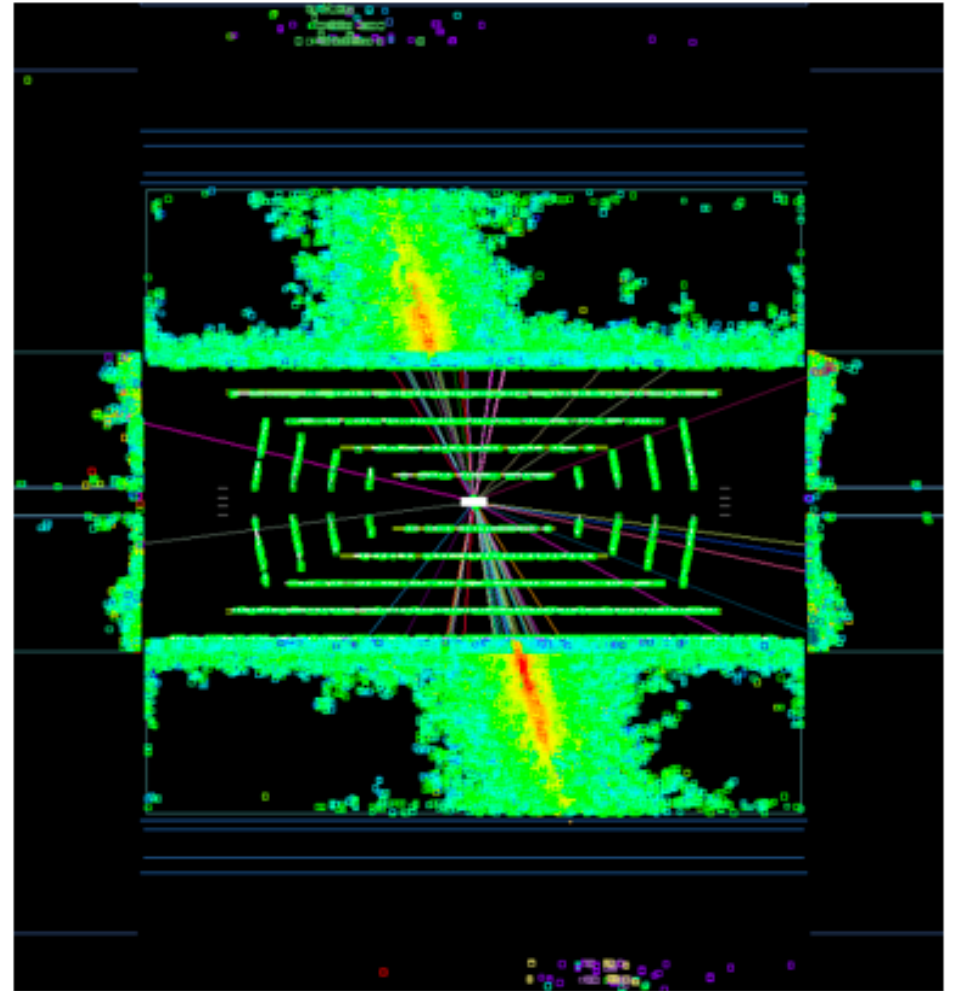
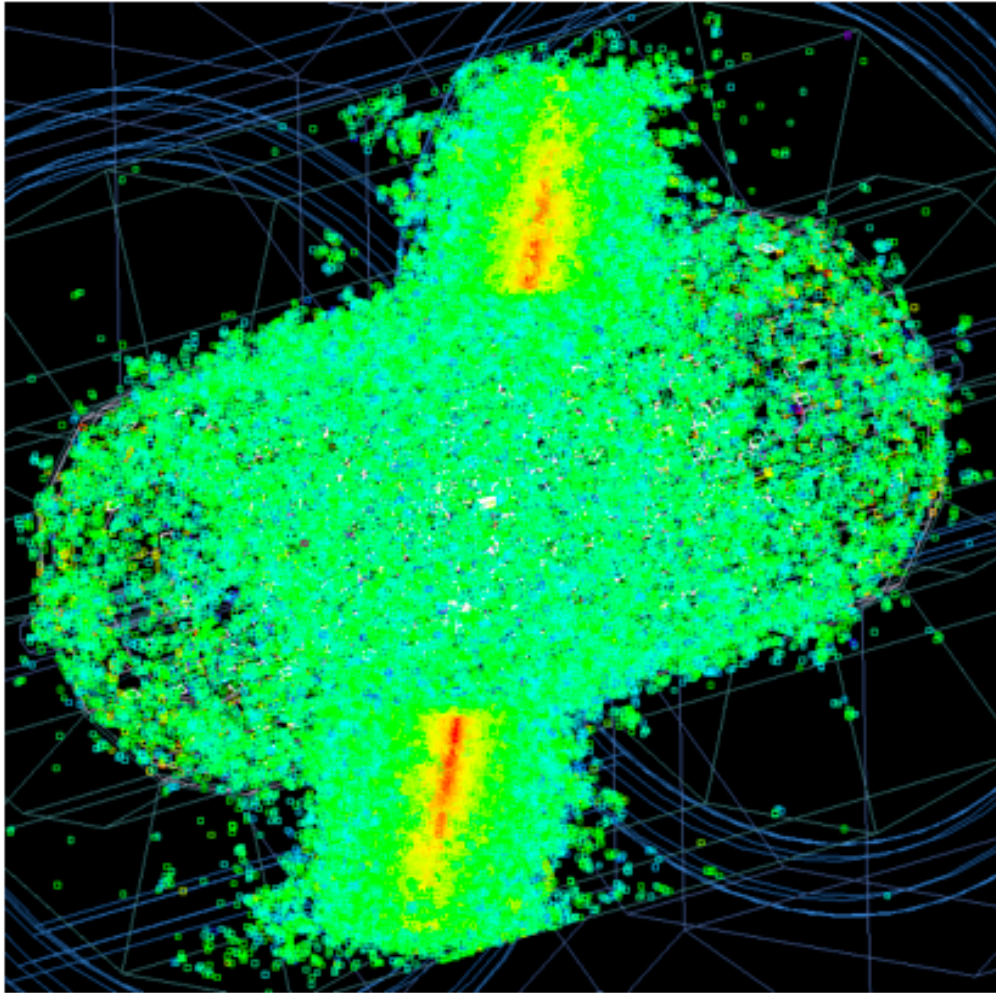


# Physics and Experiments at Future $pp$ Colliders

Ashutosh Kotwal  
Duke University



Workshop on Neutrinos at the High Energy Frontier  
University of Massachusetts, Amherst  
July 18, 2017

# Circular $pp$ Collider Physics Goals

- Testable reasons why the Standard Model must be incomplete
  - Dark Matter could be
    - Weakly-interacting particles
    - Particles interacting through Higgs portal
    - Interacting with SM particles through gravity
  - Electroweak Baryogenesis
    - Can the electroweak phase transition (formation of Higgs VEV) provide the out-of-equilibrium condition needed for matter-antimatter asymmetry observed?
  - Can the parameter space of new physics be a bounded parameter space?
    - Can it be fully covered with a 100-TeV scale  $pp$  collider?
- Naturalness – the need to explain the lightness of the Higgs mass – testing Naturalness at  $10^{-4}$



# Guidance for Detector Design

- As long as Standard Model continues to work, “higher energy is better”
- Naturalness arguments push towards higher masses
- Dark Matter, Electroweak Baryogenesis *may* relate to physics at lower masses and smaller couplings
- Other reasons that new physics may hide at low mass with weak couplings
  - “Neutral Naturalness” (partners without QCD color charge)
  - e.g. twin Higgs, Hidden Sector
  - Higgs portal to new sector (SM interactions via Higgs only)
- Implications for detector design: larger dynamic range of  $p_T$  of objects
  - Starting at  $\sim 20$  GeV leptons, photons and  $b$ -quarks (same as LHC, e.g.  $gg \rightarrow HH$ )
  - Going up to  $\sim 7$  times the highest  $p_T$  probed at LHC
- Also large rapidity range for all objects due to higher longitudinal boost

# Executive Summary

- Entering new regime on all fronts
  - Accelerator physics and design
  - Detector technology and design
- Completion of the Standard Model and its consistency with all data implies
  - Energy scale of new physics is less well-defined now than when LHC was designed
  - We must prepare for a broader range of possible new physics
- Detectors will need to be more capable on all fronts
  - Faster
  - Much higher resolution
  - Much higher granularity
  - Much more forward-detection capability
  - Much higher bandwidth, smarter triggers
- HL-LHC upgrade will provide much experience and insights

# All-Purpose Detector Goals in a Nutshell

- Maximize  $A \times \epsilon$ : all detectable particles
  - should be detected and over as much of the angular phase space as possible
  - And be well-measured over as much of their energy spectrum as possible (or of most importance to the interesting signals)
- Leptons of interest: electrons, muons and  $\tau$ -leptons
- Photons
- Quarks and gluons hadronize to jets of particles
- $b$ -quarks are special and need to be distinguished from other jets
- Undetectable particles like neutrinos and Dark Matter can only have their transverse momentum sum inferred
  - Catch all visible momentum
  - Impose transverse momentum conservation
  - Hermeticity is important

# All-Purpose Detector Goals in a Nutshell (2)

- Minimize B: reducible backgrounds from mis-identified particles
  - High rate of fragmentation pions, kaons, and photons misidentified as prompt electrons, photons and muons
  - Generic jets mis-identified as  $b$ -quark jets
  - Electrons and generic jets mis-identified as  $\tau$  leptons
  - Energy resolution of detected particles, or missed visible energy due to missing instrumentation, leads to fake missing  $p_T$  signature
  - Hermetic detectors have become very important
- Maximize  $\Delta t \times L$ : enable data-taking in high instantaneous luminosity environment
  - Large number of particles from additional (uninteresting) pp collisions
    - Can confuse/obfuscate the particles from the interesting collision
  - Total exposure of sensors to radiation flux scales with integrated luminosity and falls off with distance from collision point
    - Radiation damage causing degradation of sensor efficiency and increasing noise

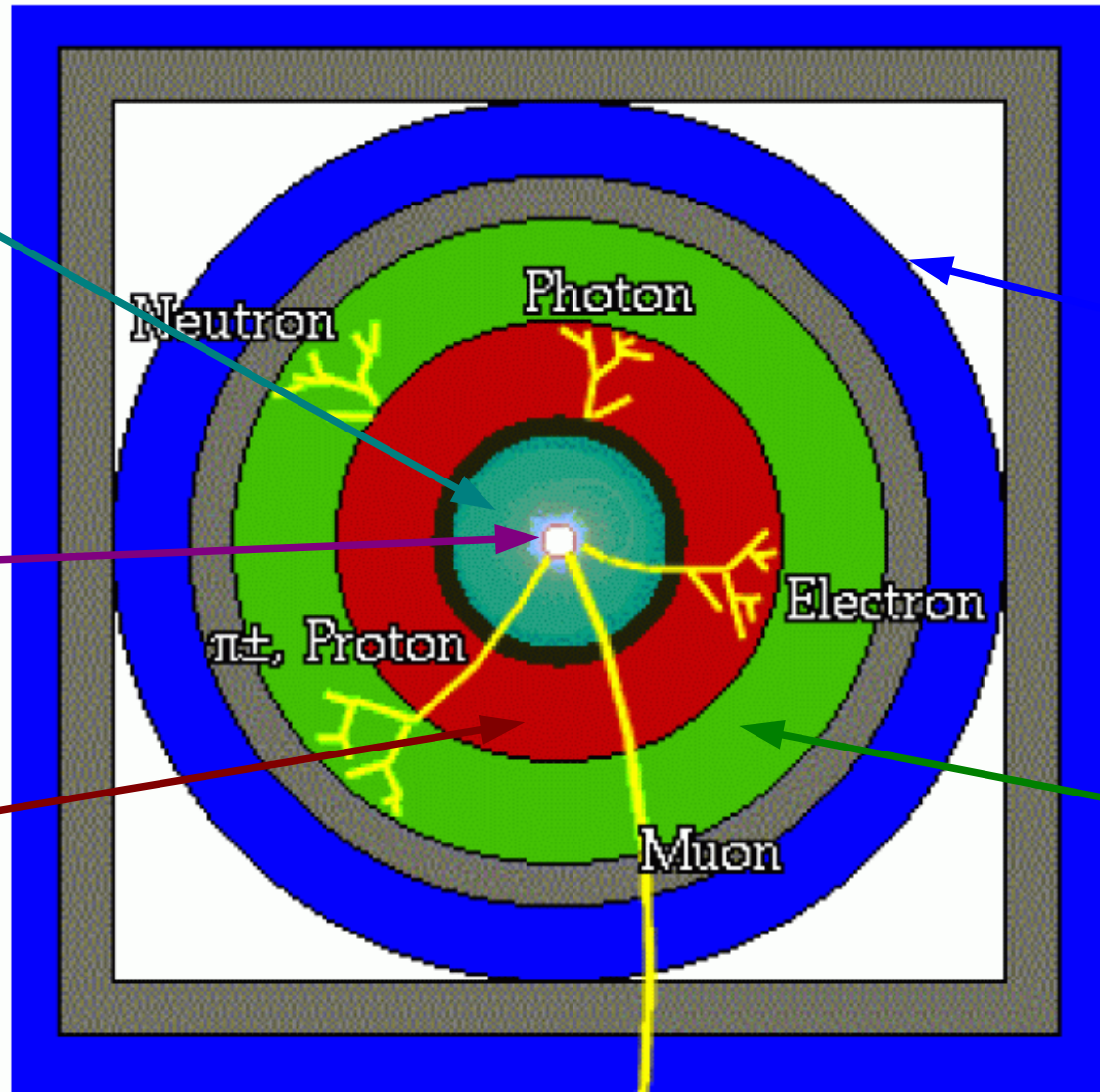


# Particle Detection

Drift chamber:  
reconstruct particle  
trajectory by sensing  
ionization in gas  
on high voltage wires

Silicon detector:  
reconstruct particle  
trajectory by sensing  
ionization in planar  
silicon sensors  
(diodes)

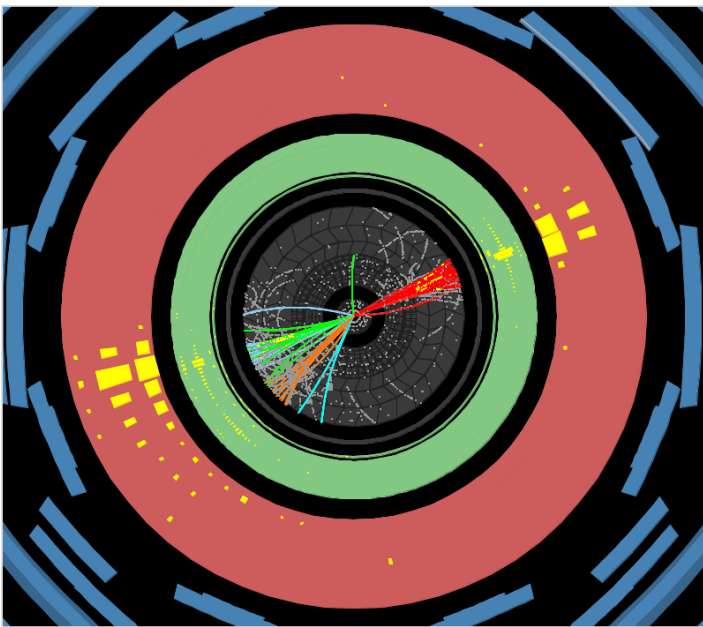
Electromagnetic  
(EM) calorimeter:  
metal sheets cause  
 $e/\gamma$  shower, sense  
light or charge



Muon chambers:  
detect penetrating  
particles behind  
shielding

Hadronic  
calorimeter:  
metal sheets  
cause hadronic  
showers, sense  
scintillator light  
or charge

# Particle Detection

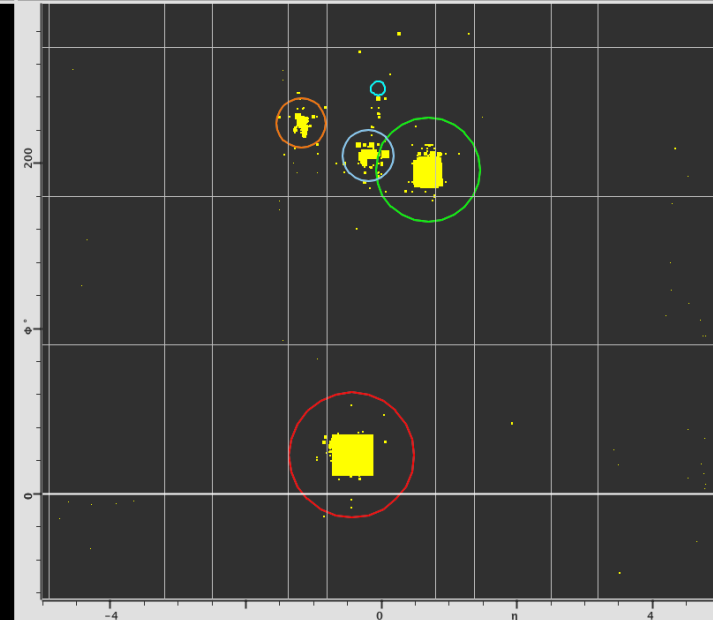
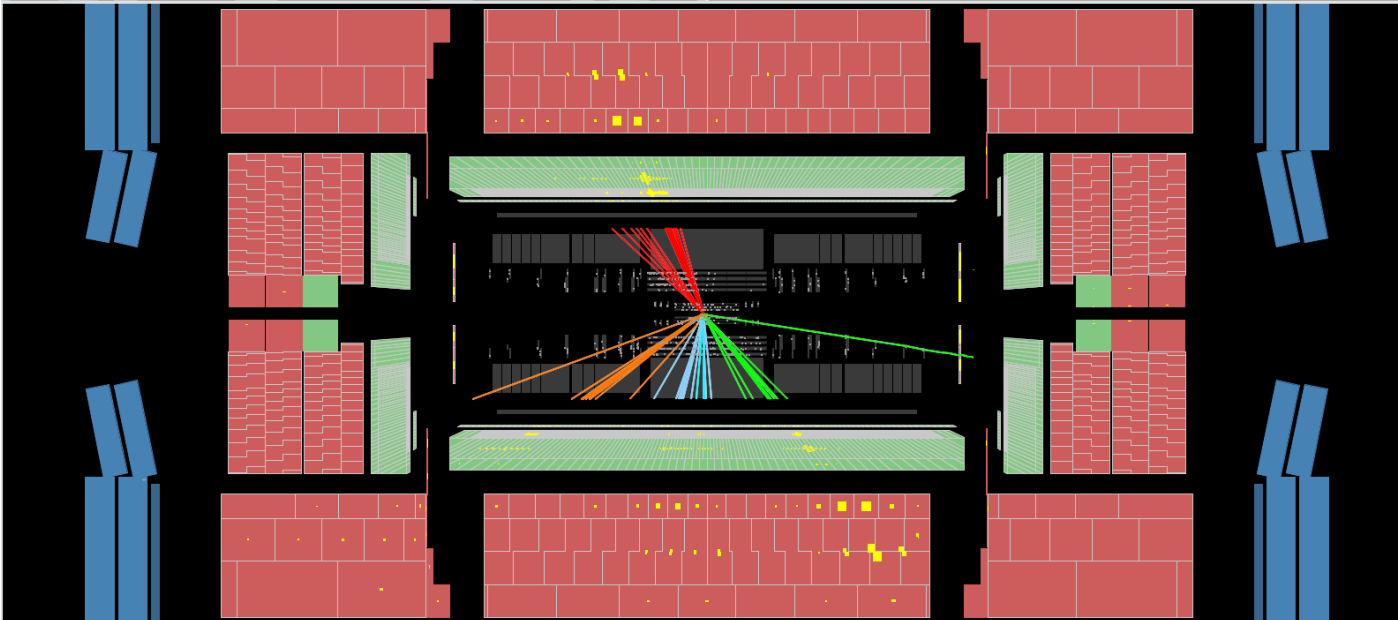
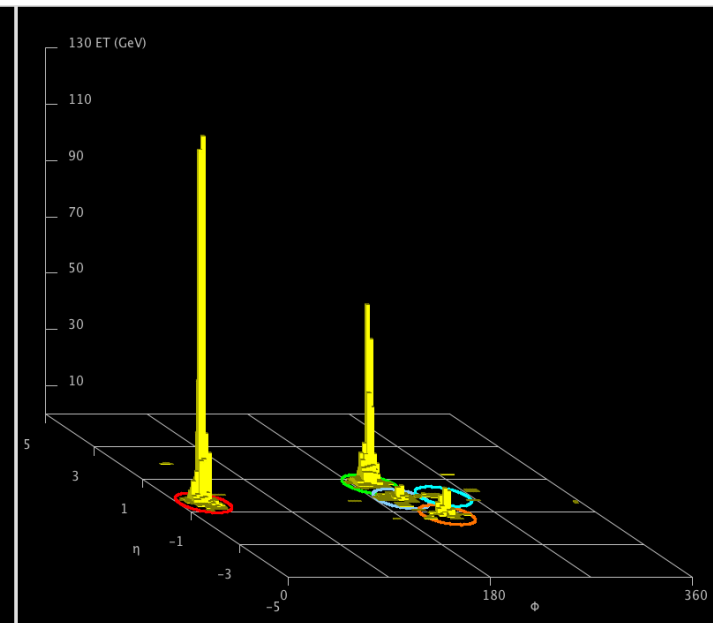


# ATLAS EXPERIMENT

Run Number: 158548, Event Number: 2486978

Date: 2010-07-04 06:46:45 CEST

## Multijet Event in 7 TeV Collisions



# Rate comparisons at 8, 14, 100 TeV

	$N_{100}$	$N_{100}/N_8$	$N_{100}/N_{14}$
<b>gg→H</b>	16 G	$4.2 \times 10^4$	110
<b>VBF</b>	1.6 G	$5.1 \times 10^4$	120
<b>WH</b>	320 M	$2.3 \times 10^4$	66
<b>ZH</b>	220 M	$2.8 \times 10^4$	84
<b>ttH</b>	760 M	$29 \times 10^4$	420
<b>gg→HH</b>	28 M		280

$$N_{100} = \sigma_{100\text{TeV}} \times 20 \text{ ab}^{-1}$$

$$N_8 = \sigma_{8\text{TeV}} \times 20 \text{ fb}^{-1}$$

$$N_{14} = \sigma_{14\text{TeV}} \times 3 \text{ ab}^{-1}$$

## Statistical precision:

- O(100 - 500) better w.r.t Run 1
- O(10 - 20) better w.r.t HL-LHC

# Magnetic Tracking



# Tracker Design – the heart of the experiment

Momentum is determined by measurement of **track curvature**  $\kappa = 1/\rho$  in B field:

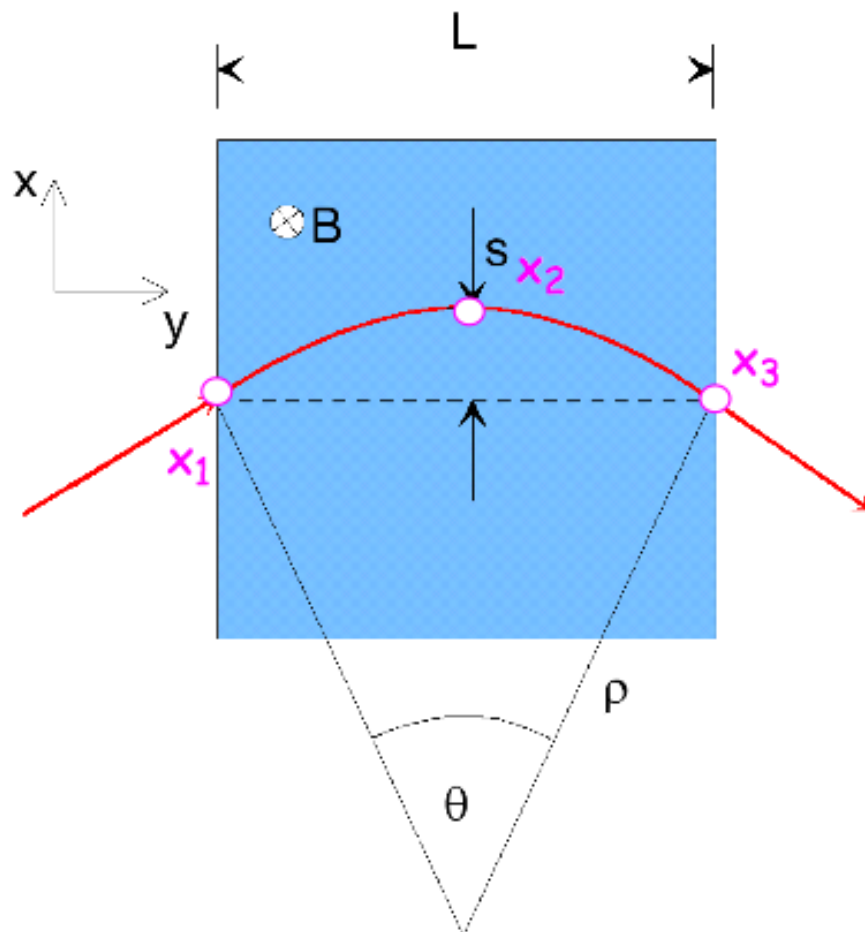
Measure **sagitta**  $s$  of the track. For the momentum component transverse to B field:

$$p_T = qB\rho$$

Units:  $p_T[\text{GeV}] = 0.3B[\text{T}]\rho[\text{m}]$

$$\frac{L/2}{\rho} = \sin\frac{\theta}{2} \approx \frac{\theta}{2} \text{ (for small } \theta) \Rightarrow \theta \approx \frac{L}{\rho} = \frac{0.3B \cdot L}{p_T}$$

$$s = \rho\left(1 - \cos\frac{\theta}{2}\right) \approx \rho\left(1 - \left(1 - \frac{1}{2}\frac{\theta^2}{4}\right)\right) = \rho\frac{\theta^2}{8} \approx \frac{0.3L^2B}{8 p_T}$$



# Relative Momentum Error

For 3 points the relative momentum resolution is given by:  $\frac{\sigma(p_T)}{p_T} = \frac{\sigma_s}{s} = \sqrt{\frac{3}{2}} \sigma_x \cdot \frac{8p_T}{0.3BL^2}$

- degrades **linearly** with **transverse momentum**
- improves **linearly** with increasing **B field**
- improves **quadratically** with **radial extension** of detector

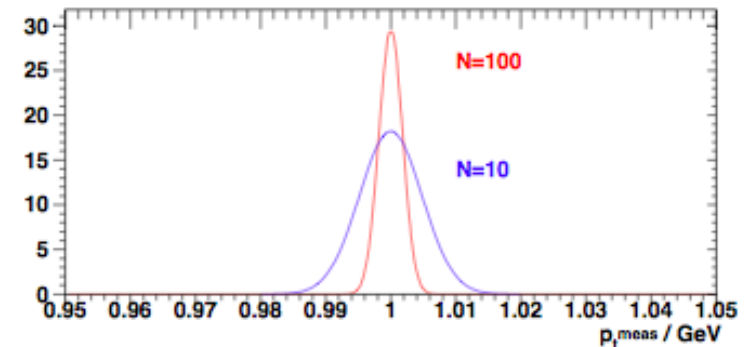
In the case of  $N$  equidistant measurements according to **Gluckstern** [NIM 24 (1963) 381]:

$$\frac{\sigma(p_T)}{p_T} = \frac{\sigma(\kappa)}{\kappa} = \frac{\sigma_x \cdot p_T}{0.3BL^2} \sqrt{\frac{720}{(N+4)}} \quad (\text{for } N \geq 10, \text{ curvature } \kappa = 1/\rho)$$

Example: For  $p_T = 1\text{GeV}$ ,  $L = 1\text{m}$ ,  $B = 1\text{T}$ ,  $\sigma_x = 200\mu\text{m}$  and  $N = 10$  one obtains:

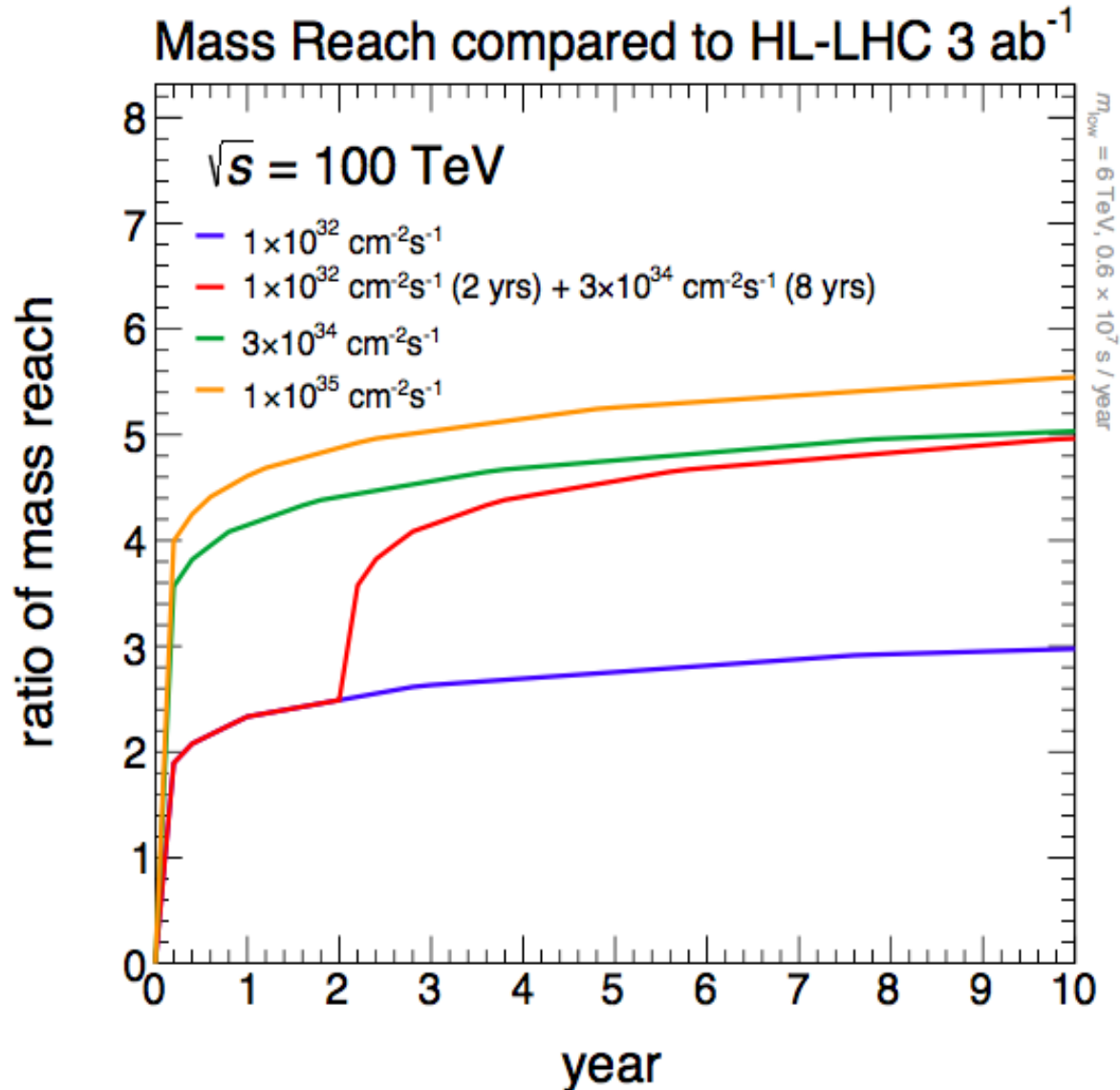
$$\frac{\sigma(p_T)}{p_T} \approx 0.5\% \quad \text{for a sagitta } s \approx 3.8\text{cm}$$

Important track detector parameter:  $\frac{\sigma(p_T)}{p_T^2}$  (%/GeV)



# Collider Luminosity and Energy

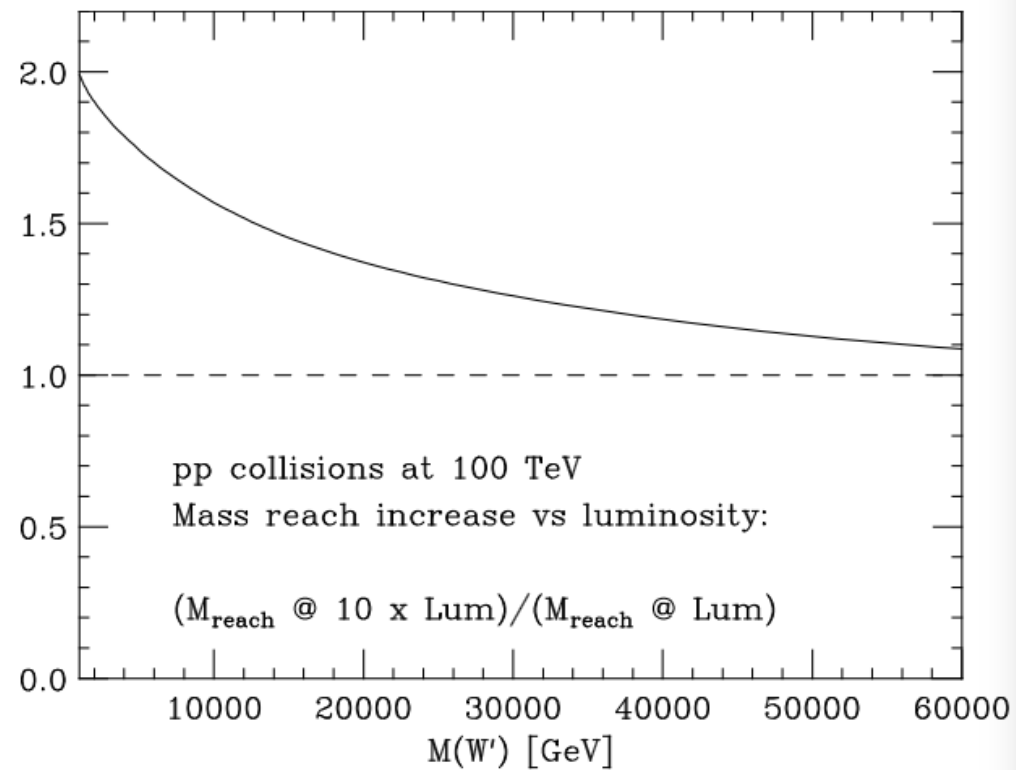
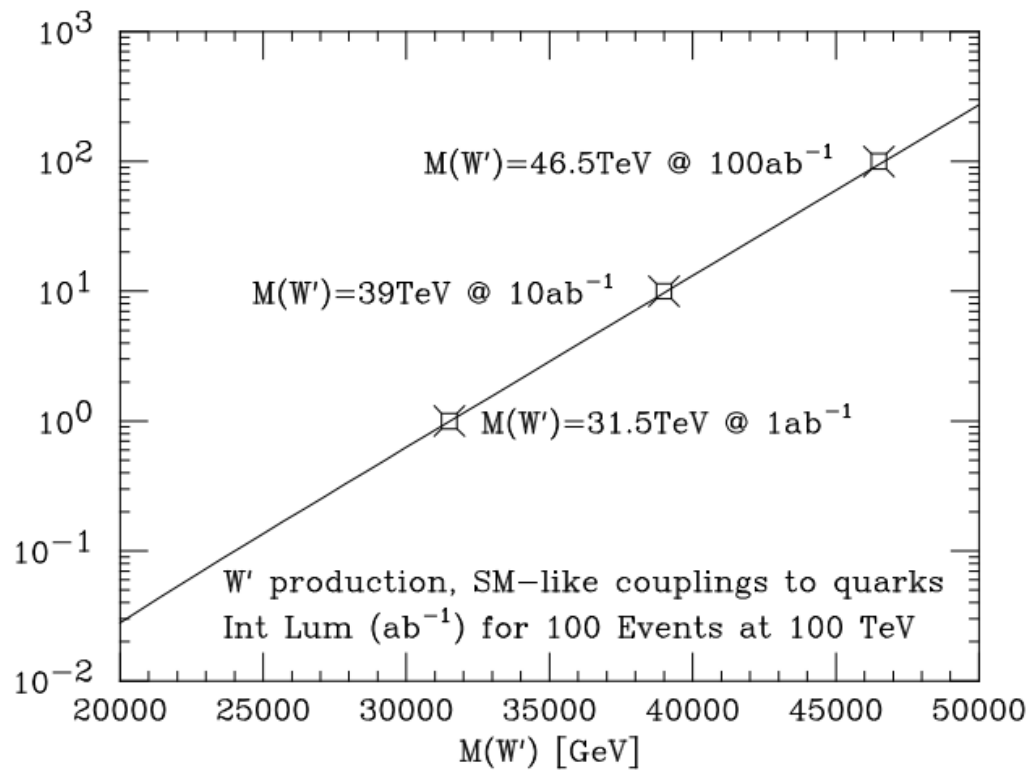
- Collider luminosity evolution for high-mass reach



(from L-T. Wang)

# Highest Mass Leptonic Resonances

- HL-LHC studies showed  $Z' \rightarrow ll$  reach up to 6.5 TeV
- **Scaling to 100 TeV collider  $\Rightarrow$  45 TeV with  $150 \text{ ab}^{-1}$  or 38 TeV with  $15 \text{ ab}^{-1}$** 
  - 7 TeV change in mass reach for factor of 10 change in luminosity



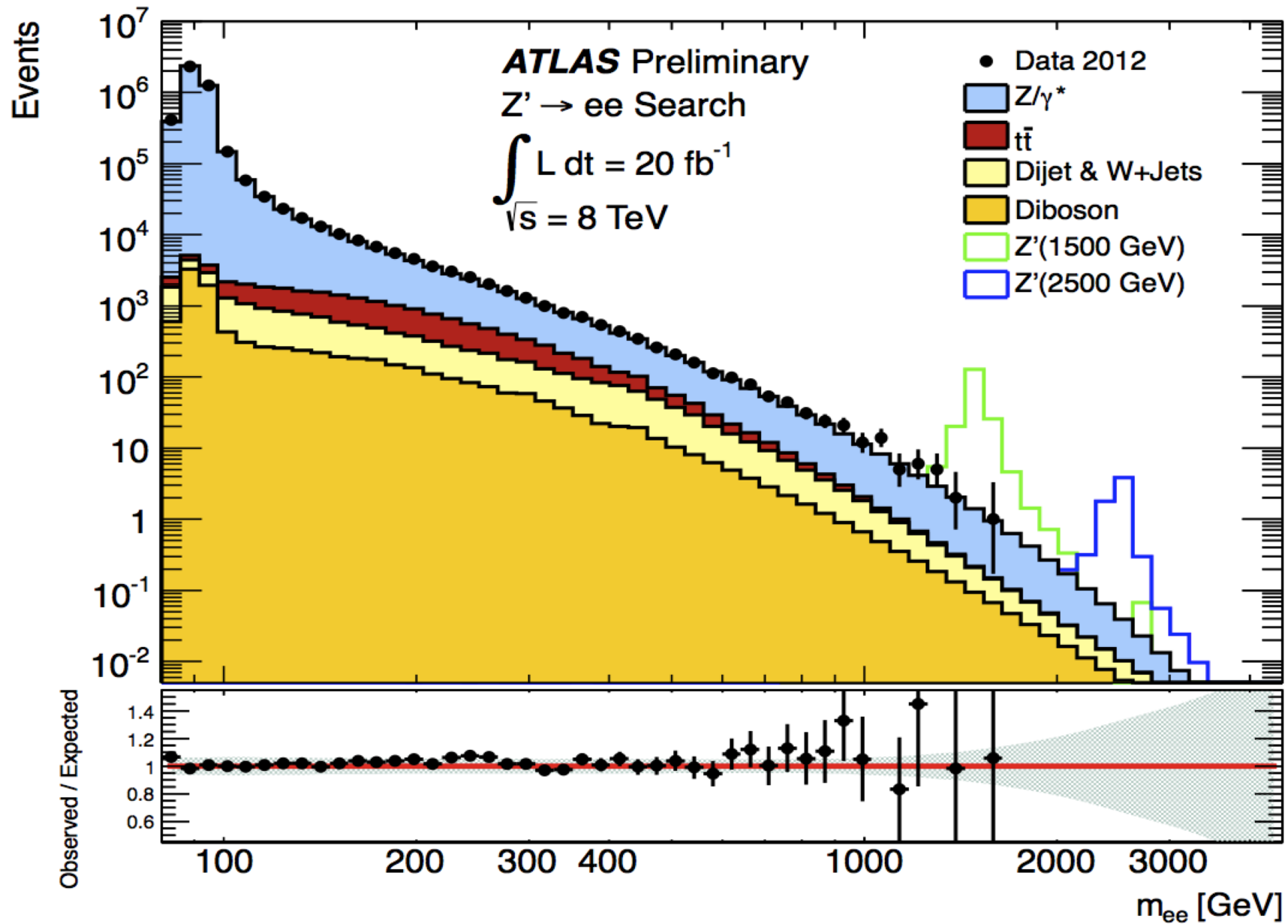
$\Rightarrow$  producing 20 TeV leptons

(from M. Mangano)



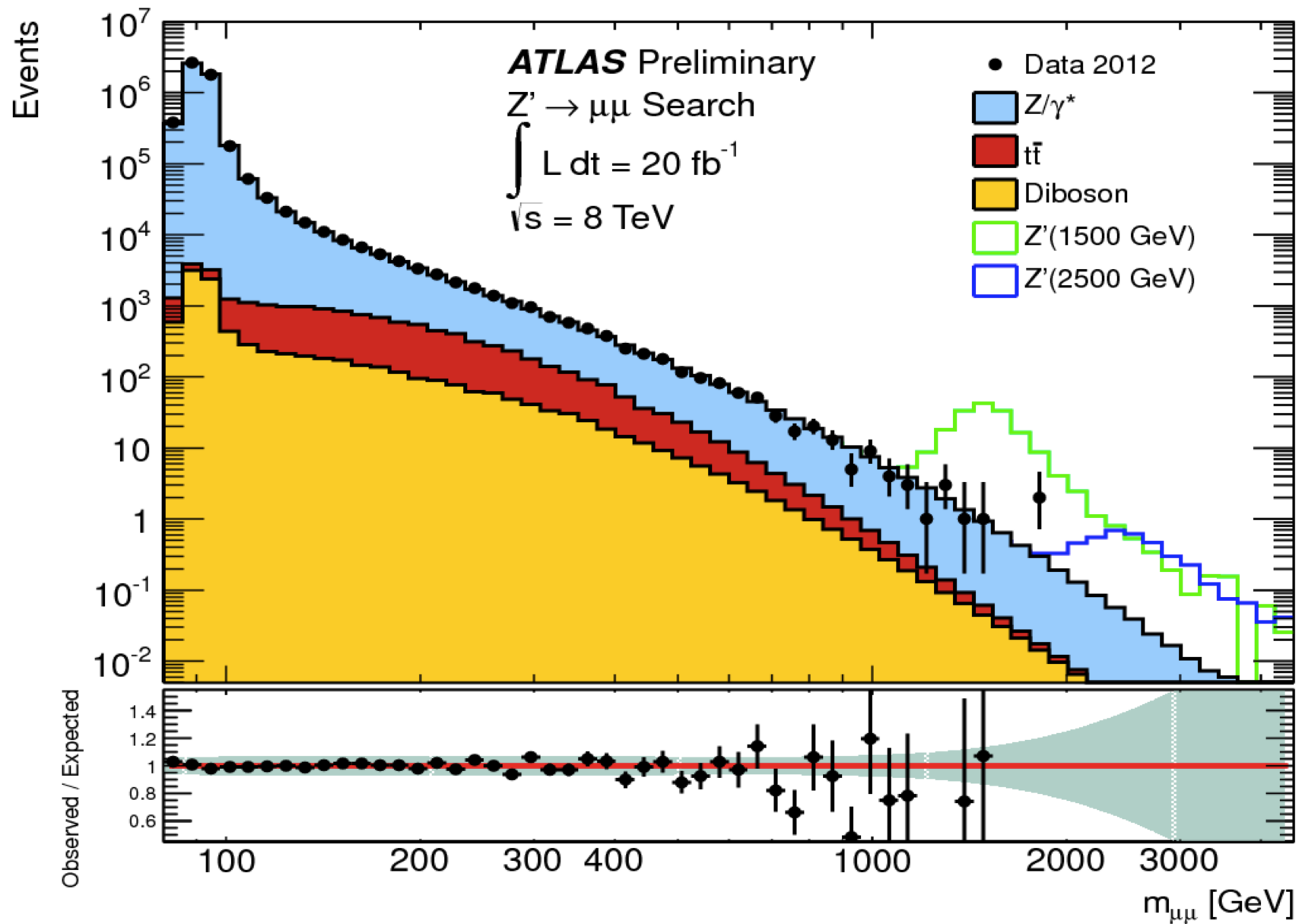
# Dielectron Mass Spectrum

Multi-TeV masses probed at LHC

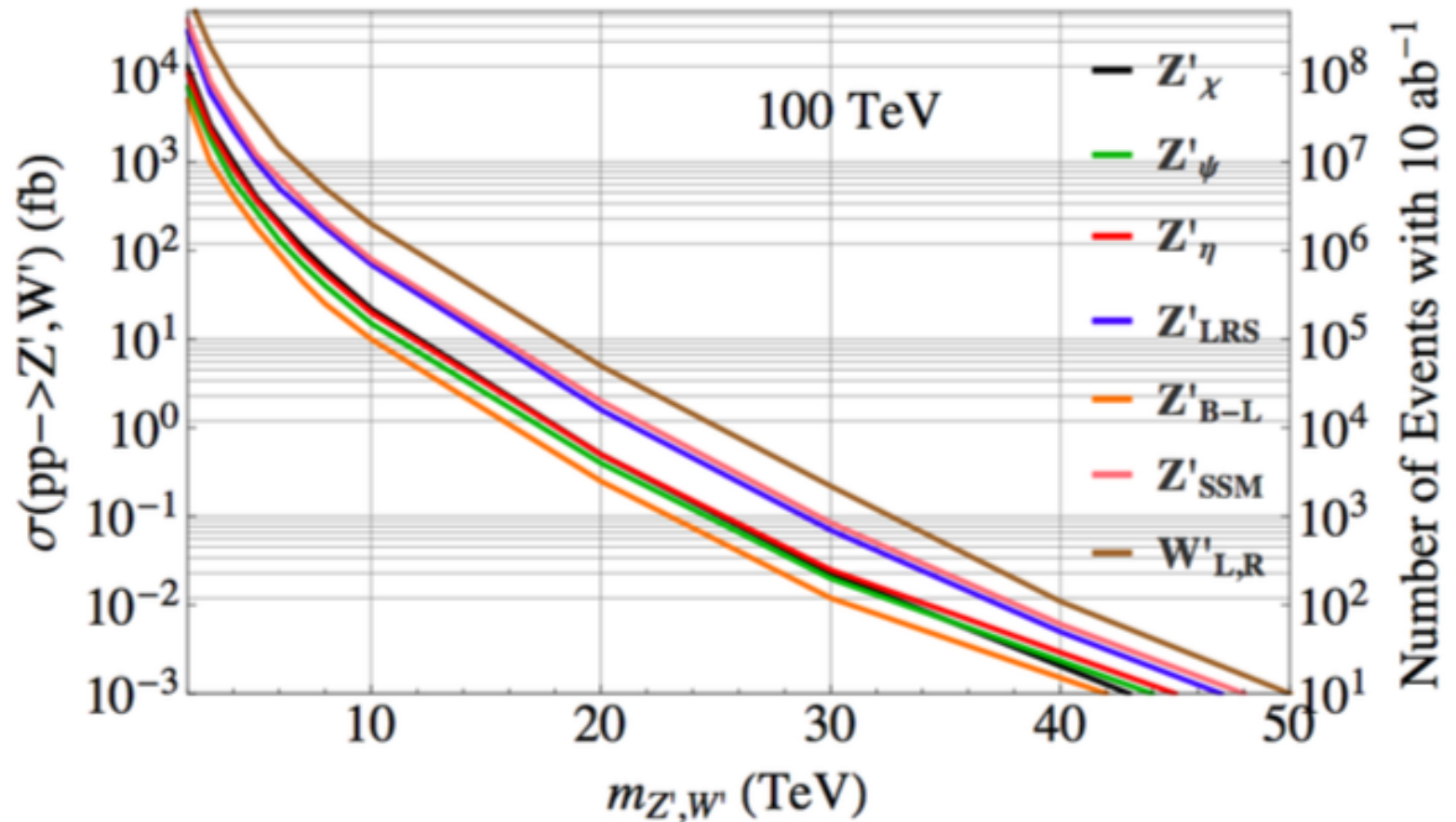


# Dimuon Mass Spectrum

Multi-TeV masses probed at LHC



# Exploring New Territory - New Weak Gauge Interactions



**Discovery reach**

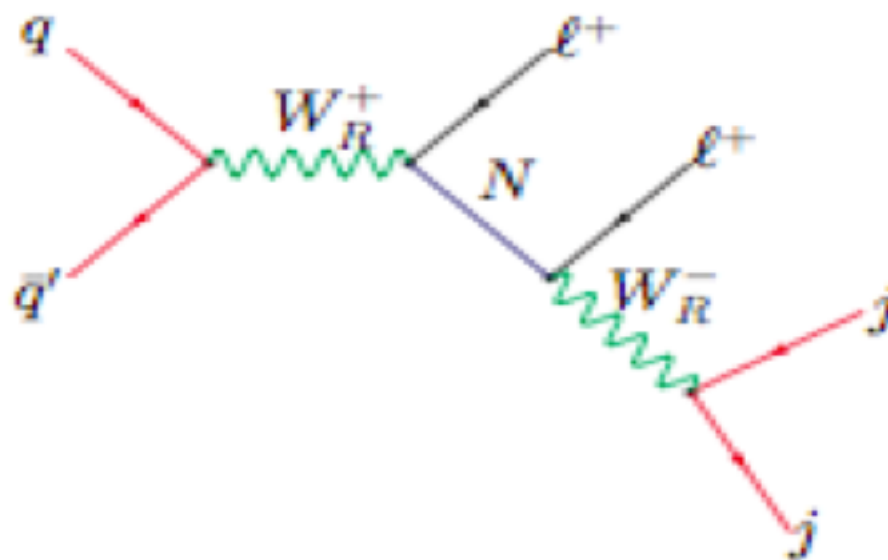
T.Rizzo, arXiv:1403.5465

Model	$1 \text{ ab}^{-1}$	$10 \text{ ab}^{-1}$	$100 \text{ ab}^{-1}$
SSM	23.8	33.3	41.3
LRM	22.6	31.5	39.5
$\psi$	20.1	29.1	37.2
$\chi$	22.7	30.6	38.2
$\eta$	20.3	29.8	38.0
I	22.4	29.2	36.2

10-fold increase in luminosity  
 $\rightarrow \sim 7 \text{ TeV}$  increase in mass reach

# Demands on $p_T$ Resolution

- High-mass dimuon resonances most demanding on tracker momentum resolution
- If universal coupling to leptons, dielectron channel is reliable
- Non-universal couplings plausible:
  - Higgs mechanism: additional Higgs bosons with  $H \rightarrow \mu\mu$
  - Left-right seesaw model of neutrino masses



$$N \rightarrow l^\pm jj$$

(Keung, Senjanovic'83)

- Prudent to maintain muon  $p_T$  resolution (%) from LHC to 7x higher  $p_T$



# Maintaining Fractional $p_T$ Resolution

- Resolution gain with number of hits on track is slow (improves as  $\sqrt{N}$ )
- Resolution improves linearly with  $BL^2 \sim$  stored magnetic field energy in tracker
- Resolution improves linearly with hit resolution

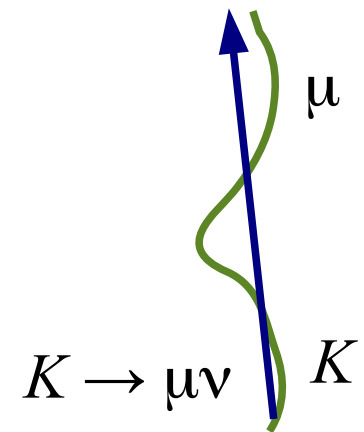
Three tracker/magnet geometries being considered:

- see Dr. Marcello Mannelli's talk at Fermilab's “Next Steps in the Energy Frontier – Hadron Collider” Workshop

<https://indico.fnal.gov/conferenceOtherViews.py?view=standard&confId=7864>

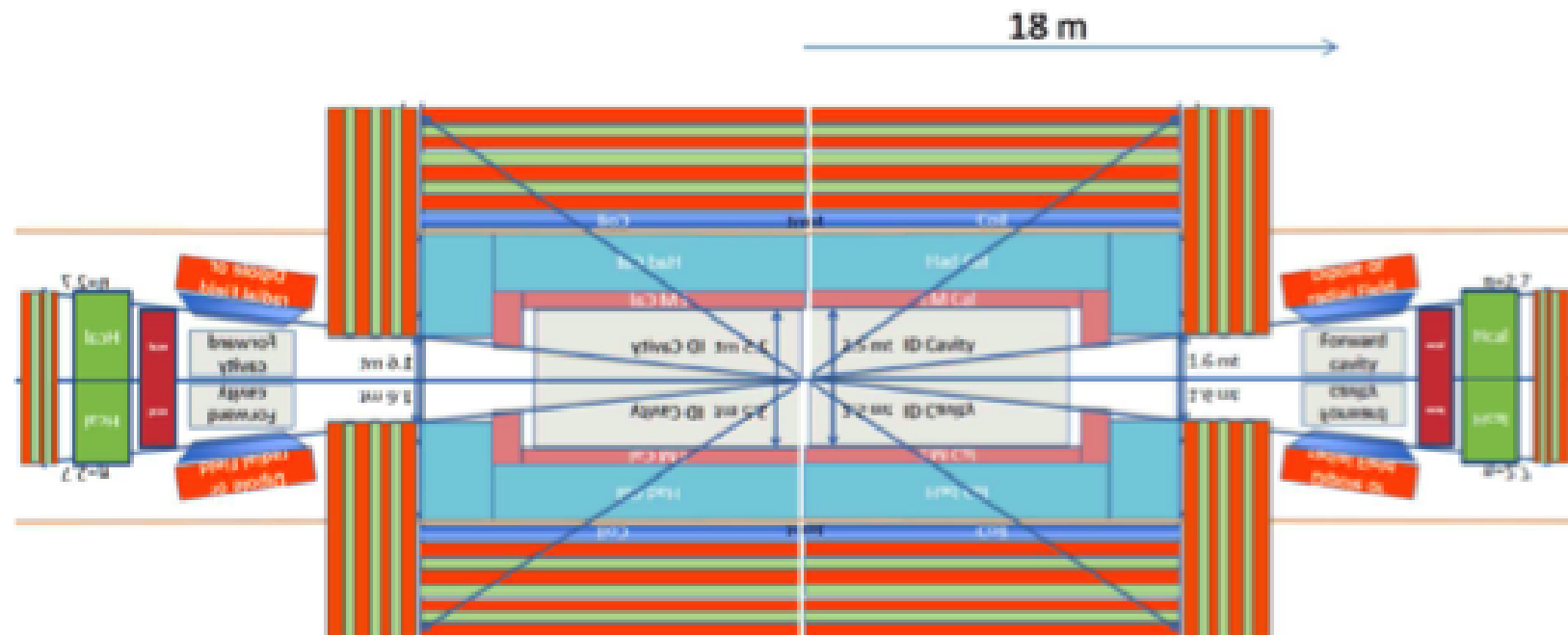
Stored energy in the tracker magnetic field in the 50-100 GJ range (similar to ITER)

Need to measure muon momentum after shielding, to eliminate mis-measured decays-in-flight with very high reconstructed  $p_T$





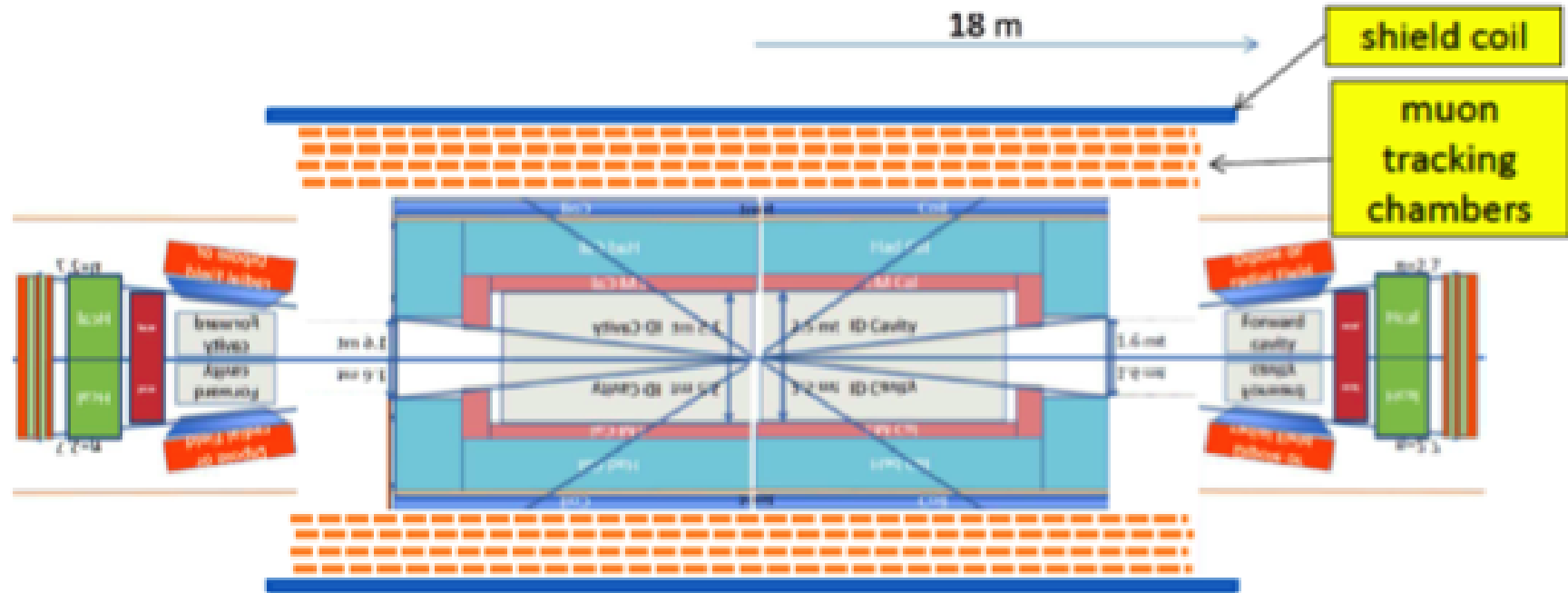
## 2. Option 1: Solenoid-Yoke + Dipoles (CMS inspired)



- ❖ **Solenoid:** 10-12 m diameter, 5-6 T, 23 m long  
+ massive Iron yoke for flux shielding and muon tagging.
- ❖ **Dipoles:** 10 Tm with return yoke placed at  $z \approx 18$  m.  
Practically no coupling between dipoles and solenoid.  
They can be designed independently at first.



## 2. Option 2: Twin Solenoid + Dipoles



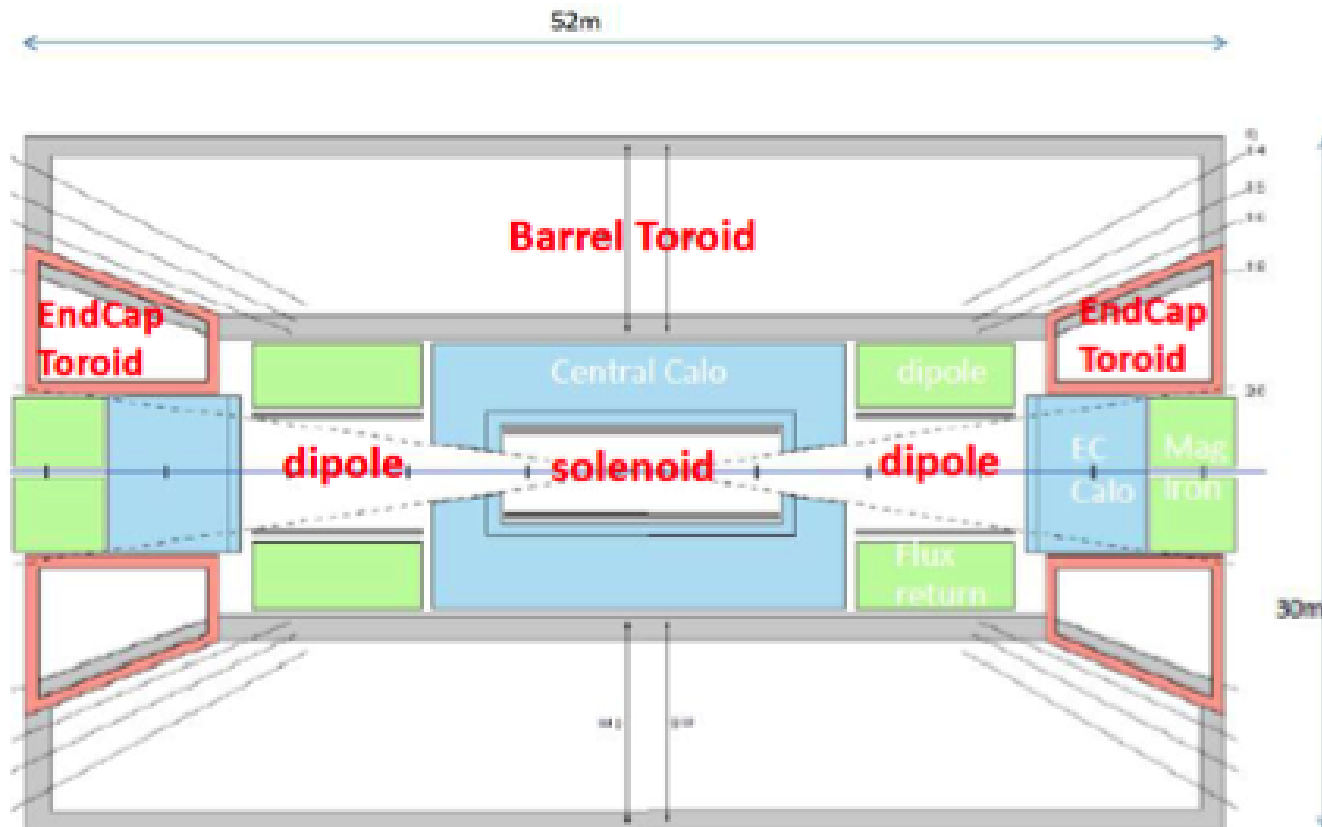
**Twin Solenoid:** a 6 T, 12 m dia x 23 m long main solenoid + an active shielding coil

**Important advantages:**

- ✓ **Nice Muon tracking space:** area with 2 to 3 T for muon tracking in 4 layers.
- ✓ **Very light:** 2 coils + structures,  $\approx 5$  kt, only  $\approx 4\%$  of the option with iron yoke!
- ✓ **Much smaller:** system outer diameter is significantly less than with iron .

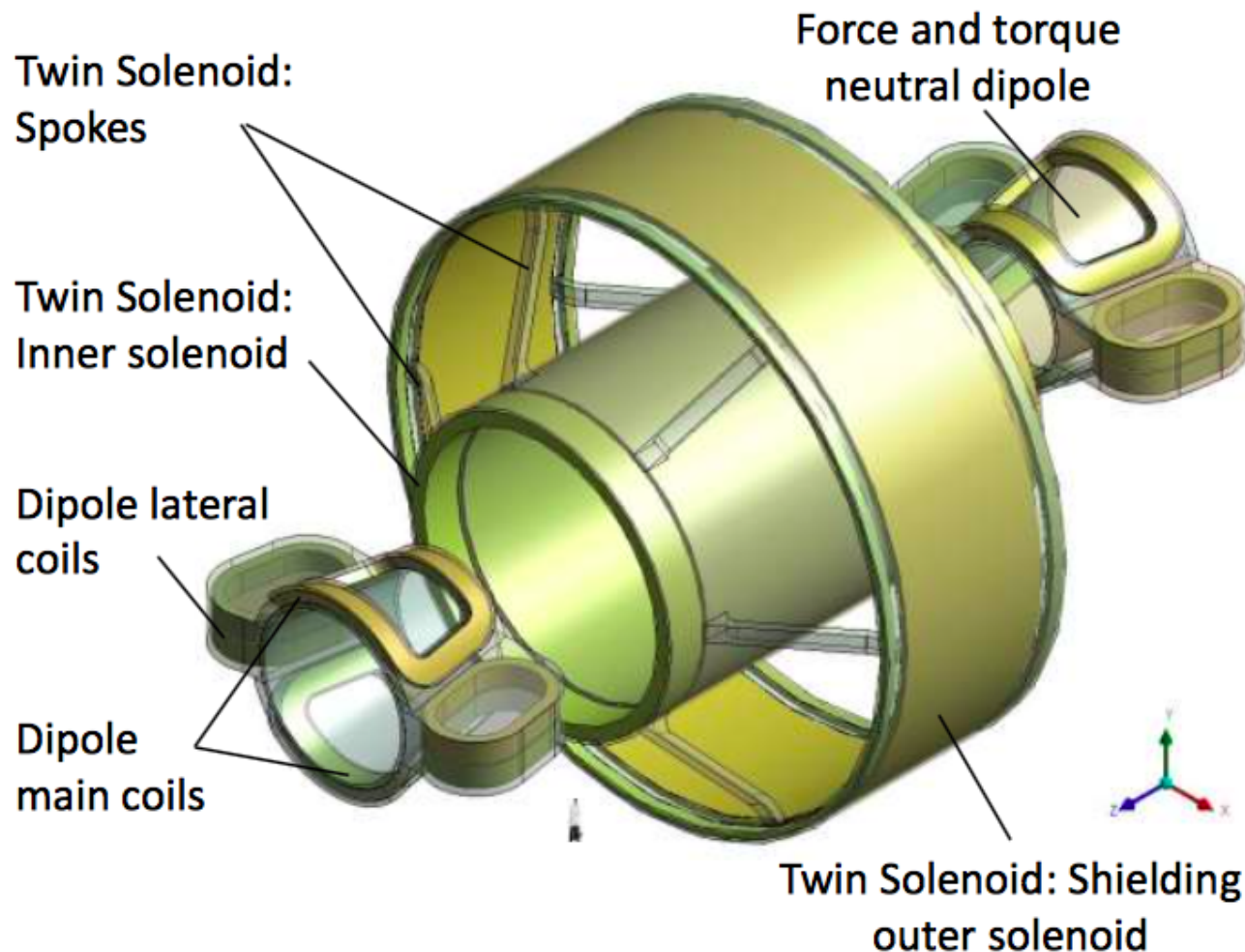


## 2. Option 3: Toroids + Solenoid + Dipoles (ATLAS +)



- ❖ 1 Air core Barrel Toroid with 7 x muon bending power  $B_z L^2$ .
- ❖ 2 End Cap Toroids to cover medium angle forward direction.
- ❖ 2 Dipoles to cover low-angle forward direction.
- ❖ Overall dimensions: 30 m diameter x 51 m length (36,000 m<sup>3</sup>).

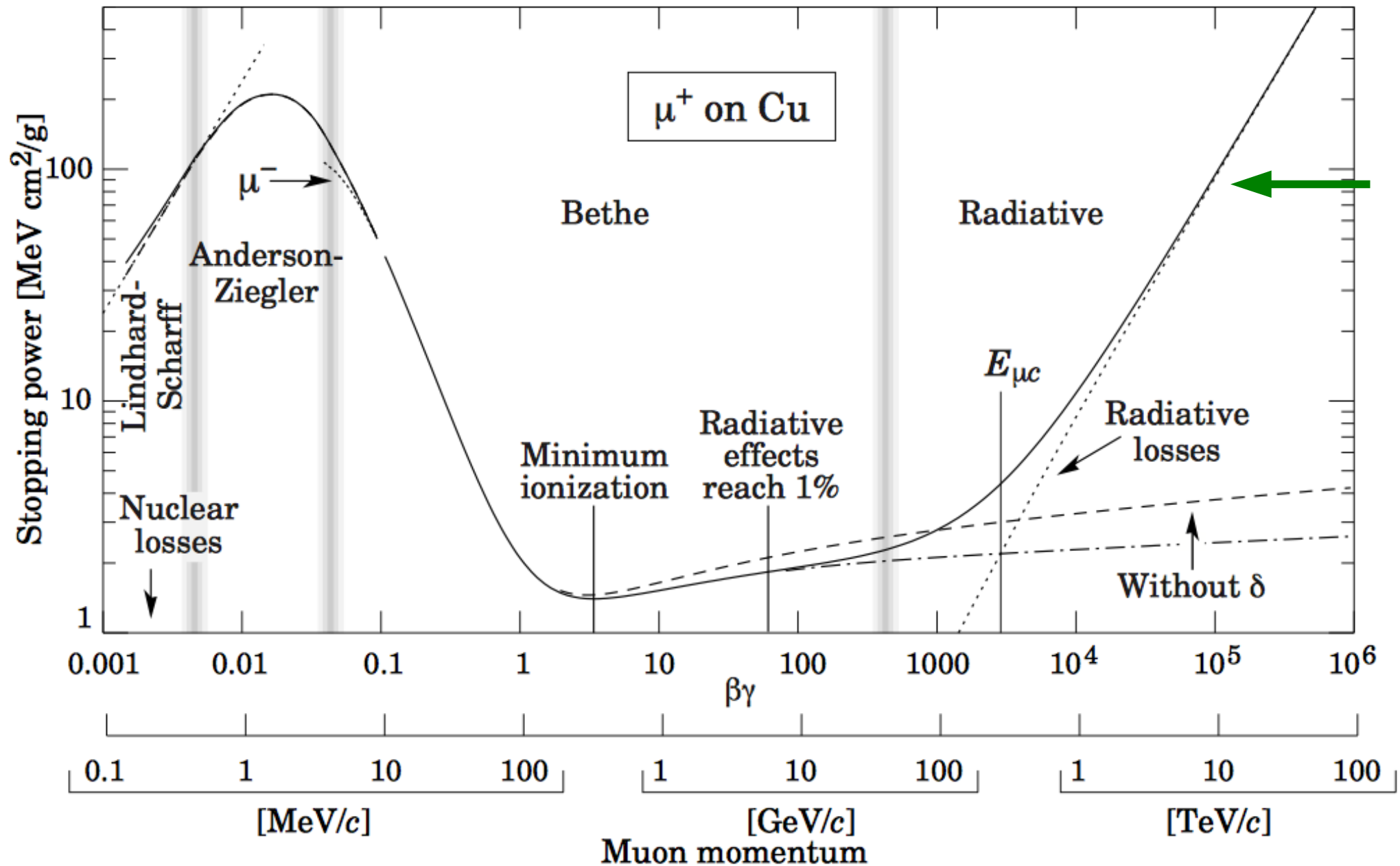
# Twin Solenoid & Dipole system – bare coils



Property	Value
TS cold mass	3.2 kt
TS vacuum vessel mass	2.4 kt
TS stored energy	53 GJ
Dipoles cold mass	2x 380 t
Dipoles vac. vessel mass	To be det.
Dipoles stored energy	2x 1.5 GJ
Free bore	12 m
Outer diameter	27 m
System length	42 m
<b>Total stored energy</b>	<b>56 GJ</b>

(from Herman ten Kate)

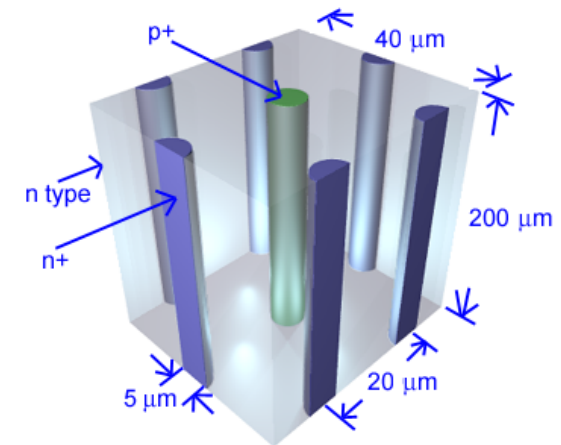
# High Energy Muon Bremsstrahlung



- For a  $\sim 10 \text{ TeV}$  muon, average energy loss  $\sim 1 \text{ GeV} / \text{cm} \sim 16 \text{ GeV} / \text{interaction length} \sim 200 \text{ GeV}$  in hadronic calorimeter, with long tailed distribution

# Improving Hit Resolution

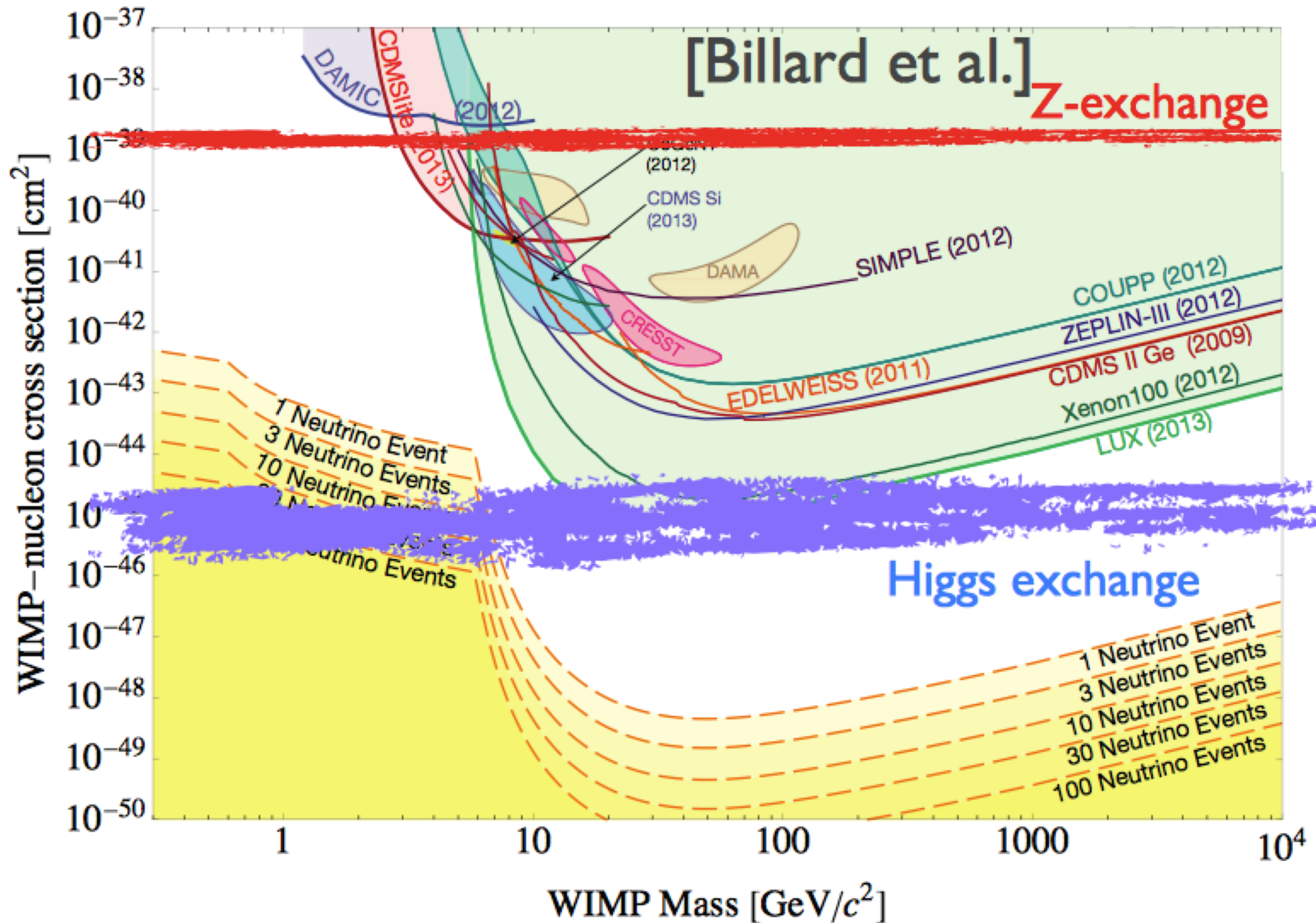
- Smaller pixels with silicon sensors have multiple advantages
  - Improved hit resolution linearly improves momentum resolution at high  $p_T$
  - Higher granularity improves two-track resolving power
    - Helps resolve close-by tracks and maintain track reconstruction efficiency in
      - high-density environment (inside boosted jets)
      - High-occupancy environment (pileup at high L)
- Issues:
  - Higher readout rate required
  - Power may be dominated by inter-pixel capacitance, which does not reduce with pixel size
    - More pixels  $\Rightarrow$  more power
- Potential solutions (3D electronics etc) under discussion





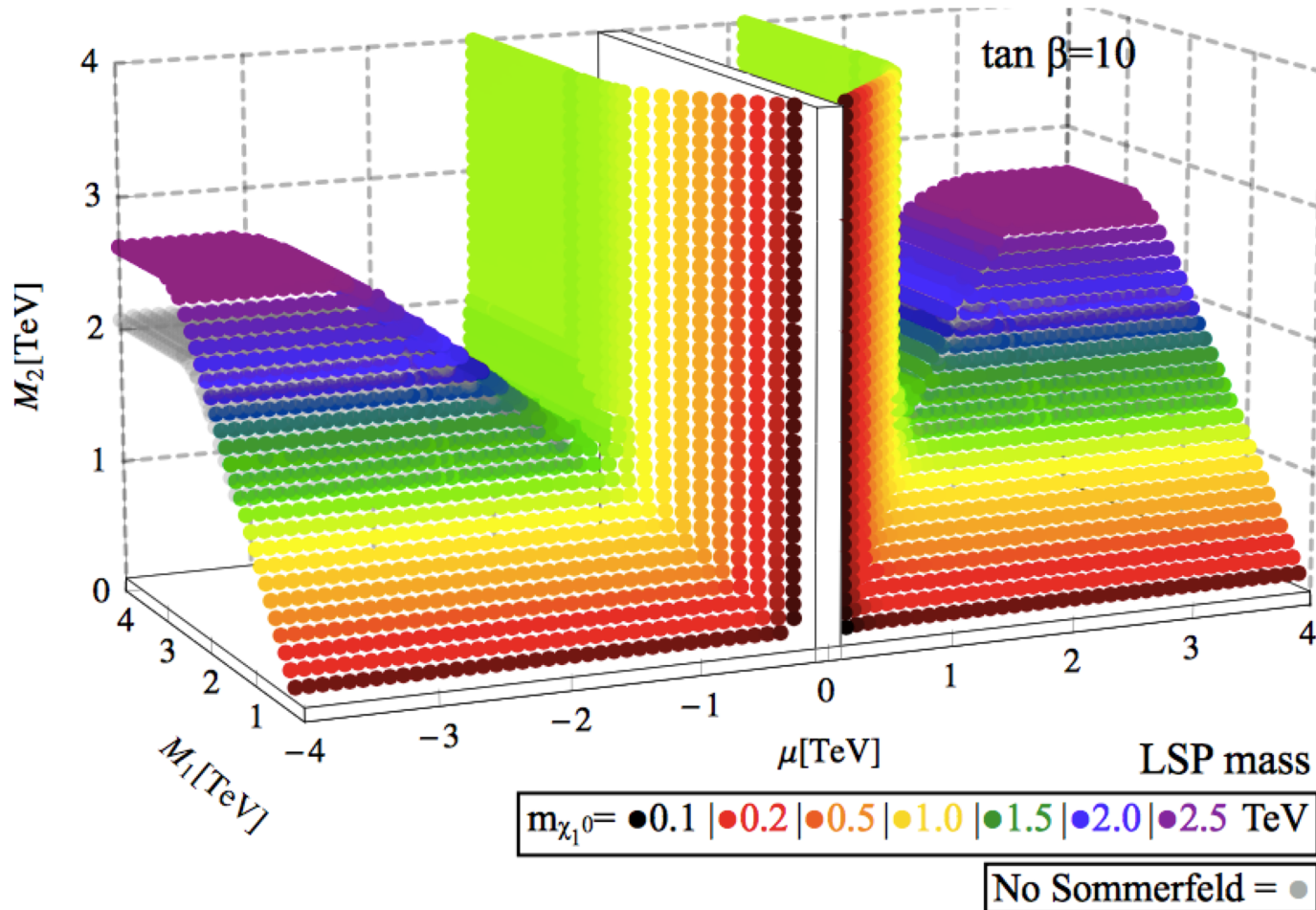
# Dark Matter

# Direct Searches for Dark Matter



# SUSY Neutralino WIMP Relic Surface

- Supersymmetric partners of photon, Z boson or Higgs boson provide generic model of weakly interacting Dark Matter
- Combinations of Neutralino mass parameters that produce the correct relic abundance, along with Dark Matter particle (LSP) mass



Bramante *et al*,

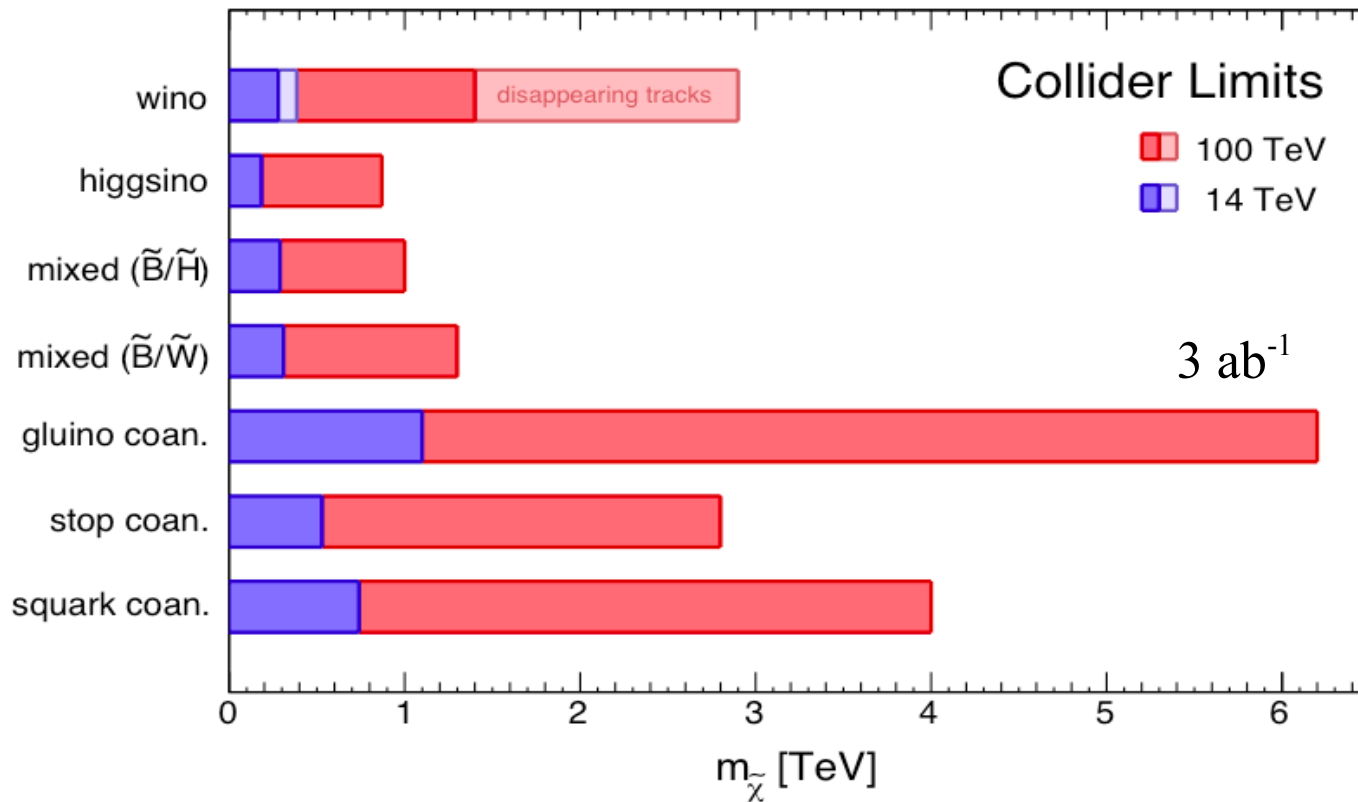
ArXiv:1510.03460

**Phys. Rev. D91 (2015)  
054015**

(in the limit that other SUSY is heavy and decoupled)

# Disappearing Track from Wino WIMP Decay

- $M_{\text{Dark Matter}} < 1.8 \text{ TeV } (g_{\text{DM}}^2/0.3)$  based on WIMP thermal relic hypothesis



M. Low, L-T Wang,  
ArXiv:1404.0682  
(mono-jet channel)

100 TeV  $pp$  collider covers most of the parameter space – **30  $\text{ab}^{-1}$  will double the mass reach**

Disappearing track: almost degenerate, long-lived  $\text{Wino}^+ \rightarrow \text{Wino}^0$   
requires robust tracking for reconstructing partial-length tracks

# Compressed Spectrum WIMPs

$$pp \rightarrow (\tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0) (\tilde{\chi}_1^\pm \rightarrow \ell^\pm \nu_\ell \tilde{\chi}_1^0) j \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 \ell^\pm \nu_\ell \gamma j$$

Bramante *et al*, **Phys. Rev. D**93 (2016) no.6, 063525

$$\begin{array}{lll} p_{T,\ell} = [10 - 60] \text{ GeV} & |\eta_\ell| < 2.5 & \\ p_{T,\gamma} = [10 - 60] \text{ GeV} & |\eta_\gamma| < 2.5 & \Delta R_{\ell\gamma} > 0.5 \\ p_{T,j} > 0.8 \text{ TeV} & |\eta_j| < 2.5 & M_{T2}^{(\gamma,\ell)} < 10 \text{ GeV} \\ \cancel{p}_T > 1.2 \text{ TeV} . & & \end{array}$$

Soft leptons and photons are crucial for this signature

# Compressed Spectrum WIMPs

$$pp \rightarrow (\tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0) (\tilde{\chi}_1^\pm \rightarrow \ell^\pm \nu_\ell \tilde{\chi}_1^0) j \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 \ell^\pm \nu_\ell \gamma j$$

Bramante *et al.*, Phys. Rev. D93 (2016) no.6, 063525

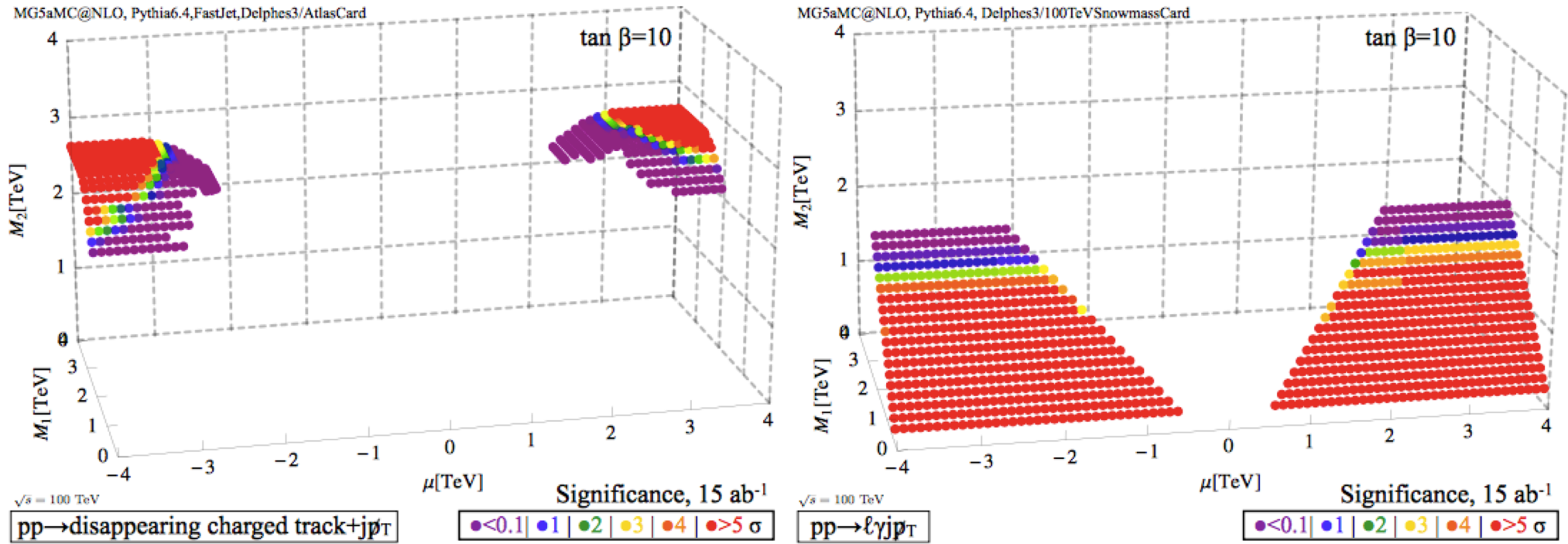


Figure 7. **Left panel:** Points on the relic neutralino surface, which will be excluded or discovered using a disappearing track search with 15  $\text{ab}^{-1}$  at a 100 TeV collider. At smaller values of  $|\mu|$  the higgsino still mixes enough to cause the mass splitting of the wino plateau to be too large for the disappearing track search to be effective. **Right panel:** Points which will be excluded or discovered using a compressed search for  $pp \rightarrow \ell^\pm \gamma jj p_T$ .



# Covering the WIMP Surface

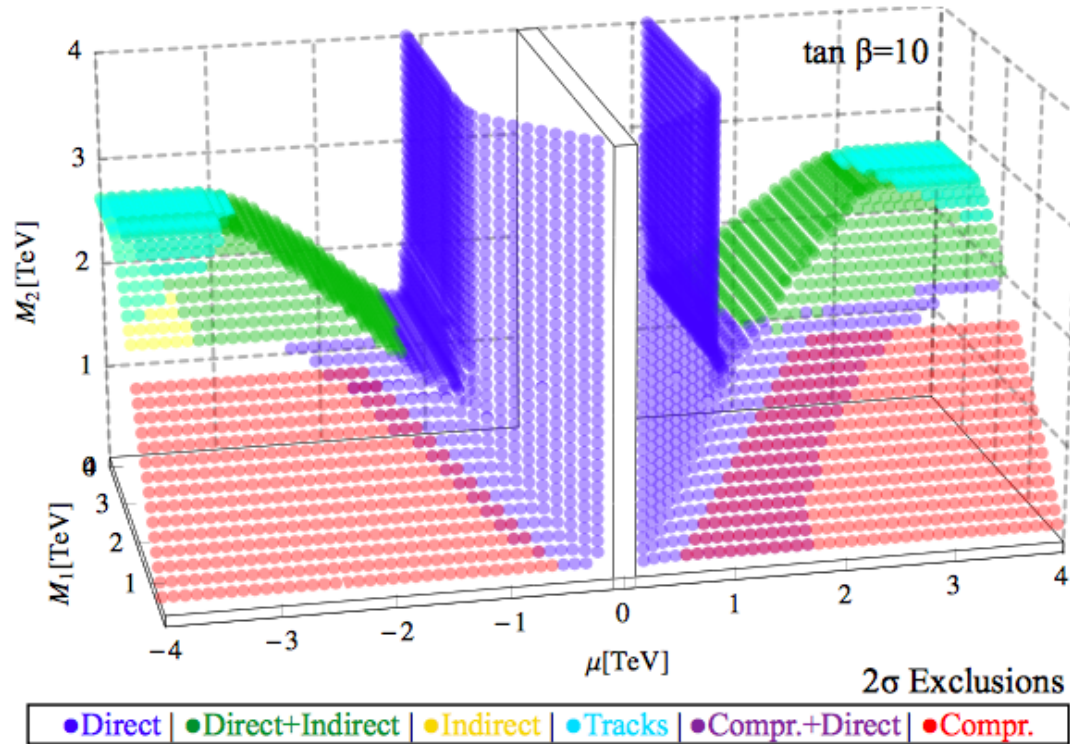


Figure 8. A combination of  $2\sigma$  exclusions from future indirect (CTA and HAWC), direct (XENON1T and LZ), and collider searches (charged tracks and compressed events at 100 TeV) are shown over the surface of thermal relic neutralinos.

Bramante *et al*, **Phys. Rev. D**93 (2016) no.6, 063525

100 TeV  $pp$  collider, combined with direct and indirect searches, covers the parameter space of WIMP satisfying relic density



# Collider vs Direct Detection Complementarity

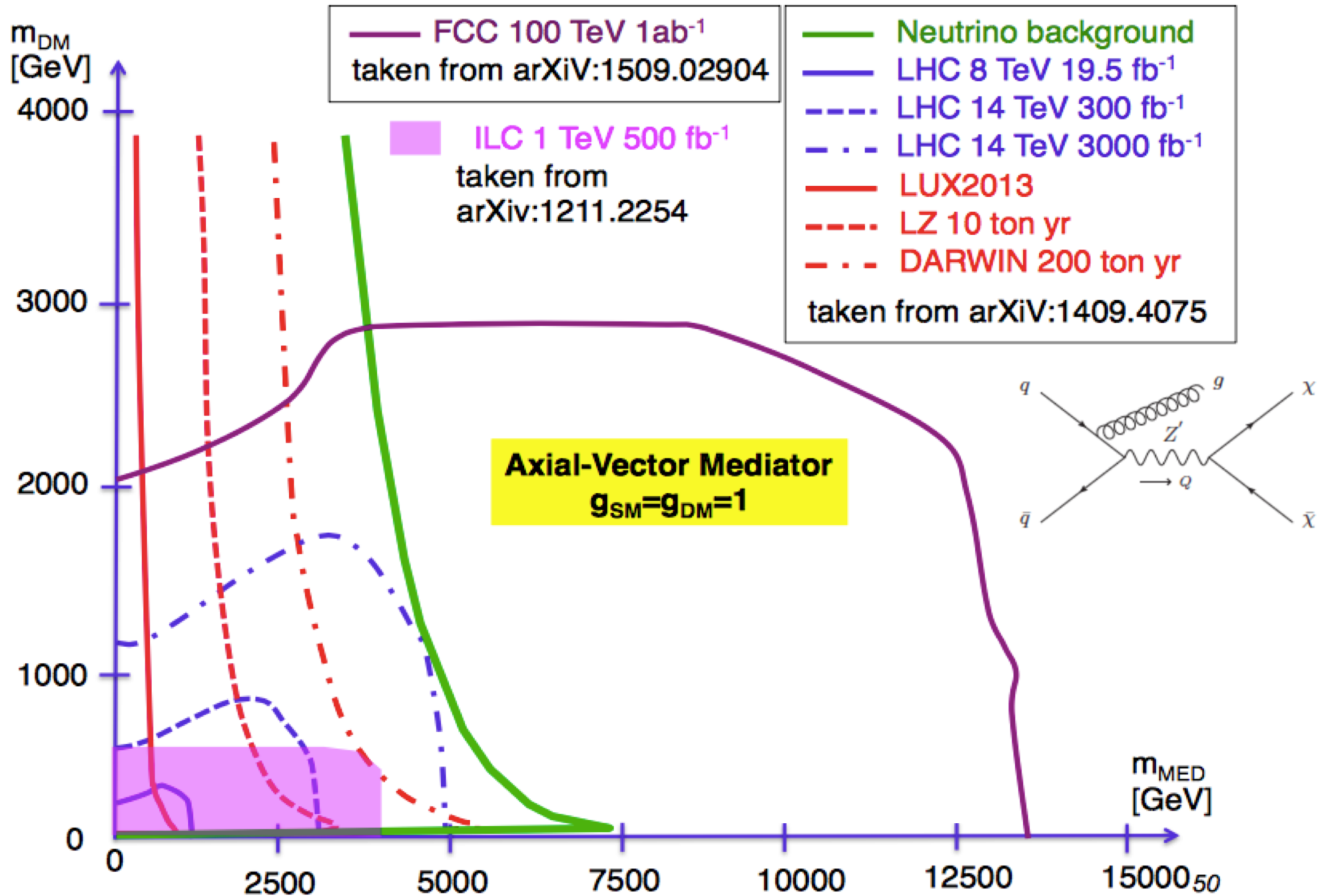
## Common ground (almost)

- **Axial-Vector mediator**  
*DD and collider are equal in overall sensitivity but probe different regions of parameter space!*
- **Scalar mediator**  
*DD and collider are equal in overall sensitivity but probe different regions of parameter space!*

## Exclusive domains (almost)

- **Vector mediator**  
*Besides very low DM masses DD wins clearly over collider*
- **Pseudo-Scalar mediator**  
*No competitive limits from DD (only from indirect detection). Collider provides limits similar in sensitivity to scalar limits*

# Collider Searches – Large Mediator Mass



# Calorimetry

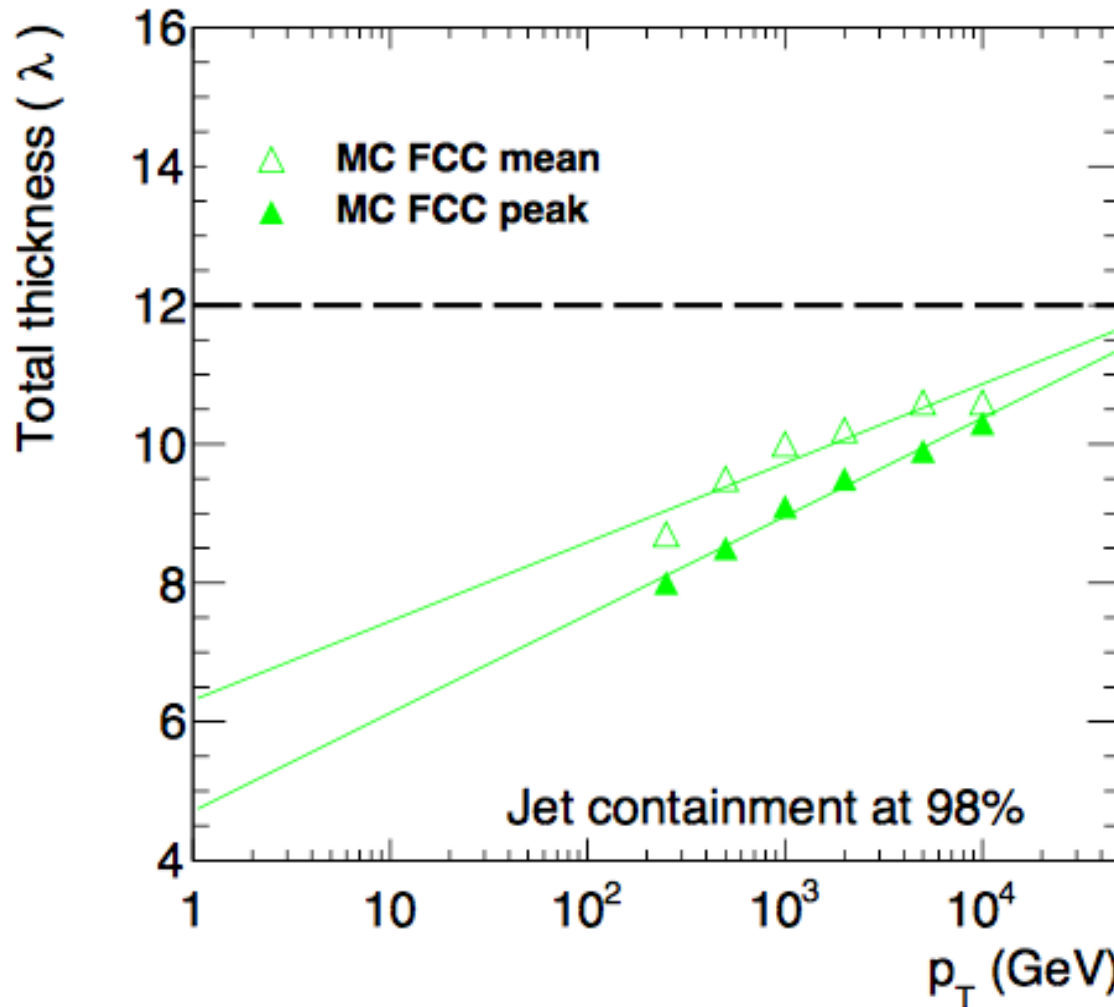
# Requirements at 100 TeV collider

The detector has to cover wide range of signatures

- Detection of high mass states
  - Dijet resonances or compositeness,  $M_{q^*} \sim 50$  TeV
  - $Z'$  or  $W'$  to leptons,  $m_{Z'} \sim 30$  TeV
  - → Deeper calorimeters, higher dynamic range
- Precision measurements of the Higgs boson properties, and Higgs in BSM production
  - Precision lepton/photon in complex events, b, c, tau tagging
  - → at least comparable to CMS/ATLAS in EM resolution and PID
- Vector boson fusion and scattering
  - Forward jets → more forward coverage, up to  $\eta=6$
- Boosted jets from Z, W, top and H
  - Jet substructures
  - → More granular calorimeters

# Calorimeter Geometry Issues

- Conveniences for going to higher energy:
  - Shower depth for full containment grows as  $\log(E)$
  - Energy resolution improves as  $\sqrt{E}$

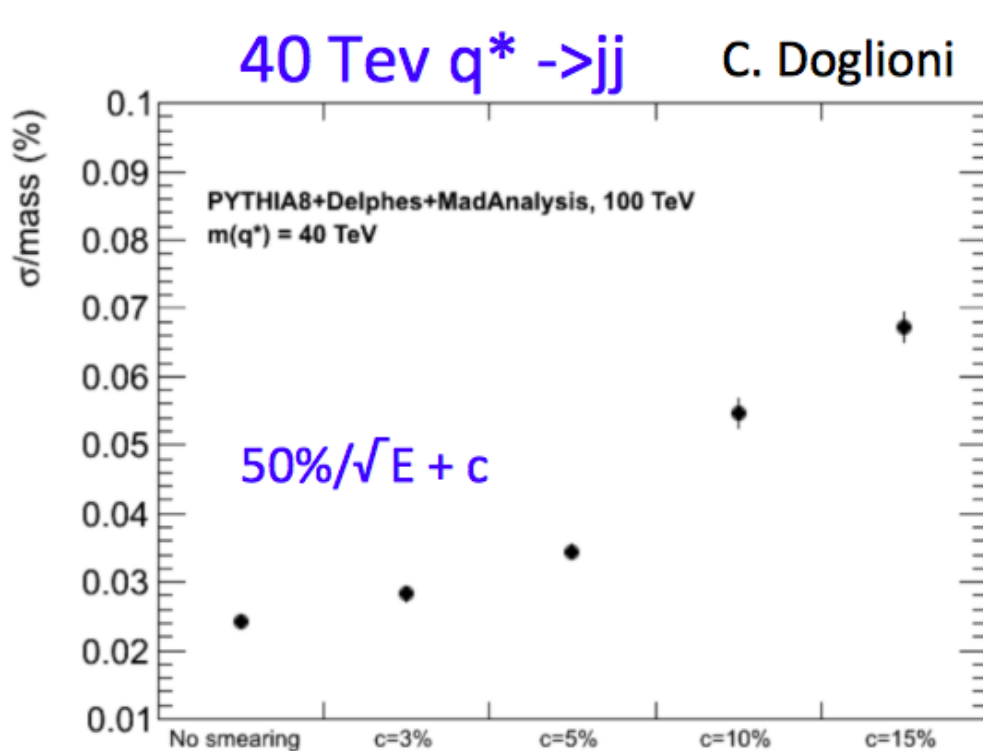


T. Carli *et al*,  
arXiv:1604.01415

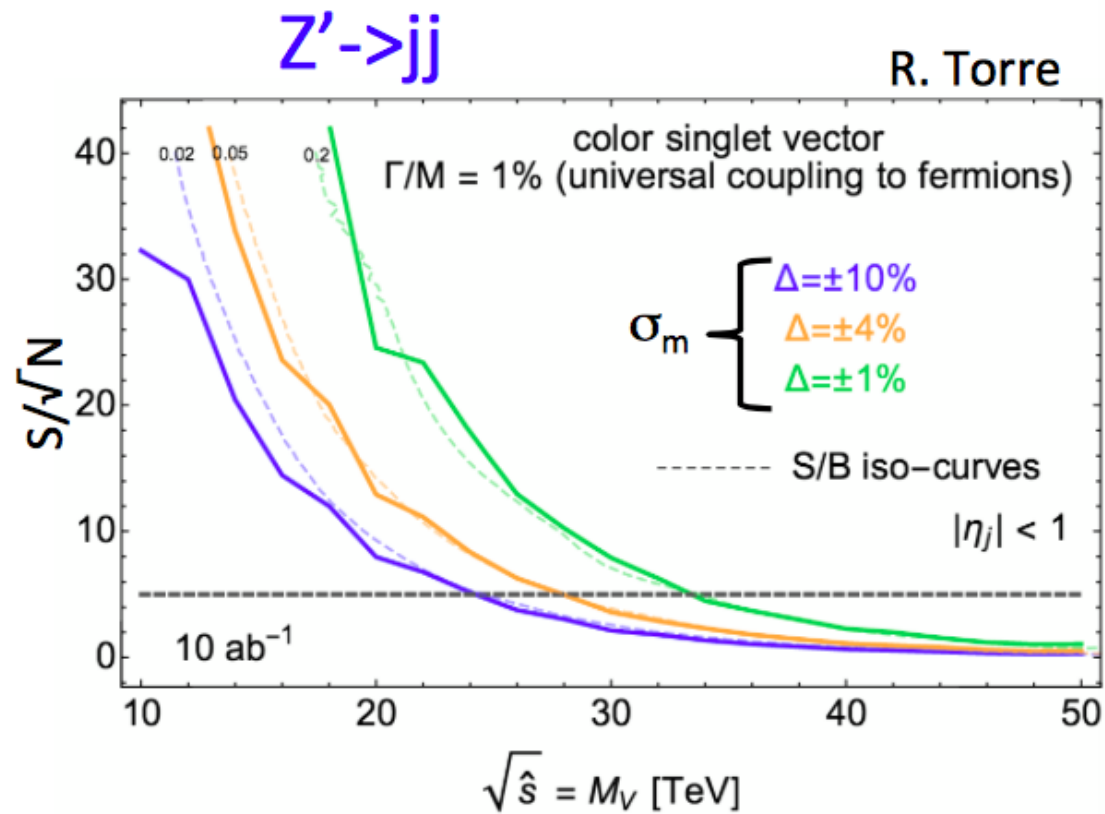
11-12 interaction lengths  
needed – space constraints  
(coil radius is expensive)

- Dynamic range of electronics readout required scales linearly with collider energy

# Effect of HCAL Energy Resolution on Dijet Resonances



HCAL resolution constant term  $C$



Jet resolution  $\sim 2\text{-}3\%$  needed for multi TeV dijet resonances

- Extend  $Z' \rightarrow jj$  discovery potential by 10TeV between  $\sigma_m = 10\%$  to 1%
- Constant term will dominate at TeV energies ( $\sigma/E = a/\sqrt{E} \oplus c$ )
- Good shower containment is mandatory!

(from Ana Henriques)

# Calorimeter Granularity

- Granularity is a KEY issue: all decay products will be boosted closer together
  - 5 TeV resonance  $\rightarrow$  HH  $\rightarrow$  4  $\tau$  produces 1 TeV  $\tau$ -lepton
    - Photons within  $\tau$ -jet are separated by  $\sim 2$  mm
    - $\tau$ -leptons from Higgs separated by  $\sim 10$  cm
  - 20 TeV resonance  $\rightarrow$   $tt$ , top decay products separated by  $\sim 3$  cm
  - 10 TeV Zprime  $\rightarrow$  WW, boosted W  $\rightarrow$  jets separated by  $\sim 3$  cm
- Tracking particles inside jets can be crucial
- Exploit particle flow algorithms to the fullest, push experience from CMS and ILC detector design effort



# GEANT Simulations

- Strategy:
  - Focus on high-granularity calorimeters
  - Resolve highly-boosted vector and Higgs bosons, top quarks,  $\tau$ -leptons
- GEANT4 simulations with ILCSoft (installed by S. Chekanov at Argonne with some help from SLAC, PNNL)
- Geometry tuning and sample generation (Chekanov and AVK)
- Analysis by Nhan Tran (Fermilab CMS postdoc), Shin-Shan Yu (Asst. Prof. in Taiwan), Sourav Sen (Duke graduate student)
- Lindsey Gray (Fermilab CMS) is our Particle Flow Algorithm expert consultant
- Samples created on OSG on 1-week timescale – need more analysts !

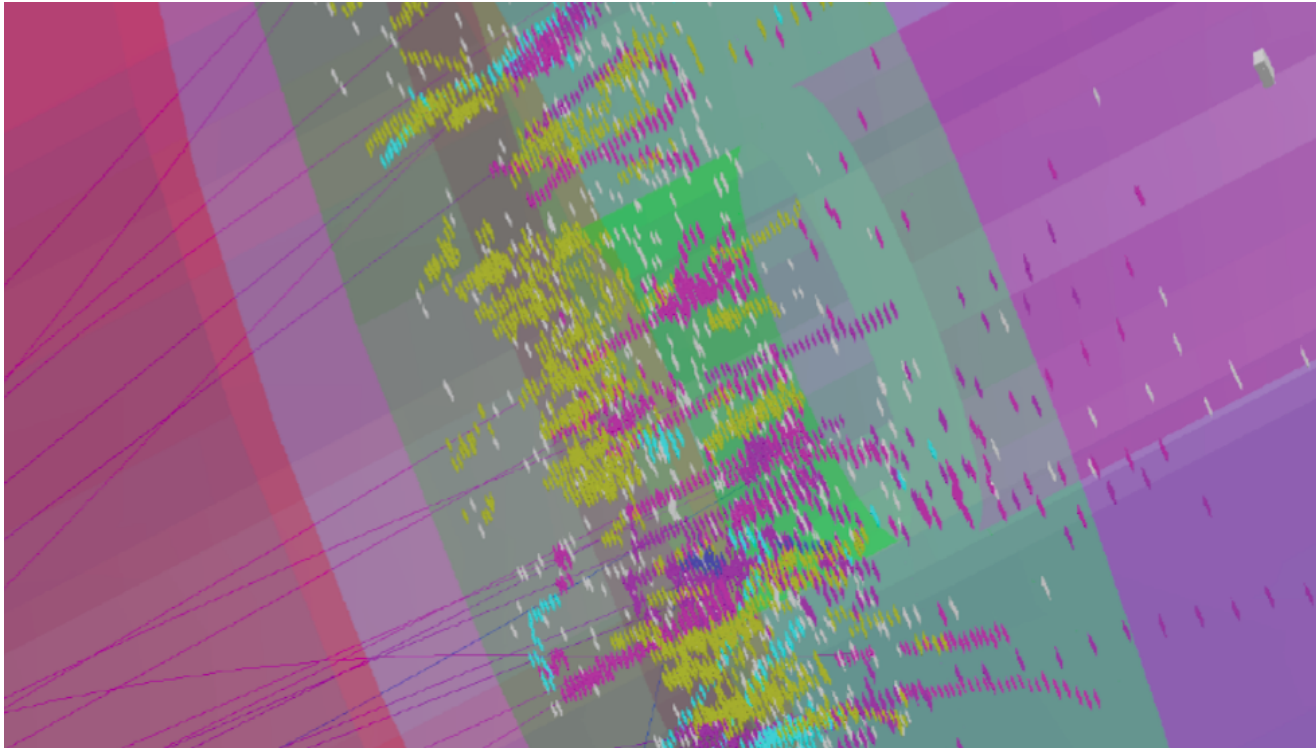
# Silicon High Granularity Calorimeter

Good cluster energy resolution

Very detailed topographical information

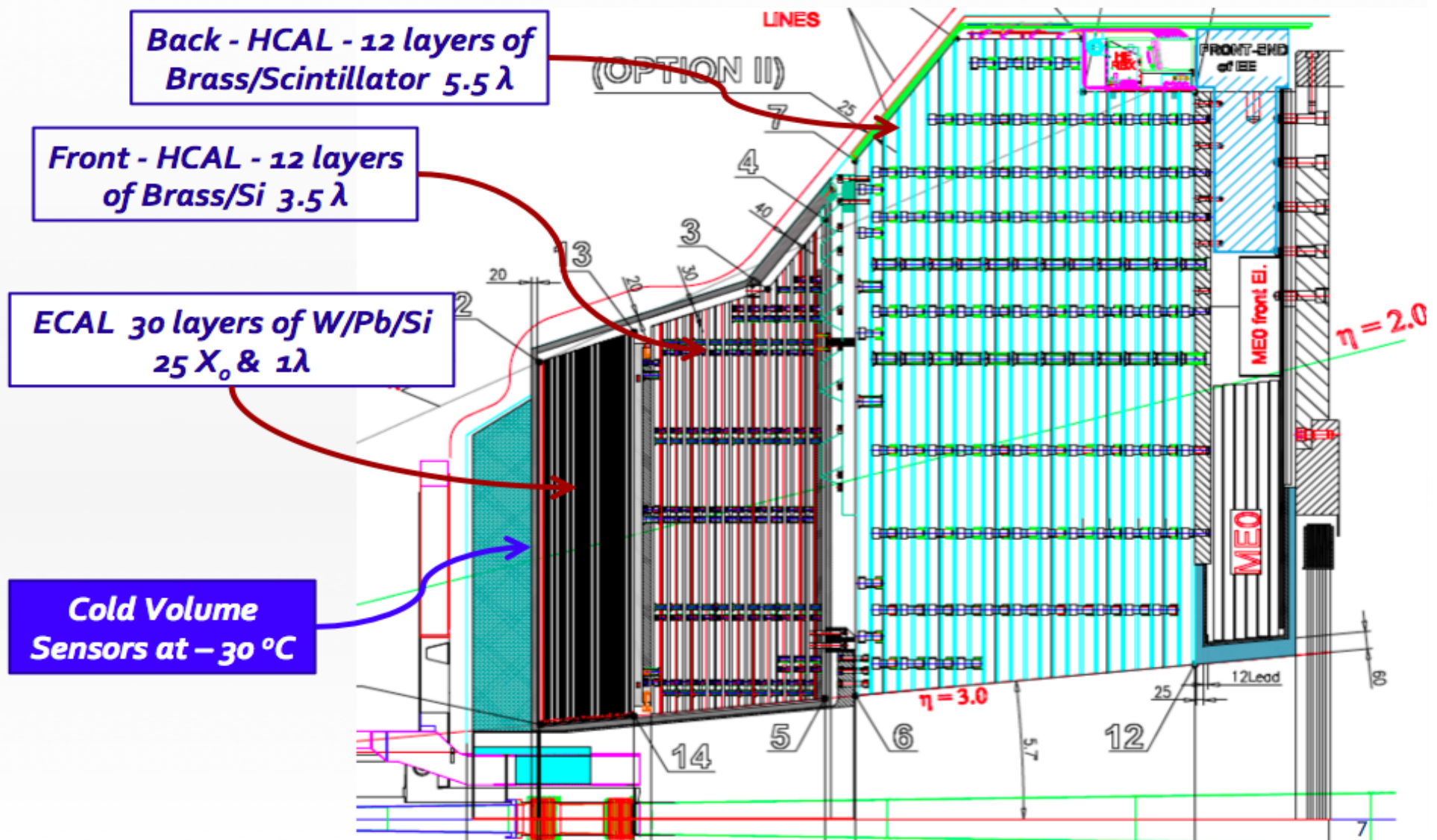
Excellent two particle cluster resolving power

Suitable for particle flow reconstruction in a high particle density environment



# Proposal – Si-HGC for CMS Endcap

## CMS Calorimeter Concept



# **Geant4 simulation of a high-granular calorimeter for TeV-scale boosted particle**

**S. Chekanov**  
*HEP/ANL*

*FCC Week. April 11-15, 2016*  
*Rome, Italy*

*With contributions from:*

A.Kotwal (Fermilab/Duke), L.Gray (Fermilab), J.Strube (PNNL), N.Tran (Fermilab), S. Yu (NCU), S.Sen (Duke), J.Repond (ANL), J.McCormick (SLAC), J.Proudfoot (ANL), A.M.Henriques Correia (CERN), C.Solans (CERN), C.Helsens (CERN)

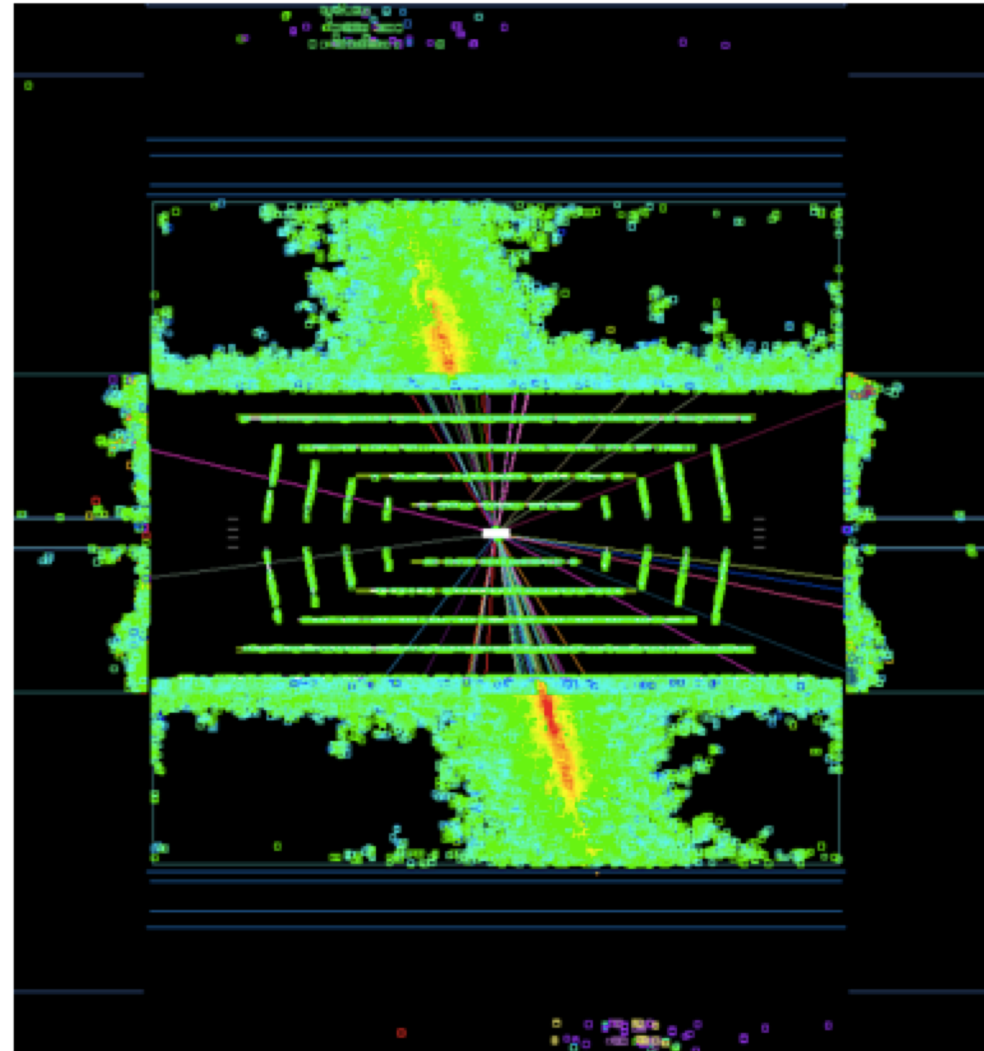
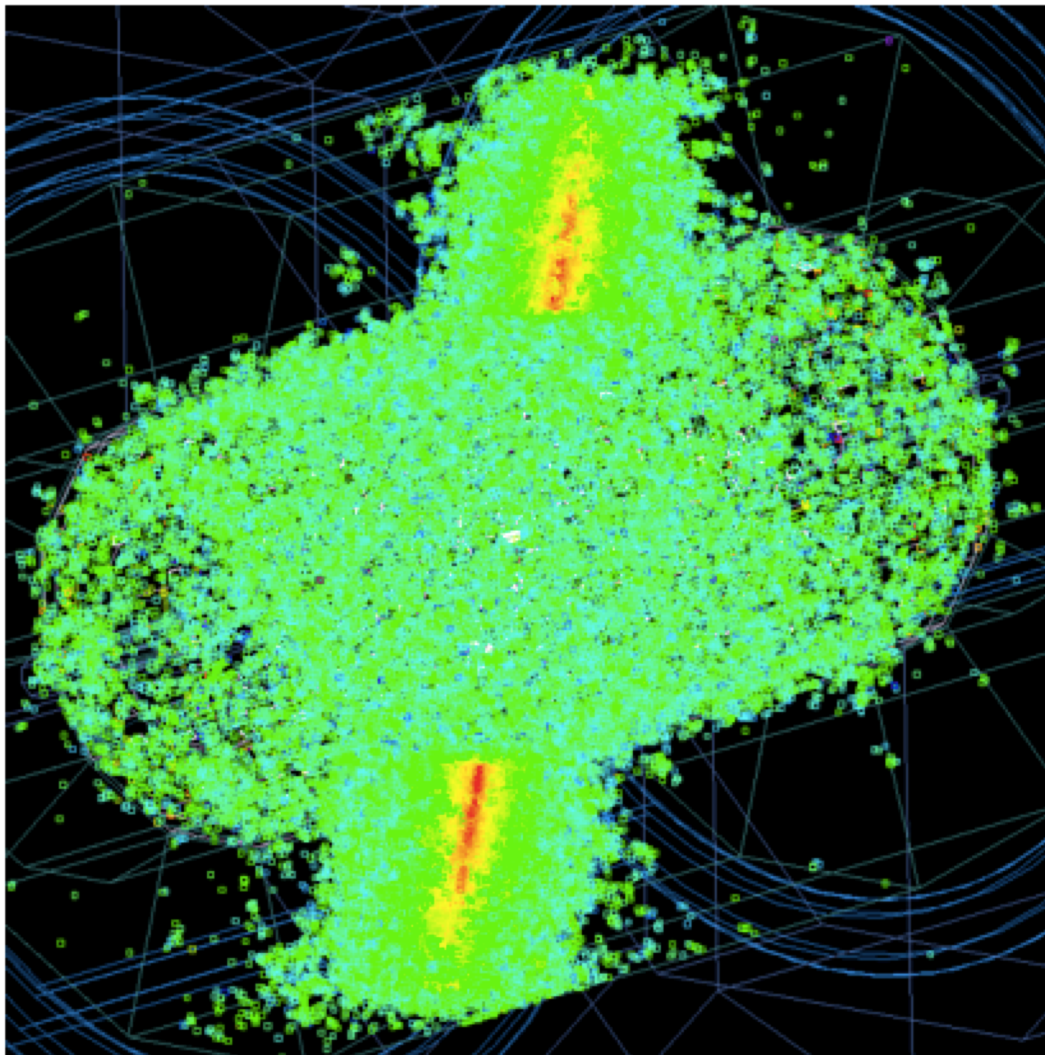
See Sergei Chekanov's talk in BOOST2017



# GEANT Simulation of Scintillator / Iron HCAL and Silicon Tracker

5 TeV hadronic  $W \rightarrow$  dijet decay with 4 cm x 4 cm scintillator readout

Background simulation in progress, will investigate different pad sizes and higher  $p_T$



Generated on OSG by S. Chekanov

# GEANT Simulation of Silicon/Tungsten EM Calorimeter

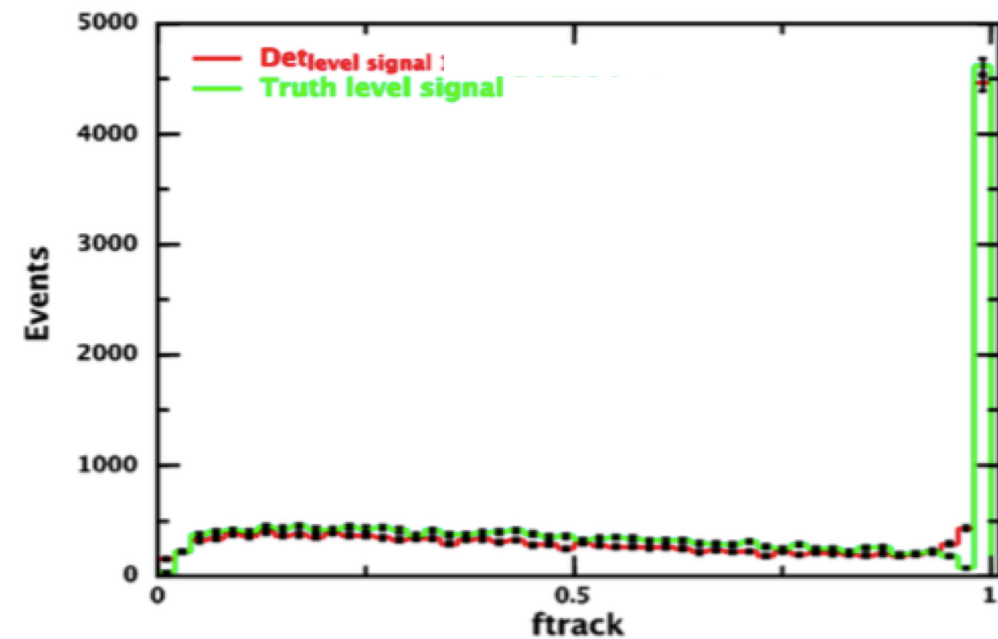
500 GeV hadronic  $\tau$ -lepton decays with 4mm x 4mm silicon pads

Background simulation in progress, will investigate larger pad sizes and higher  $p_T$

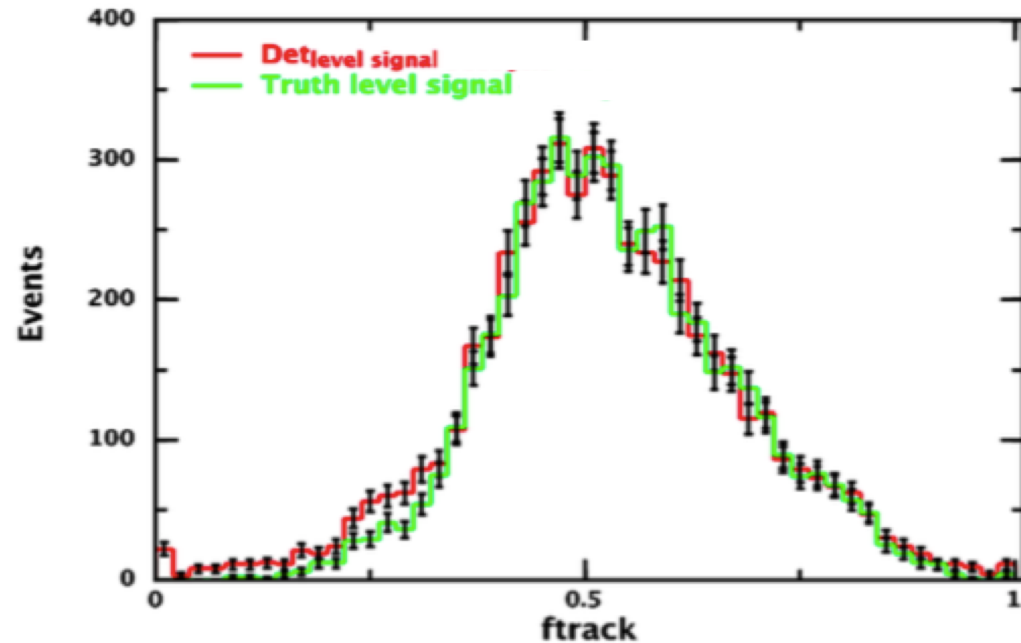
$f_{\text{track}}$  (leading track momentum fraction)

$= (\text{pT of highest pT track in core region } (\Delta R < \text{core})) / (\text{Total } E_T \text{ deposited in } \Delta R < \text{core})$

core = 0.1



1 prong



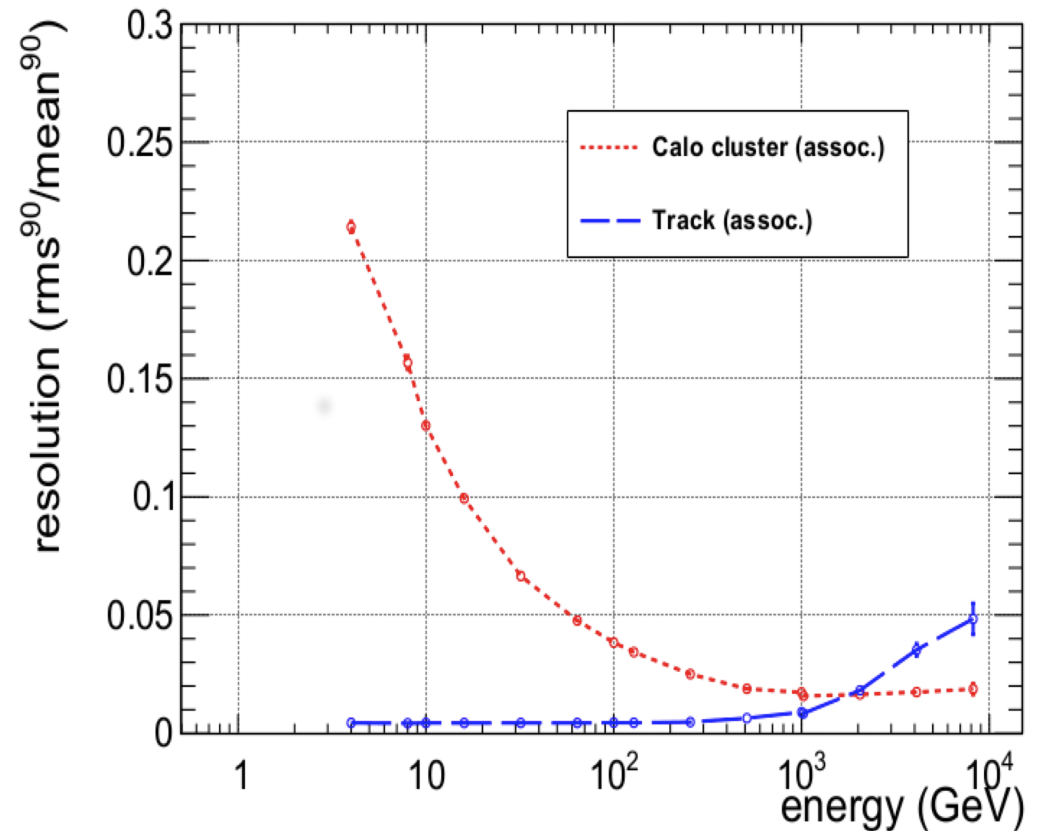
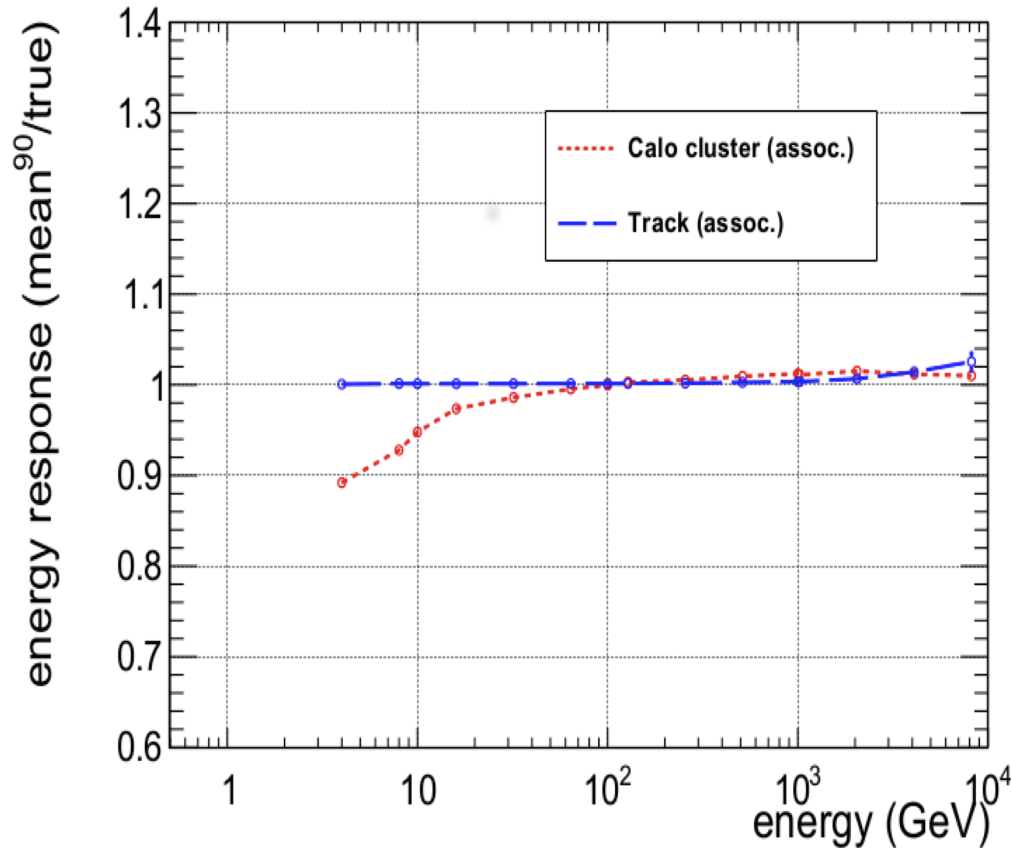
3 prong

Analysis by Sourav Sen (Duke graduate student)

Higgs  $\rightarrow \tau\tau$  is an important channel to complement  $\gamma\gamma$  and  $b\bar{b}$

# GEANT Simulation: Si/W ECAL & Scintillator/Iron HCAL

## Single pion response and resolution

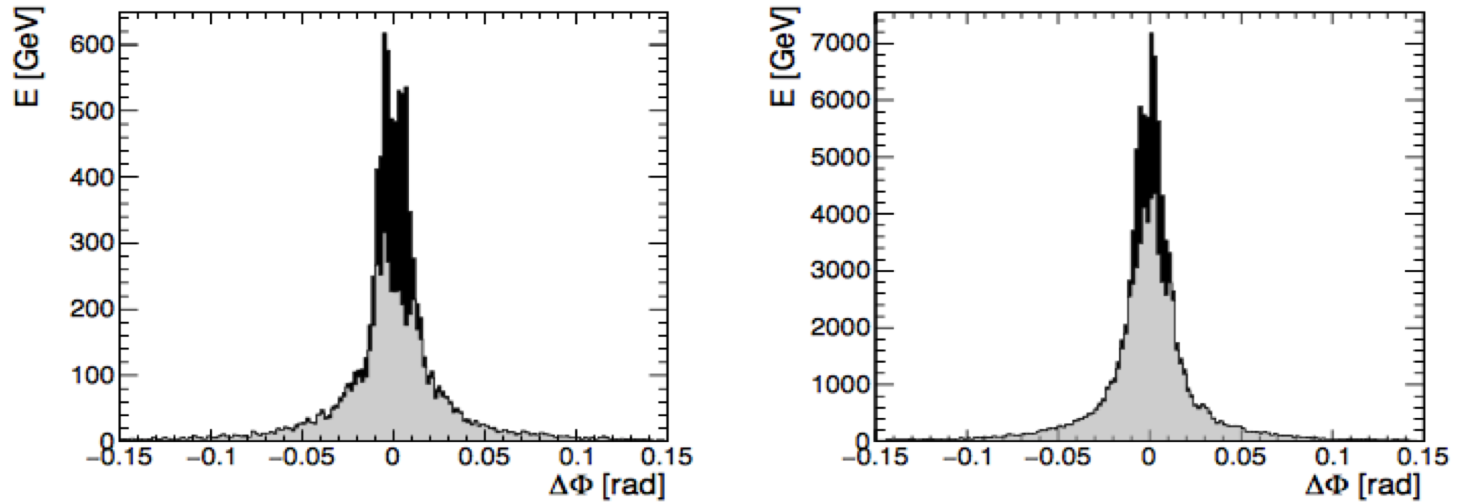


- Analysis by S. Yu, N. Tran and S. Chekanov
- First look at boosted object discriminating variables
- Published in **JINST 12 (2017) no.06, P06009**

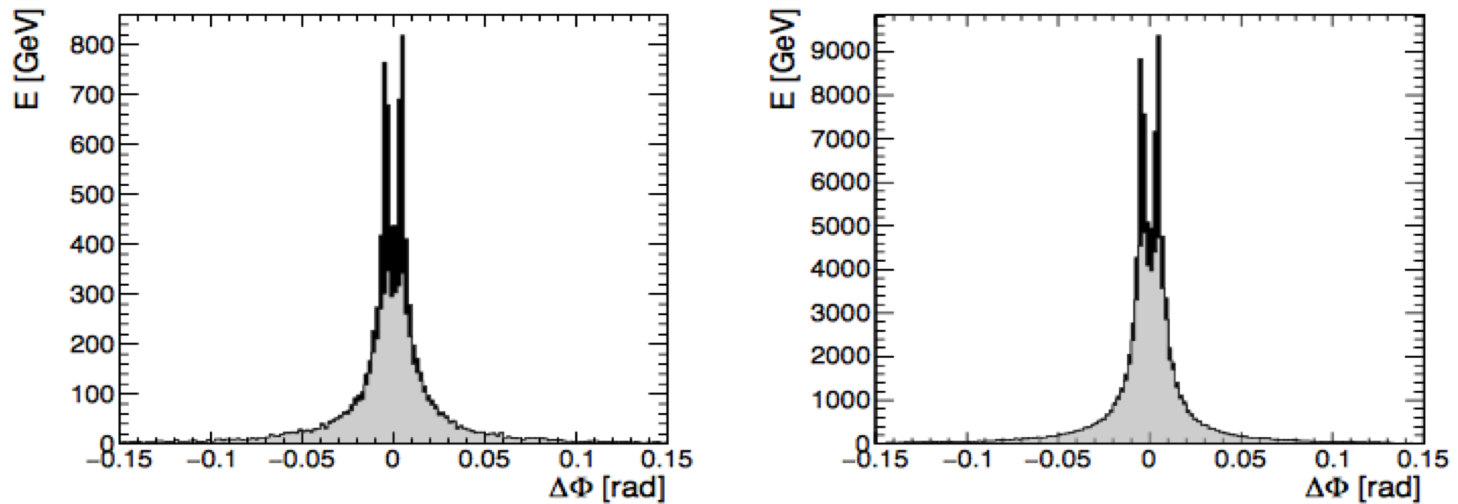


# GEANT Simulation: Silicon/Tungsten EMCAL & Iron/Scintillator HCAL

Dual  $K_L^0$  spatial separation (generated  $\Delta\phi = 10$  mrad)



(b)  $5 \times 5$  cm HCAL cells and  $2 \times 2$  cm ECAL cells



(c)  $1 \times 1$  cm HCAL cells and  $3 \times 3$  mm ECAL cells

Figure 14: Azimuthal distribution of energy deposition for pair of incident  $K_L^0$  particles at 100 GeV (left) and 1000 GeV (right), with the angular separation of  $\Delta\phi^K = 0.009$  rad. Electromagnetic calorimeter cells are indicated in black while hadronic calorimeter cells are indicated in gray.

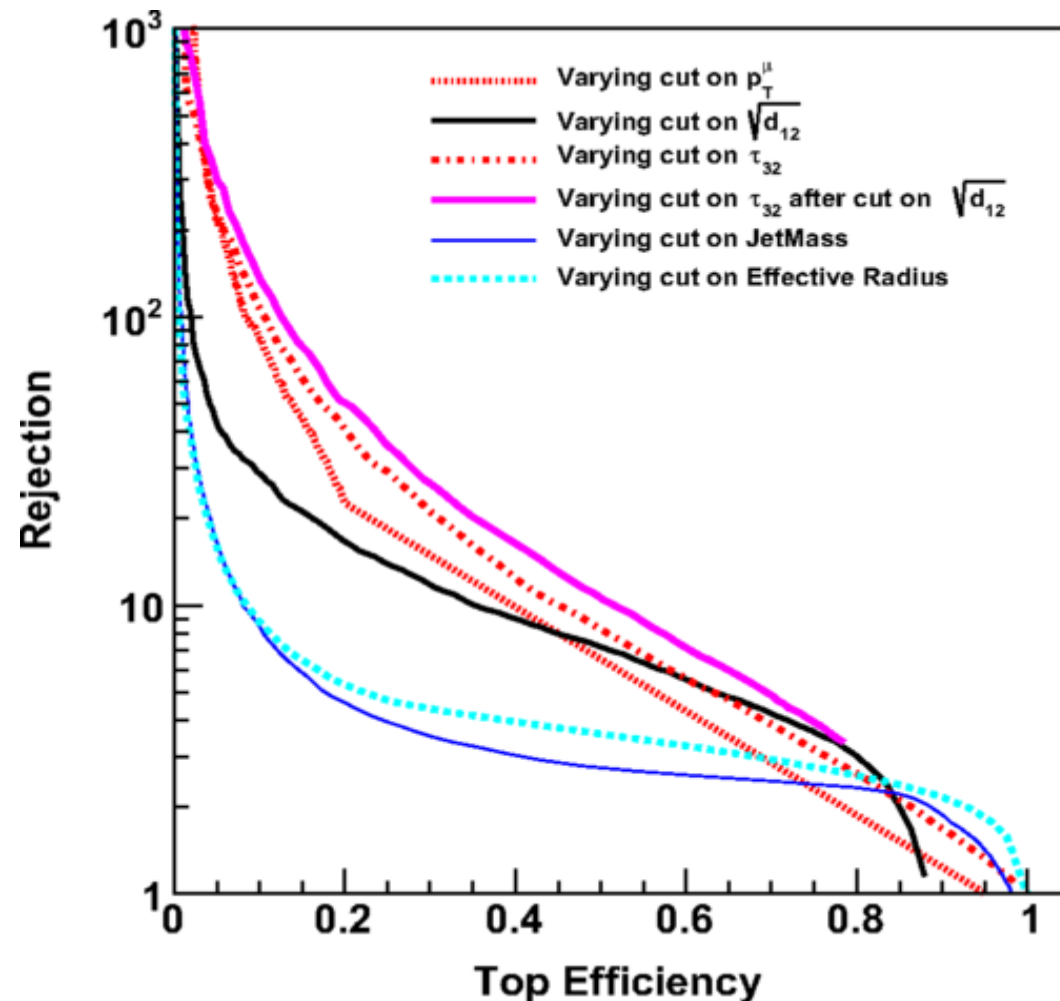
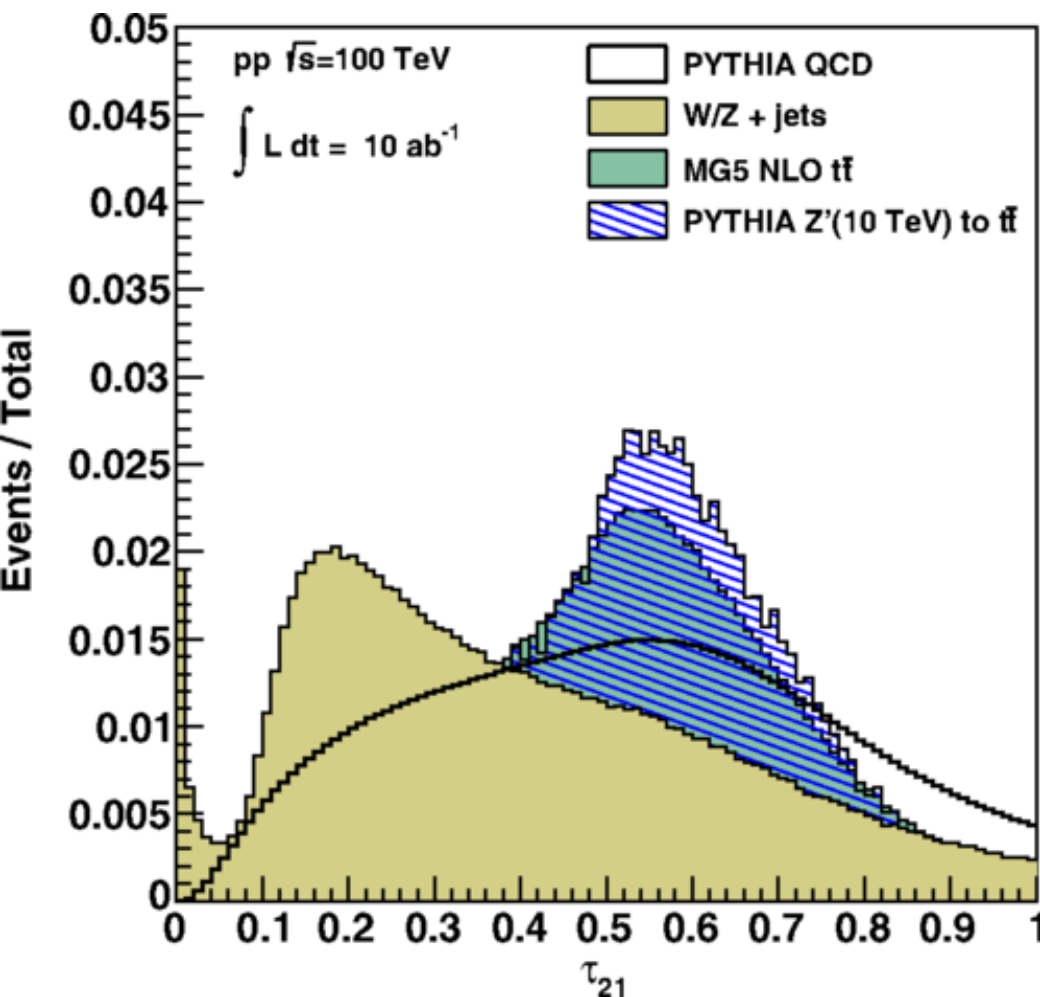
Analysis by  
Nhan Tran

# Granularity Requirements for Boosted Top Quarks

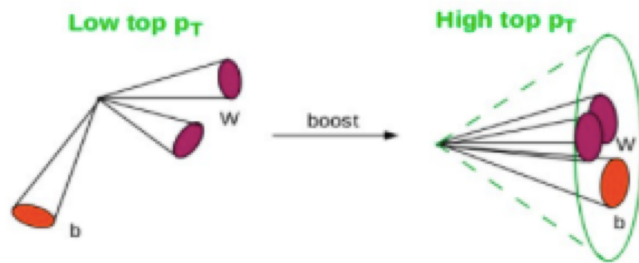
Sensitivity to new high-mass states decaying to  $t\bar{t}$  at a 100 TeV collider

B. Auerbach, S. Chekanov, J. Love, J. Proudfoot, and A. V. Kotwal  
Phys. Rev. D **91**, 034014 – Published 17 February 2015

20 TeV colored  
resonances discoverable via  
boosted hadronic decays



# Effect of HCAL transversal segmentation on jet sub-structure



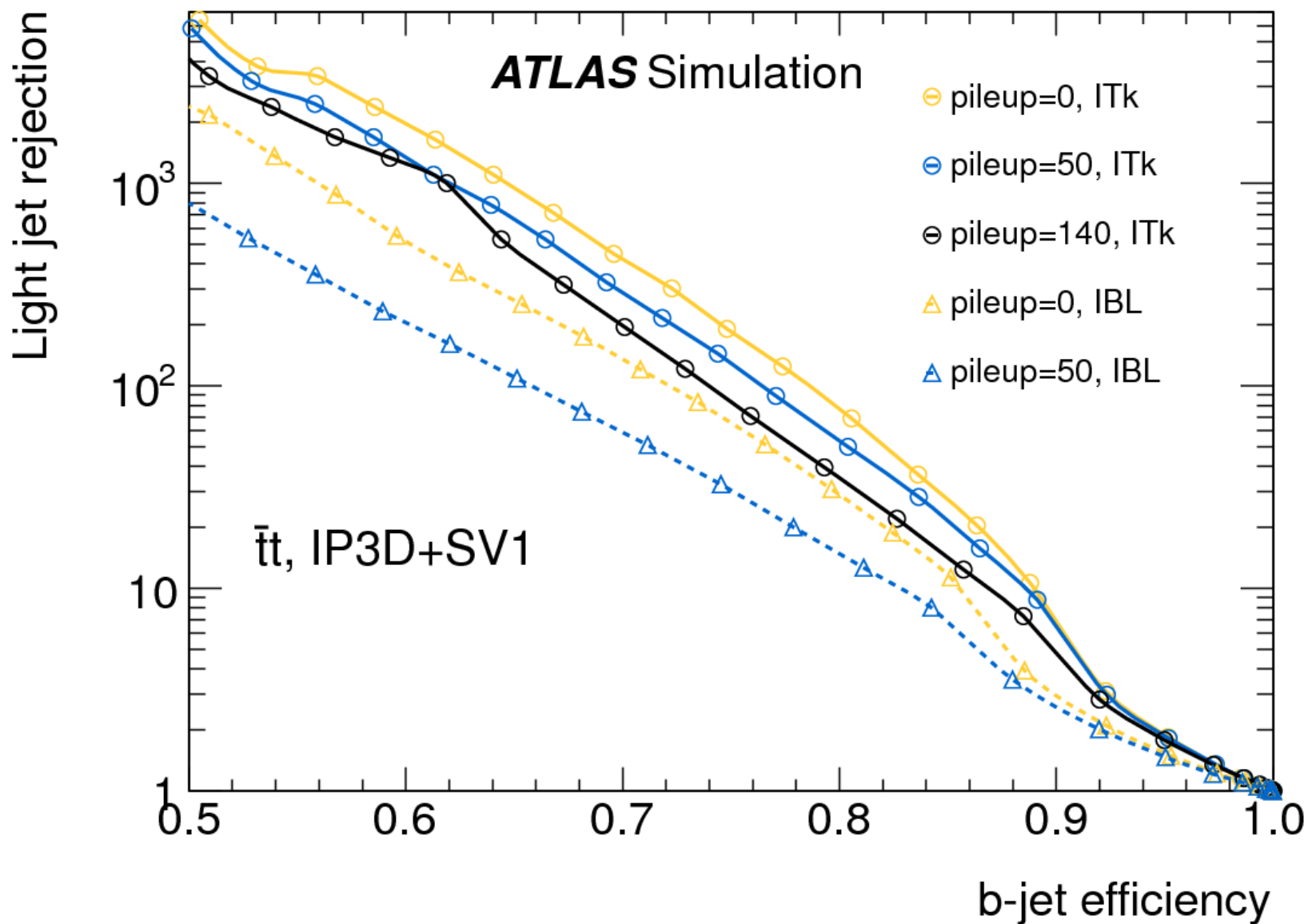
Full GEANT simulations of jet response, resolution and substructure discrimination in progress by Shin-Shan Yu, Nhan Tran, S. Chekanov *et al*

- Improve  $\sigma_m$  of sub-jettiness variables compared to  $\Delta\eta \times \Delta\phi = 0.1 \times 0.1$  for high  $P_T$  jets by:
  - 80% for  $\Delta\eta \times \Delta\phi = 0.05 \times 0.05$
  - 120% for  $\Delta\eta \times \Delta\phi = 0.025 \times 0.025$

**Need at least 2-4 times better granularity than ATLAS/CMS  $\Delta\eta \times \Delta\phi = 0.1 \times 0.1 \rightarrow 0.025 \times 0.025$**

*b*-tagging

# *b*-tagging Design Performance for HL-LHC



IBL = current,      ITk = HL-LHC design (3  $\rightarrow$  4 pixel layers, smaller pixels)

# *b*-tagging

- FCC stage 1 plans to deliver  $\sim 3 \text{ ab}^{-1}$ 
  - Similar conditions as HL-LHC, pileup  $\sim 200$  at 25 ns bunch crossing
- FCC stage 2 plans to deliver  $\sim 15 \text{ ab}^{-1}$ 
  - Pileup  $\sim 1000$ 
    - or 5 ns bunch crossing? If very fast detectors have no out-of-time pileup
- Need to achieve same *b*-tagging performance in higher-density environments
  - Highly boosted top quarks and Higgs bosons from heavy resonance decays
  - Width of b-jet  $\sim 300$  microns at 2 cm radius
  - Need to resolve tracks with factor x5 higher local density than LHC

Forward rapidity coverage



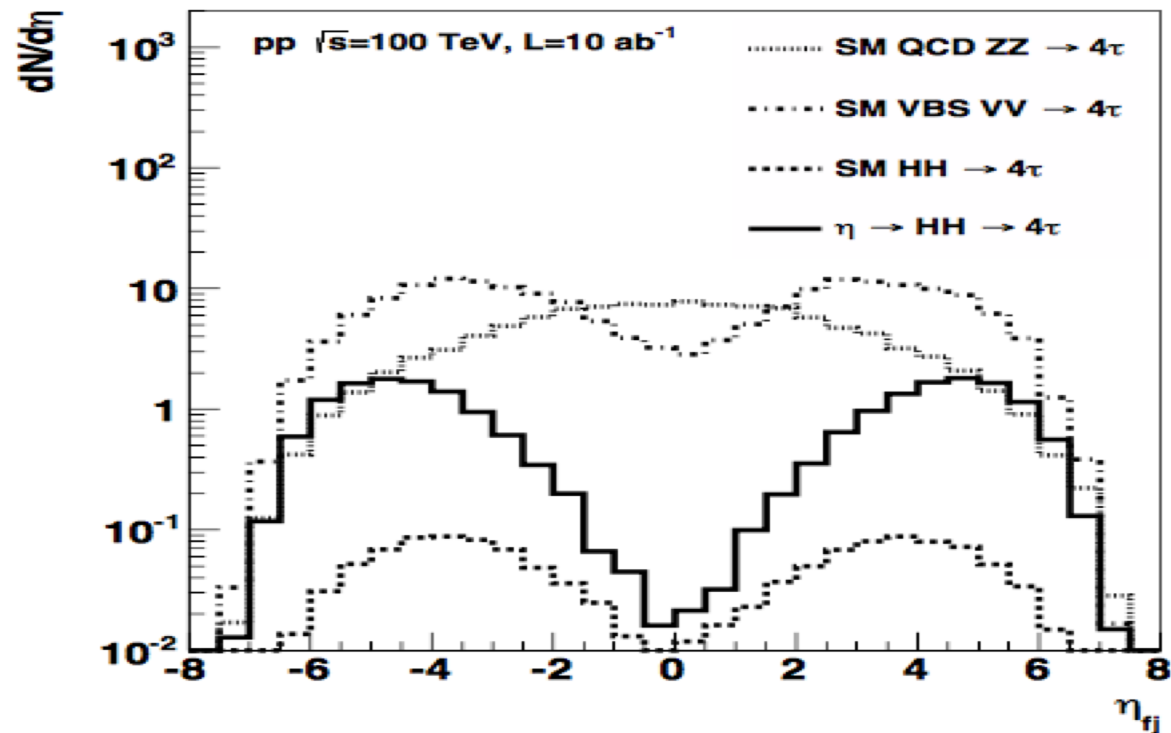
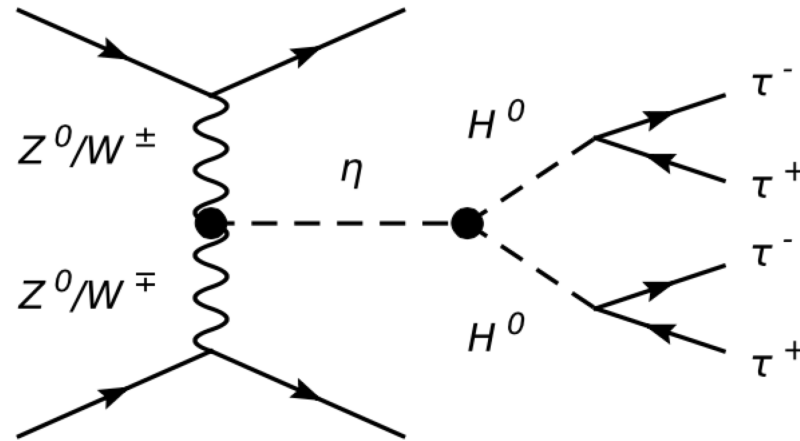
# Why is the Higgs Boson So Light?

- Old idea: Higgs doublet (4 fields) is a Goldstone mode generated from the spontaneous breaking of a larger global symmetry
  - Higgs boson and  $W_L$ ,  $Z_L$  are all Goldstone bosons from, eg. Spontaneously breaking global  $SO(5) \rightarrow SO(4)$
  - Examples: Holographic Higgs, Little Higgs models...
  - Electroweak vev “ $v$ ” is small compared to  $SO(5)$  breaking scale “ $f$ ”
- Vector boson scattering topology
  - Quarks emit longitudinal vector bosons which interact with new (presumably strong) dynamics
  - Quarks scatter by small angle in the forward direction

# Longitudinal Vector Boson Scattering

Double Higgs Boson Production in the  $4\tau$  Channel from Resonances in Longitudinal Vector Boson Scattering at a 100 TeV Collider

AVK, S. Chekanov, M. Low  
Phys.Rev. D91 (2015) 114018



(a) The pseudo-rapidity distributions of the forward jets.

# Forward Jet Coverage for Longitudinal VBS

$$V_L V_L \rightarrow \eta \rightarrow HH$$

AVK, S. Chekanov, M. Low

TABLE II.  $5\sigma$  discovery mass reach for the  $\eta \rightarrow HH \rightarrow 4\tau$  resonance, at a  $pp$  collider with  $\sqrt{s} = 100$  TeV and  $\mathcal{L} = 10 \text{ ab}^{-1}$ , for various cuts values on minimum  $p_T$  of the forward jets. The fractional width of the  $\eta$  resonance is set to  $\Gamma/M = 20\%$ .

$p_T^{\min}$ (GeV)	30	50	70	90	110
$m_\eta$ (TeV)	3.53	2.90	2.35	1.92	1.56

- Lower  $p_T$  threshold on forward tagging jets is preferred
  - Reject pileup jets with good tracking in forward direction
  - Resolve overlapping pileup jets with higher granularity / spatial resolution (*a la* CMS high-granularity endcap calorimeter for HL-LHC)

# Vector Boson Scattering

Double Higgs Boson Production in the  $4\tau$  Channel from Resonances in Longitudinal Vector Boson Scattering at a 100 TeV Collider

AVK, S. Chekanov, M. Low

Phys.Rev. D91 (2015) 114018

TABLE III.  $5\sigma$  discovery mass reach for the  $\eta \rightarrow HH \rightarrow 4\tau$  resonance, at a  $pp$  collider with  $\sqrt{s} = 100$  TeV and  $\mathcal{L} = 10 \text{ ab}^{-1}$ , for various cuts values on the maximum rapidity ( $y$ ) of the forward jets. The fractional width of the  $\eta$  resonance is set to  $\Gamma/M = 20\%$ .


$y^{\text{max}}$	8	7	6	5	4
$m_\eta$ (TeV)	2.9	2.9	2.81	2.42	1.75

Want jet rapidity coverage up to 6

# Origin of Baryon Asymmetry

POSSIBLE EXPLANATIONS...

$$\frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-9} \text{ (from BBN)}$$

⇒ **Baryogenesis at EW Scale**  **TESTABLE!**

⇒ ...

SAKHAROV CONDITIONS *(for dynamical generation of baryon asymmetry)*

*B* Violation ✓ *Sphalerons*

V. A. Kuzmin, V. A. Rubakov, M. Shaposhnikov, Phys. Lett. B155 (1985) 36

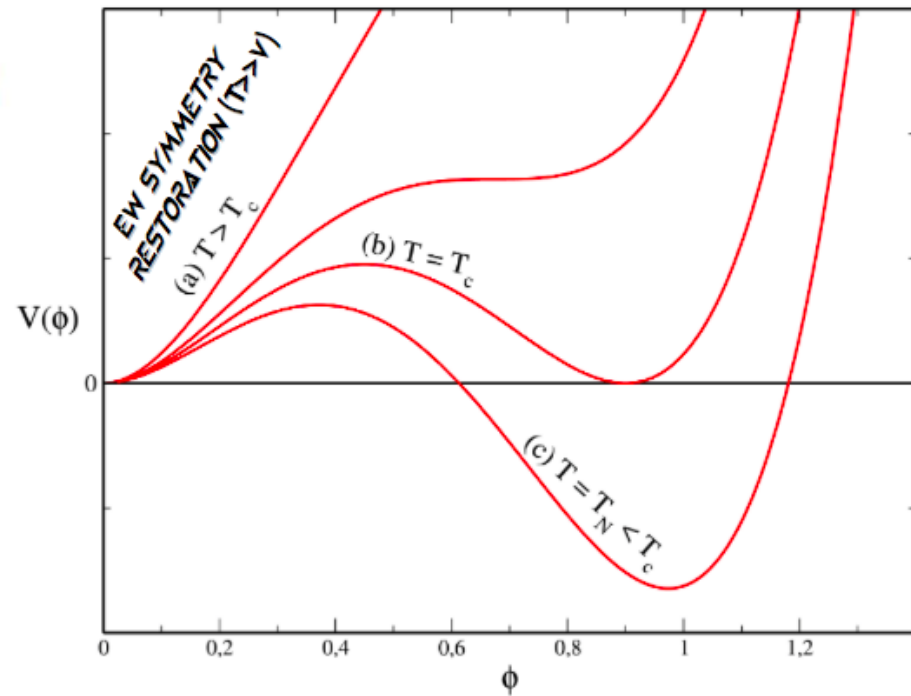
*C/CP* Violation ✗ *not enough*

Departure from Thermal Equilibrium ✗ *not enough*

# Baryon Asymmetry and Electroweak Phase Transition

1<sup>st</sup> Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$  Discontinuous

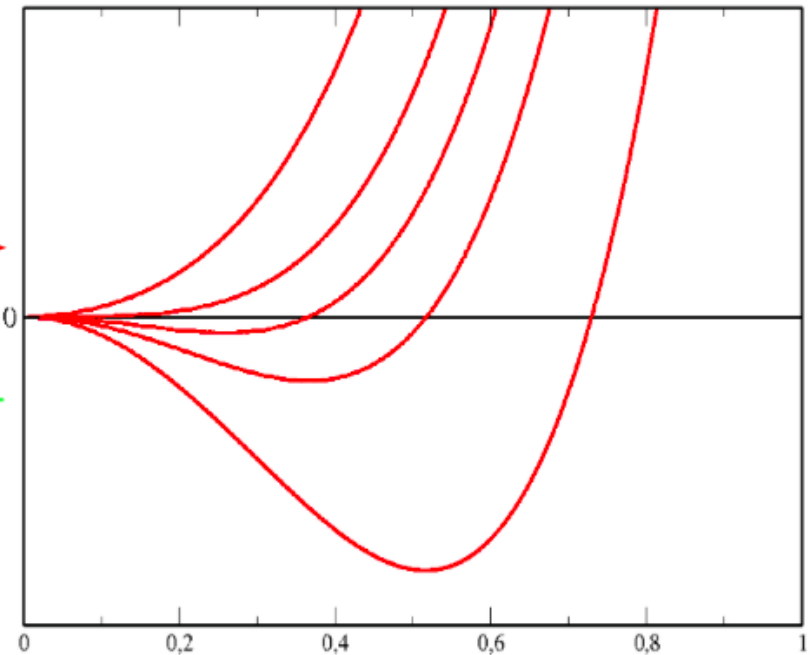


LARGER  $M_H$

NEW BOSONS

2<sup>nd</sup> Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$  Continuous



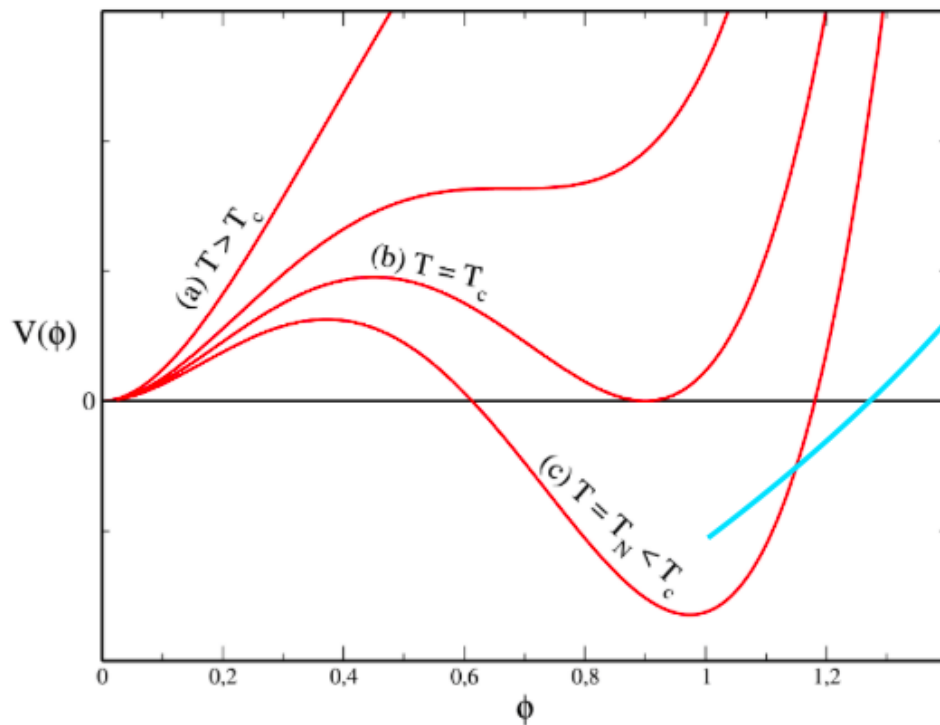
*In the SM ( $m_h = 125$  GeV) EW Phase Transition Smooth CrossOver*

*K. Kajantie, M. Laine, K. Rummukainen, M. Shaposhnikov, Phys. Rev. Lett. **77** (1996) 2887*

# Baryon Asymmetry and Electroweak Phase Transition

1<sup>st</sup> Order:

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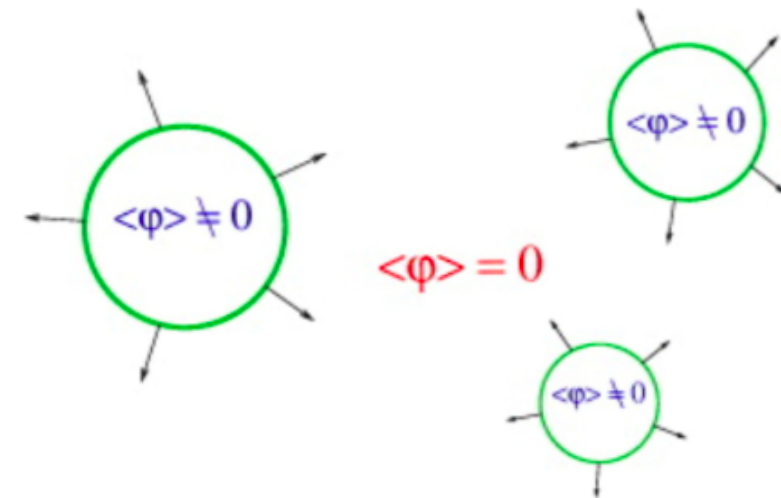


Nucleation of True Vacuum Bubbles  
(in False Vacuum Sea)

*J. S. Langer, Ann. Phys. **54** (1969) 258*

*S. R. Coleman, Phys. Rev. D **15** (1977) 2929*

*A. D. Linde, Nucl. Phys. B **216** (1983) 421*

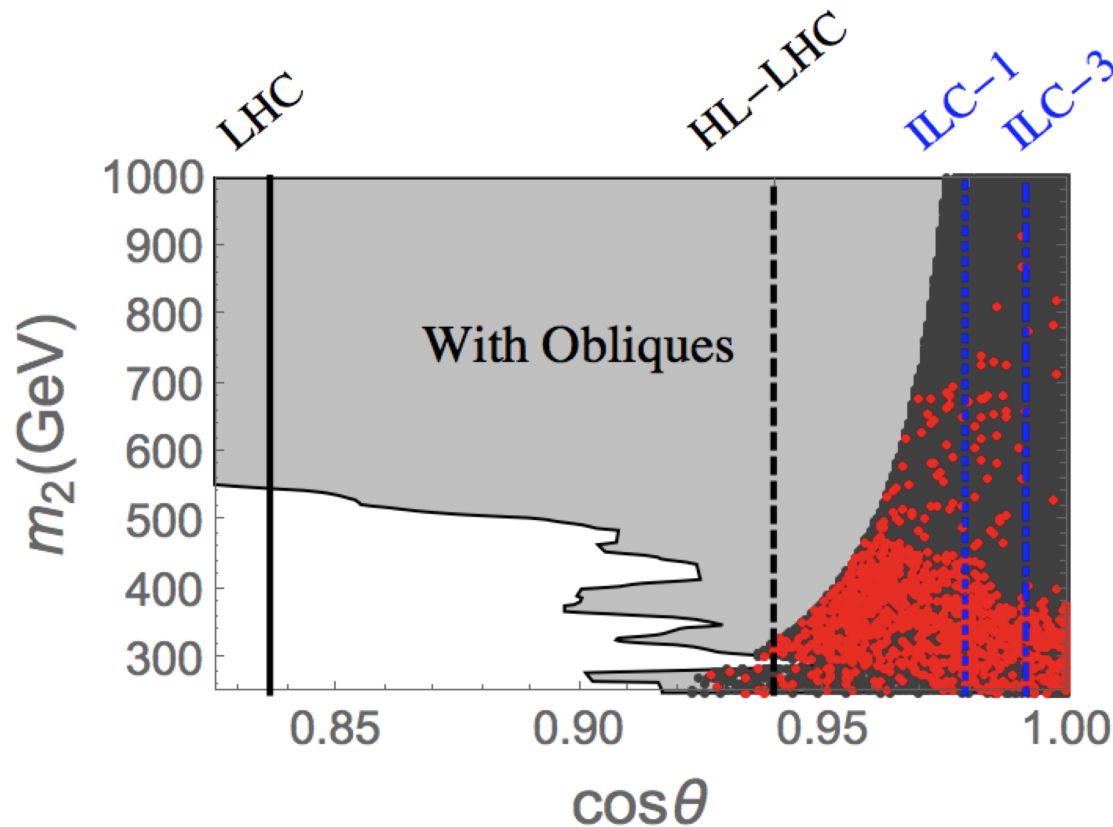


SUDDEN CHANGE IN HIGGS VEV



# First Order Phase Transition

$$V(H, S) = -\mu^2 (H^\dagger H) + \lambda (H^\dagger H)^2 + \frac{a_1}{2} (H^\dagger H) S + \frac{a_2}{2} (H^\dagger H) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4$$



(from P. Winslow)

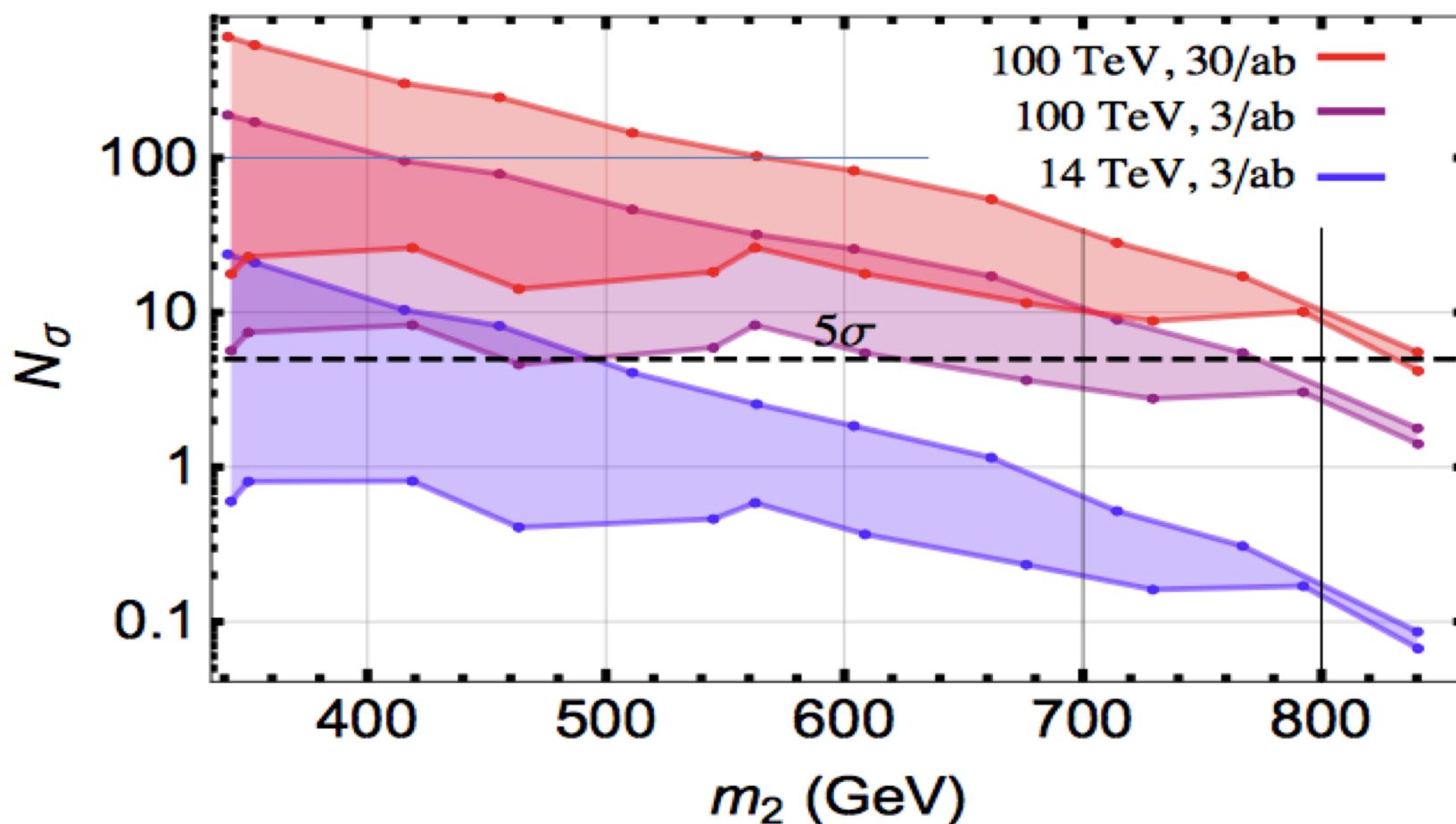
S. Profumo, M. J. Ramsey-Musolf, C. L. Wainwright and P. Winslow, arXiv:1407.5342

Can TeV-scale new physics associated with 1<sup>st</sup> order phase transition be completely covered by a  $pp$  collider?

# Inducing First-Order Electroweak Phase Transition

$$V(H, S) = -\mu^2 (H^\dagger H) + \lambda (H^\dagger H)^2 + \frac{a_1}{2} (H^\dagger H) S + \frac{a_2}{2} (H^\dagger H) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4$$

$$S \rightarrow HH \rightarrow \gamma\gamma bb \text{ and } 4\tau$$



(AVK, P. Winslow,  
J. M. No,  
M. J. Ramsey-Musolf,  
**Phys.Rev. D94 (2016),  
035022)**

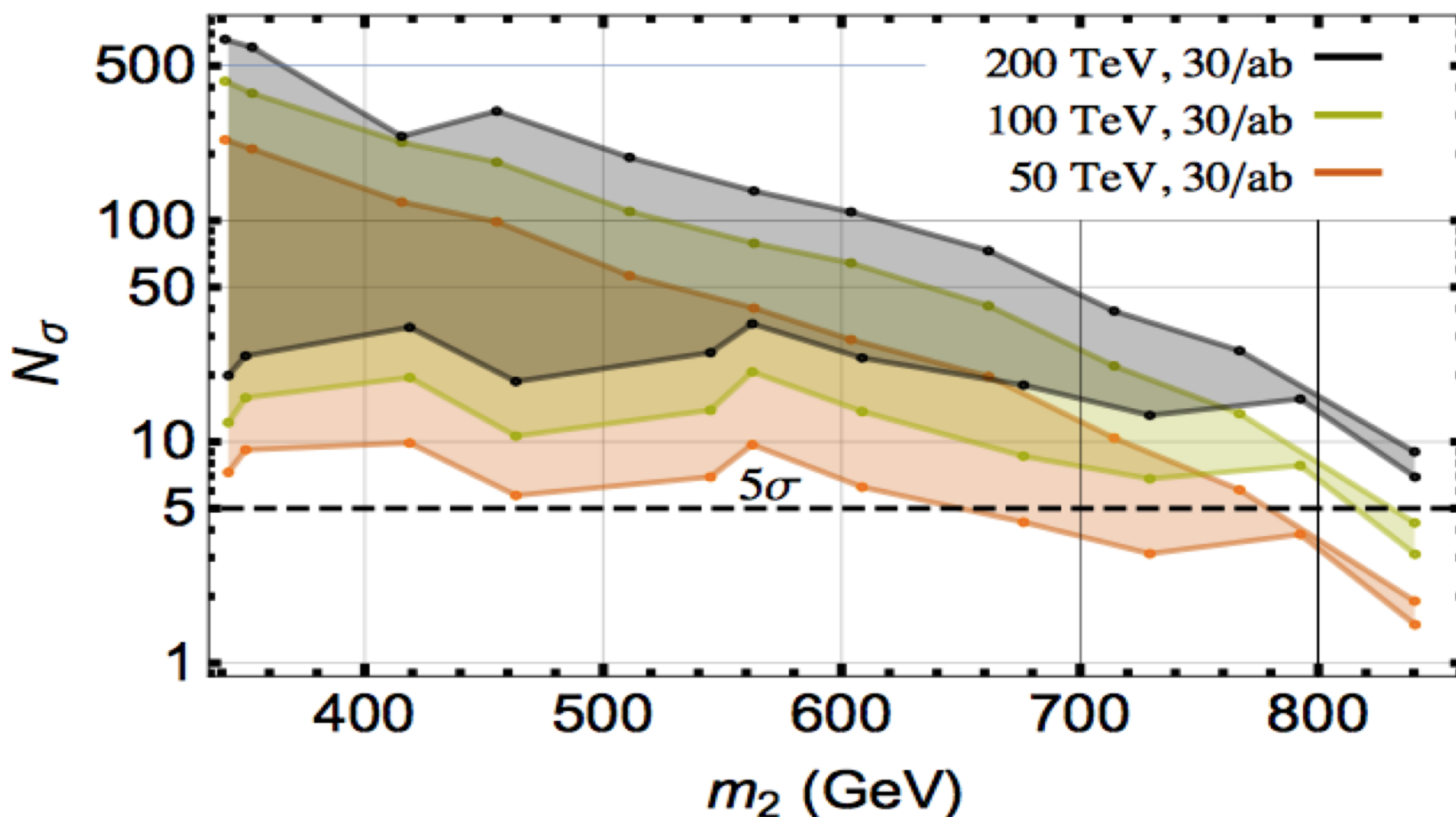
Discovery potential across entire parameter space with next collider

# Inducing First-Order Electroweak Phase Transition

Assumptions: photon,  $b$ -quark and  $\tau$ -lepton efficiency = 75%

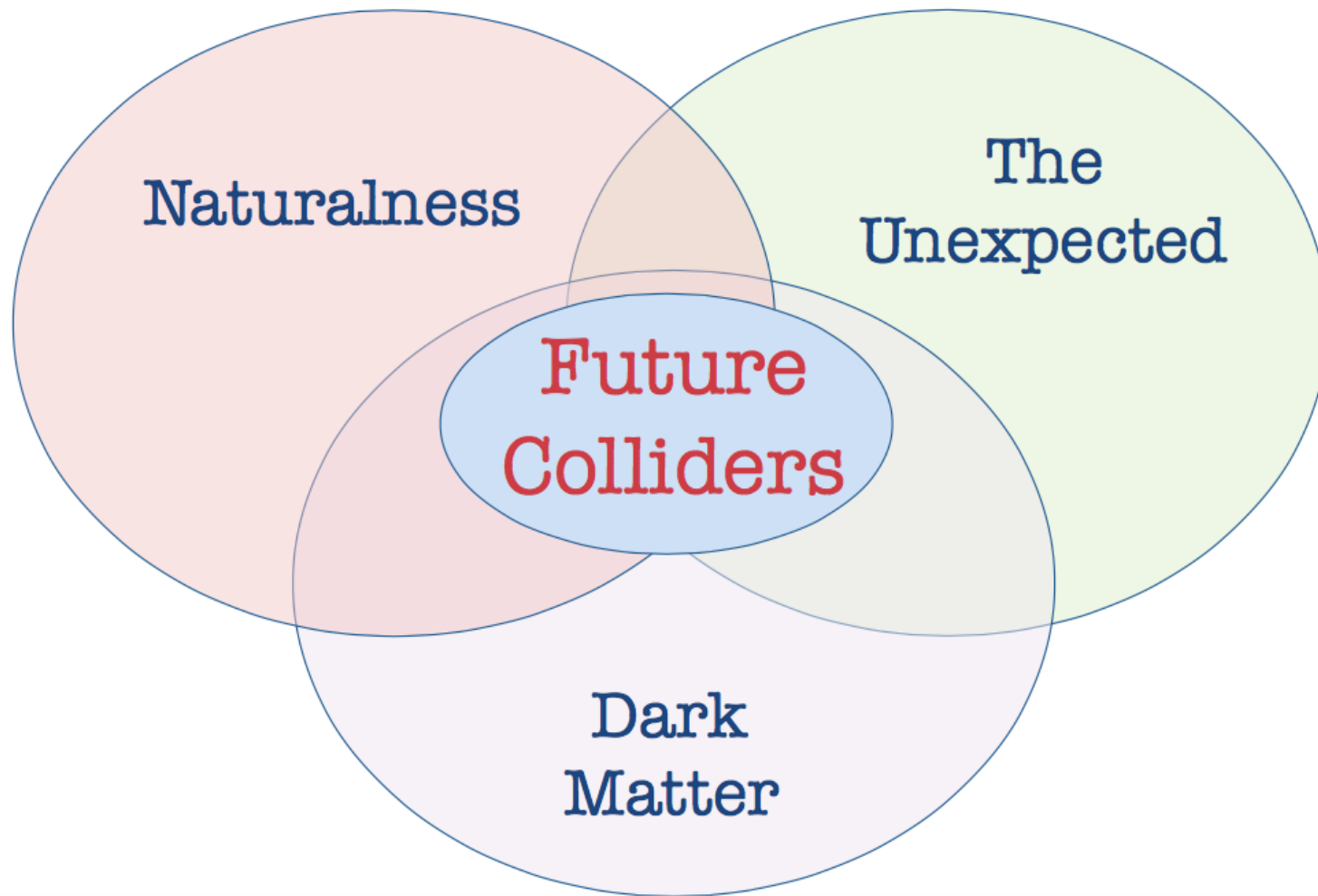
Jets  $\rightarrow b$ -quark and  $\tau$ -lepton fake rate = 2%

$S \rightarrow HH \rightarrow \gamma\gamma bb$  and  $4\tau$



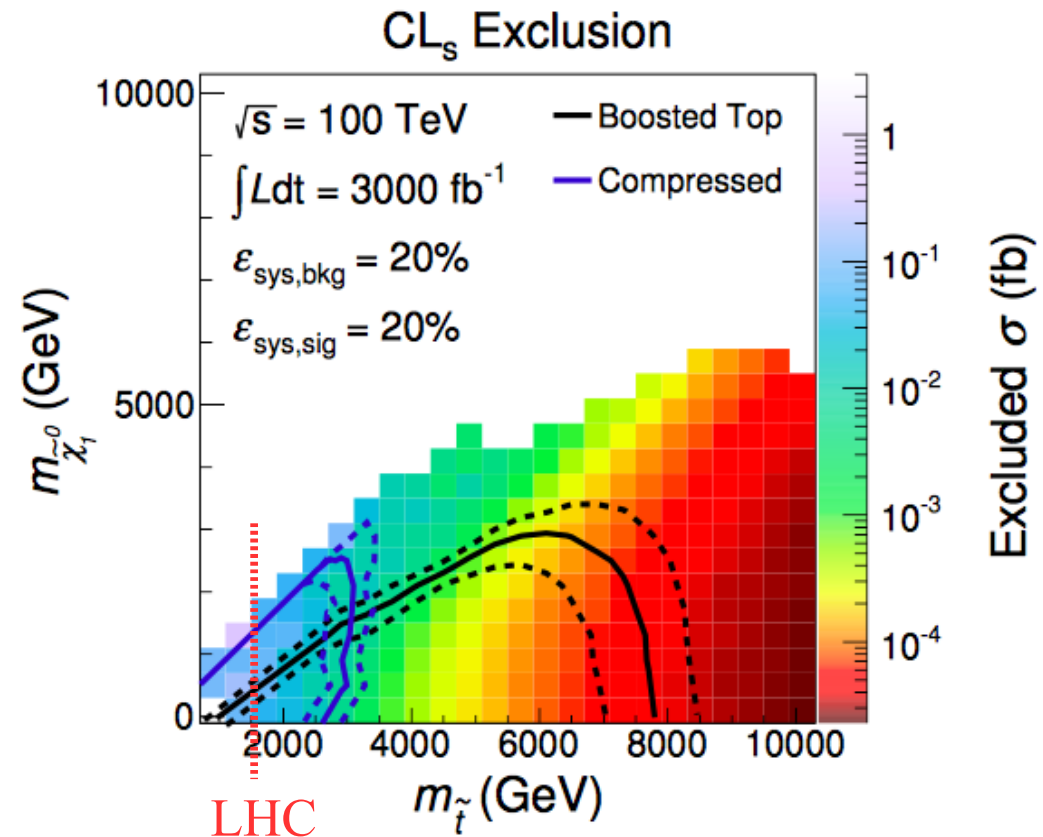
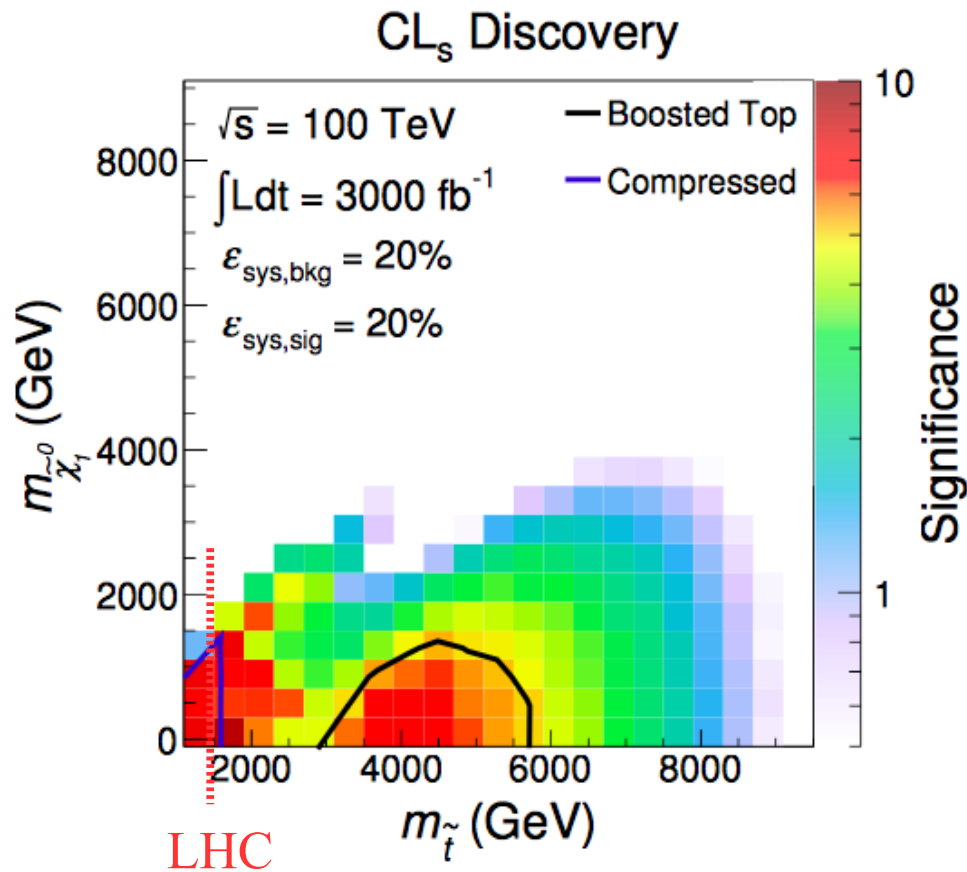
(AVK, P. Winslow,  
J. M. No,  
M. J. Ramsey-Musolf,  
**Phys.Rev. D94 (2016),  
035022)**

Discovery potential across entire parameter space with next collider



# Supersymmetric Colored Top Partner Sensitivity

(Cohen *et al*, 2014)

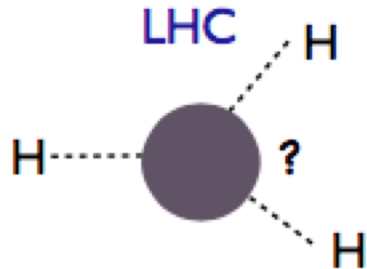


$$\text{Fine-tuning} \sim m_{\text{stop}}^2 \sim 10^{-4}$$

A big jump beyond LHC

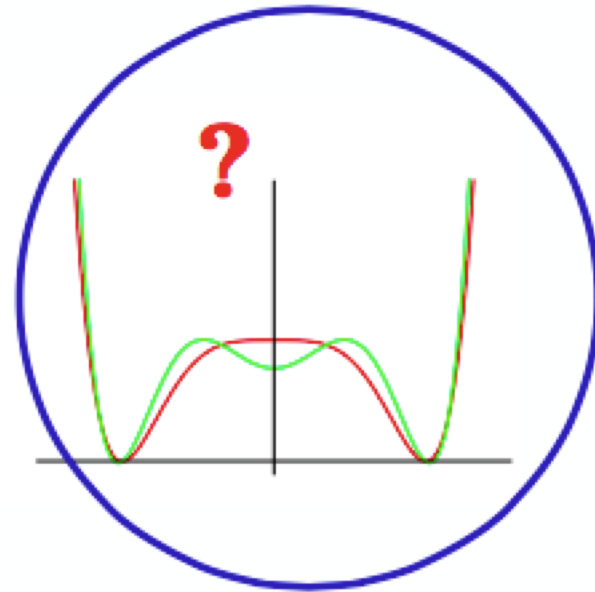
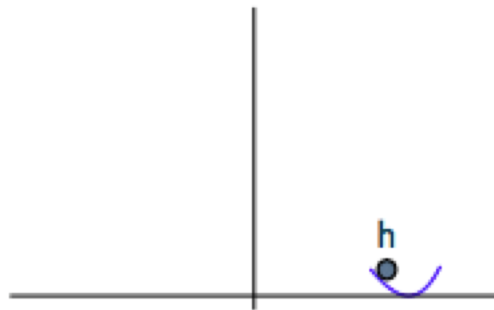
Discovering or eliminating “natural” low-energy SUSY

# Higgs Self-Coupling



Unique type of coupling for spin-0 scalars  
Not seen before in nature!

Measuring it well is crucial to  
answer this question.



Expect  $O(1)$  deviations from SM in self-coupling coefficient

# Measuring the Higgs Self-Coupling

- $gg \rightarrow HH$  (most promising?) ,  $qq \rightarrow HHqq$  (via VBF)
- Reference benchmark process:  $HH \rightarrow b\bar{b} \gamma\gamma$
- Goal: 5% (or better) precision for SM selfcoupling

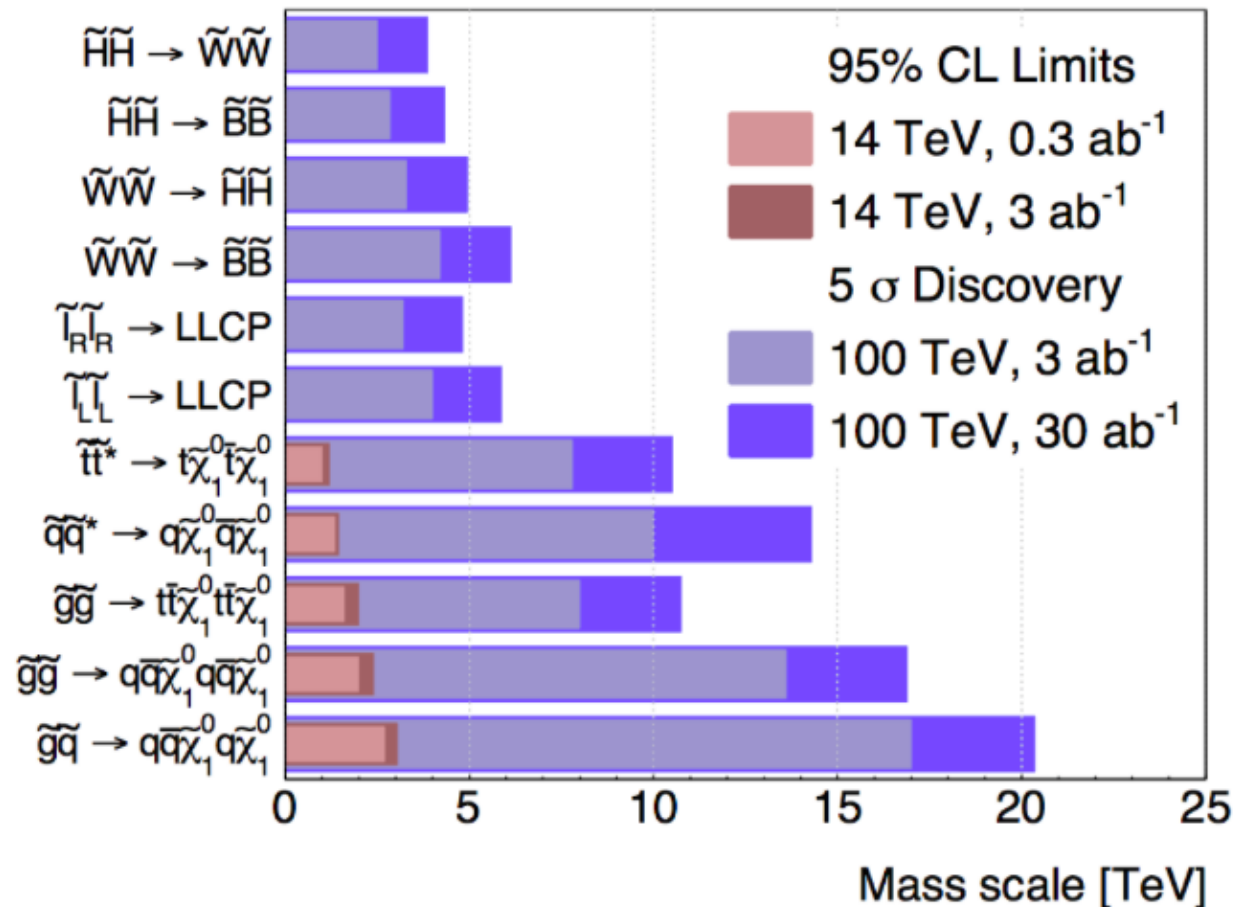
$HH \rightarrow b\bar{b} \gamma\gamma$	Barr,Dolan,Englert,Lima, Spannowsky JHEP 1502 (2015) 016	Contino, Azatov, Panico, Son arXiv:1502.00539	He, Ren Yao arXiv:1506.03302
FCC@100TeV 3/ab	30~40%	30%	15%
FCC@100TeV 30/ab	10%	10%	5%
$S/\sqrt{B}$	8.4	15.2	16.5
Details	<ul style="list-style-type: none"> <li>✓ <math>\lambda_{HHH}</math> modification only</li> <li>✓ <math>c \rightarrow b</math> &amp; <math>j \rightarrow \gamma</math> included</li> <li>✓ Background systematics</li> <li>○ <math>b\bar{b} \gamma\gamma</math> not matched</li> <li>✓ <math>m_{\gamma\gamma} = 125 \pm 1</math> GeV</li> </ul>	<ul style="list-style-type: none"> <li>✓ Full EFT approach</li> <li>○ No <math>c \rightarrow b</math> &amp; <math>j \rightarrow \gamma</math></li> <li>✓ Marginalized</li> <li>✓ <math>b\bar{b} \gamma\gamma</math> matched</li> <li>✓ <math>m_{\gamma\gamma} = 125 \pm 5</math> GeV</li> <li>✓ Jet / <math>W_{had}</math> veto</li> </ul>	<ul style="list-style-type: none"> <li>✓ <math>\lambda_{HHH}</math> modification only</li> <li>✓ <math>c \rightarrow b</math> &amp; <math>j \rightarrow \gamma</math> included</li> <li>○ No marginalization</li> <li>✓ <math>b\bar{b} \gamma\gamma</math> matched</li> <li>✓ <math>m_{\gamma\gamma} = 125 \pm 3</math> GeV</li> </ul>

**Work in progress to compare studies, harmonize performance assumptions, optimize, etc  
⇒ ideal benchmarking framework**



# Exploring New Territory – Squarks and Gluinos

## Summary from FCC Report:



Squark & gluino discovery potential up to 10-20 TeV

Full exploration of “low-scale” SUSY

# Summary

# Summary

- Entering new regime on detector design and technology
- Completion of the Standard Model and its consistency with all data implies
  - Energy scale of new physics is less well-defined now than when LHC was designed
  - We must prepare for a broader range of possible new physics
  - Specialized, targeted detectors risky as target signatures are unconstrained
  - Prudent to continue CDF & D0 (Run 2), ATLAS & CMS general-purpose detector philosophy
- Detectors will need to be more capable on all fronts
  - Faster
  - Larger dynamic range
  - Much higher resolution
  - Much higher granularity
  - Much more forward-detection capability
  - Much higher bandwidth, smarter triggers
- HL-LHC upgrade will provide experience and insights, but need to look beyond

# Summary

- Entering new regime on detector design and technology
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  - We must prepare for a broader range of possible new physics
  - Specialized, targeted detectors risky as target signatures are unconstrained
  - Prudent to continue CDF & D0 (Run 2), ATLAS & CMS general-purpose detector philosophy
- Need improved capabilities
  - Better track momentum resolution
  - Maintain/improve  $b$ -tagging at high jet  $p_T$  and high track density
  - Improve hadronic  $\tau$ -lepton identification efficiency  $\rightarrow$  high-granularity EMCAL
  - Boosted H/W/Z/top substructure  $\rightarrow$  high-granularity HCAL
  - Extend forward jet coverage to rapidity  $\sim 6$  for vector boson scattering
  - Extend forward tracking for rejecting top quark background and suppressing forward pileup jets

# Physics Case Studies and Seminars in US

- Seminar + Brainstorming Session Thursday 1 PM CST via ReadyTalk/Indico on some “hot topic” relevant for FCC-*hh*
  - Announcement on Fermilab Today / Labwide Calendar & Mailing list
  - VLHCPHYSICS@fnal.gov (or email me at kotwal@fnal.gov)
- Theme workshop series
  - Dark Matter (December 4-6, 2015 @ Fermilab)
  - Electroweak Baryogenesis (September 17-19, 2015 @ Univ. Mass Amherst)
  - New Symmetries
  - High-Granularity Calorimetry – GEANT simulation
- Resources:
  - Full analysis chain available for MADGRAPH + PYTHIA showering → Ntuples → repository → C++ analysis code
  - Argonne HEP analysis cluster for CPU and Ntuple storage
  - Quick ramp-up for anyone to pursue any model and channel of interest
  - Need experimentalists with analysis experience collaborating with theorists<sub>72</sub>



<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/FutureHadroncollider>

<http://cepc.ihep.ac.cn/preCDR/volume.html>

<http://tlep.web.cern.ch>

<http://lhec.web.cern.ch>

backup



# Timing

# Collider Luminosity and Sensor Timing

Luminosity is a measure of how often protons get close enough to interact

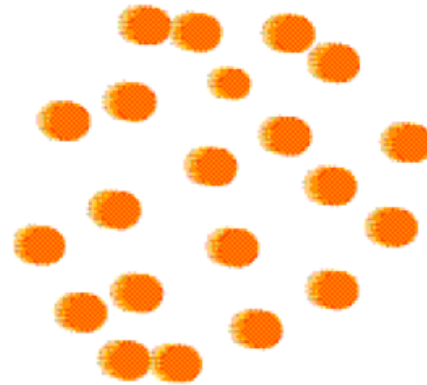
$$L = f \frac{n_1 n_2}{4\pi s_x s_y}$$

**f** = beam crossing frequency

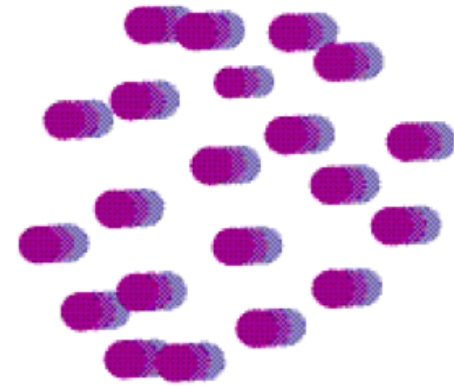
**n** = protons/bunch

**s** = transverse beam size

**L**  $\sim 10^{34}$  crossings/cm<sup>2</sup>/sec



50 ns  $\rightarrow$  25 ns at LHC



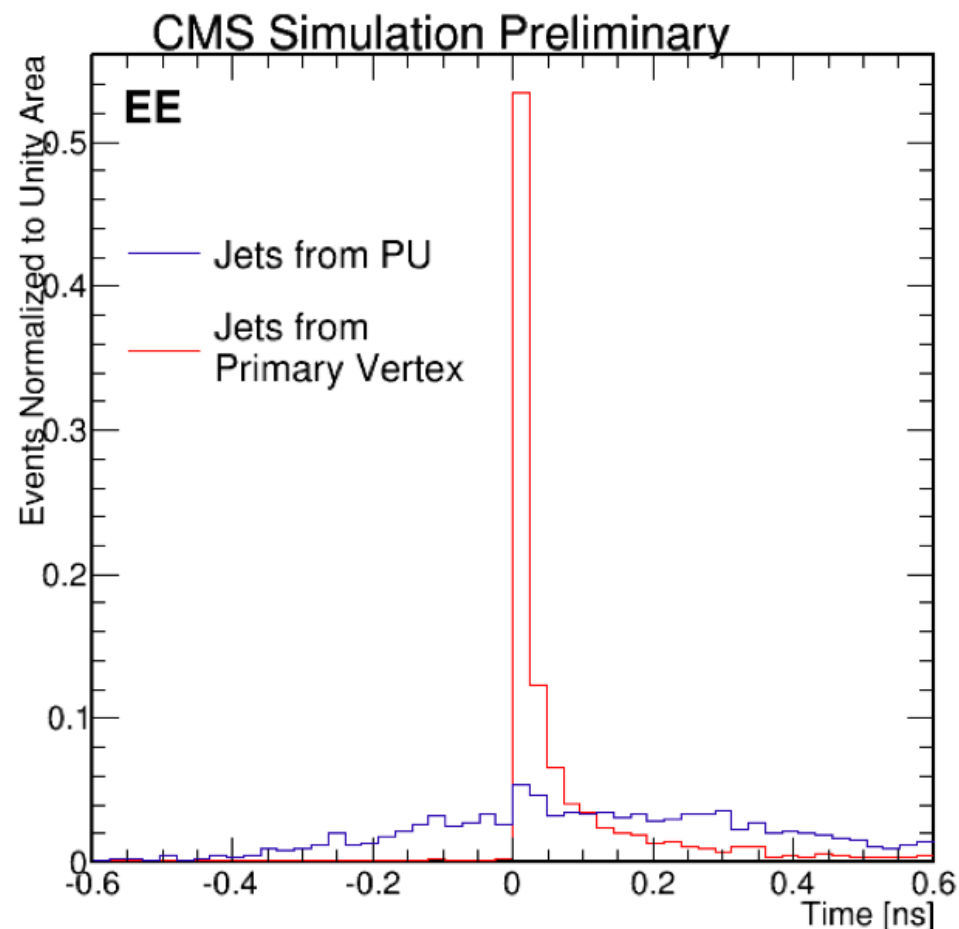
Reducing pileup by reducing  $n$  requires increasing  $f \Rightarrow$  faster detectors

5 ns option to be considered

Beam power increases in inverse proportion to crossing time (unless  $s$  reduced)

# ECAL CLEAN-UP USING TIMING

- **Effect of timing cut** on  $\Sigma E_T^{ECAL}$  variable
  - sum of all ECAL hits with  $E > 1\text{ GeV}$ .
- $O(30\text{ ps})$  resolution detector simulated
- Require ECAL timing (time-of-flight subtracted) within a **90 ps window**
- Most of the **PU extra energy gone**
  - able to almost recover no PU conditions
- Timing-based selection looks **promising for high PU environment**



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