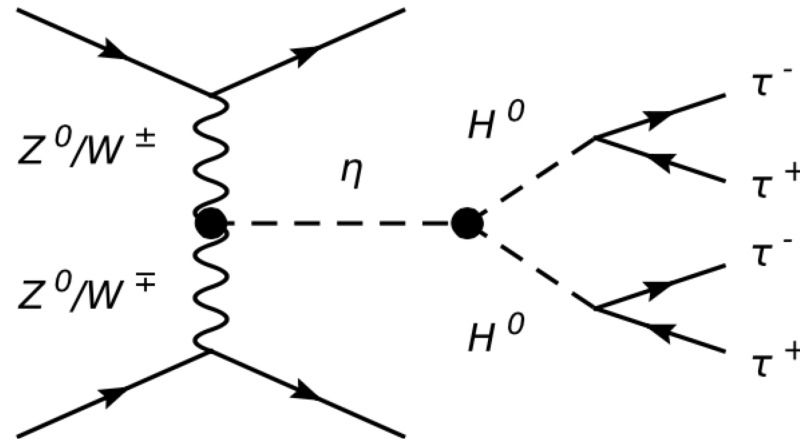
Why is the Higgs Boson So Light?

- Old idea: Higgs doublet (4 fields) is a Goldstone mode generated from the spontaneous breaking of a larger global symmetry
 - Higgs boson and W_L , Z_L are all Goldstone bosons from, eg. Spontaneously breaking global $SO(5) \rightarrow SO(4)$
 - Examples: Holographic Higgs, Little Higgs models...
 - Electroweak vev “ v ” is small compared to $SO(5)$ breaking scale “ f ”
- Vector boson scattering topology
 - Quarks emit longitudinal vector bosons which interact with new (presumably strong) dynamics
 - Quarks scatter by small angle in the forward direction

Longitudinal Vector Boson Scattering

Double Higgs Boson Production in the 4τ Channel from Resonances in Longitudinal Vector Boson Scattering at a 100 TeV Collider

AVK, [S. Chekanov](#), M. Low
Phys.Rev. D91 (2015) 114018



(a) The pseudo-rapidity distributions of the forward jets.

Forward Jet Coverage for Longitudinal VBS

$$V_L V_L \rightarrow \eta \rightarrow HH$$

AVK, S. Chekanov, M. Low

TABLE II. 5σ discovery mass reach for the $\eta \rightarrow HH \rightarrow 4\tau$ resonance, at a pp collider with $\sqrt{s} = 100$ TeV and $\mathcal{L} = 10 \text{ ab}^{-1}$, for various cuts values on minimum p_T of the forward jets. The fractional width of the η resonance is set to $\Gamma/M = 20\%$.

p_T^{\min} (GeV)	30	50	70	90	110
m_η (TeV)	3.53	2.90	2.35	1.92	1.56

- Lower p_T threshold on forward tagging jets is preferred
 - Reject pileup jets with good tracking in forward direction
 - Resolve overlapping pileup jets with higher granularity / spatial resolution (*a la* CMS high-granularity endcap calorimeter for HL-LHC)

Vector Boson Scattering

Double Higgs Boson Production in the 4τ Channel from Resonances in Longitudinal Vector Boson Scattering at a 100 TeV Collider

AVK, [S. Chekanov](#), M. Low

Phys.Rev. D91 (2015) 114018

TABLE III. 5σ discovery mass reach for the $\eta \rightarrow HH \rightarrow 4\tau$ resonance, at a pp collider with $\sqrt{s} = 100$ TeV and $\mathcal{L} = 10 \text{ ab}^{-1}$, for various cuts values on the maximum rapidity (y) of the forward jets. The fractional width of the η resonance is set to $\Gamma/M = 20\%$.

y^{max}	8	7	6	5	4
m_η (TeV)	2.9	2.9	2.81	2.42	1.75

Want jet rapidity coverage up to 6

Summary

Physics Conclusions

- Circular proton-proton colliders at very high energy provide unprecedented discovery potential
- New territory explored with precision measurements and direct searches is strongly motivated for
 - Solving the mysteries associated with the Higgs boson
 - Discovering WIMP Dark Matter
 - Understanding the electroweak phase transition and discovering the conditions for electroweak baryogenesis
- Potential for big surprises and discovery of unexpected new principles of nature

Detector Summary

- Entering new regime on detector design and technology
- Completion of the Standard Model and its consistency with all data implies
 - Energy scale of new physics is less well-defined now than when LHC was designed
 - We must prepare for a broader range of possible new physics
 - Specialized, targeted detectors risky as target signatures are unconstrained
 - Prudent to continue CDF & D0 (Run 2), ATLAS & CMS general-purpose detector philosophy
- Need improved capabilities
 - Better track momentum resolution
 - Maintain/improve b -tagging at high jet p_T and high track density
 - Improve hadronic τ -lepton identification efficiency \rightarrow high-granularity EMCAL
 - Boosted H/W/Z/top substructure \rightarrow high-granularity HCAL
 - Extend forward jet coverage to rapidity ~ 6 for vector boson scattering
 - Extend forward tracking for rejecting top quark background and suppressing forward pileup jets

More Challenges

- Readout bandwidth driven by high granularity
 - Wireless transmission ???
- Pileup of ~ 1000 additional interactions: handle with precision timing?
- Triggering
 - challenging to trigger on disappearing tracks and long-lived particles

Signatures of displaced decays

5

- ☐ Inner Tracker green
- ☐ EM Calorimeter Blue/green
- ☐ Hadronic calorimeter Blue
- ☐ Muon system Grey

Displaced decay signatures

1. Decay in muon system - jet
2. Two body decay (lepton jet)
3. Decay in HCAL of - jet
4. Emerging jets
5. Inner Tracker decay to jets
6. Decay to jets in the IT
7. Disappearing (invisible) LLP
8. Non-pointing $\gamma \rightarrow e^+e^-$

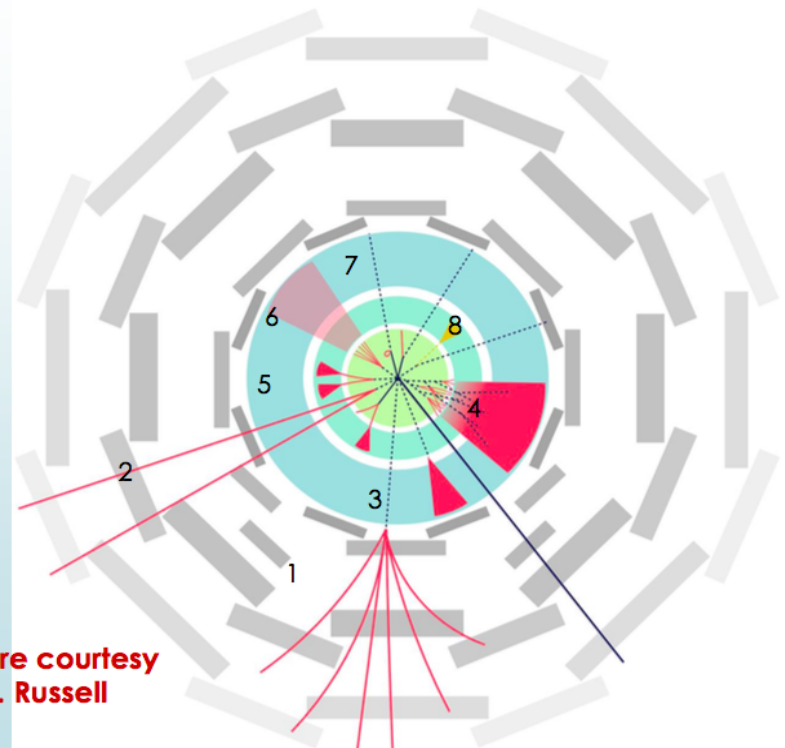


Figure courtesy
of H. Russell