

## Model Hamiltonian

$H = \epsilon_1 n_1 + \epsilon_2 n_2 + U_1 n_{1\uparrow} n_{1\downarrow} + U_2 n_{2\uparrow} n_{2\downarrow}$   
Orbital 1 and 2 with strong intra-level Coulomb repulsion ( $\epsilon_1 \leq \epsilon_2$ )

$+ U_{12} n_1 n_2$   
Inter-level Coulomb repulsion

$+ \omega_0 b^\dagger b$   
Local phonon mode

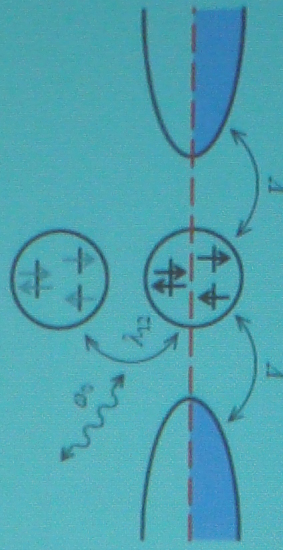
$+ \lambda_{12} (b^\dagger + b) \sum_{\sigma} (d_{1\sigma}^\dagger d_{2\sigma} + d_{2\sigma}^\dagger d_{1\sigma})$   
Phonon-assisted inter-orbital tunneling

$+ \sum_{k,\sigma} \epsilon_k c_{k\sigma}^\dagger c_{k\sigma}$   
Leads

$+ \sum_{j=1,2} V_j \sum_{k,\sigma} (c_{k\sigma}^\dagger d_{j\sigma} + d_{j\sigma}^\dagger c_{k\sigma})$   
Tunneling between molecule and leads

Regime of interest

$$U_j > |\epsilon_j|, \omega_0 \gg \Gamma_j = 2\pi \rho(E_F) V_j^2$$





## E-e and e-ph competition

$$d_{e\sigma} = (d_{1\sigma} + d_{2\sigma}) / \sqrt{2}$$

Define:

$$d_{o\sigma} = (d_{1\sigma} - d_{2\sigma}) / \sqrt{2}$$

Terms involving phonon mode can be written as:

$$\begin{aligned} \omega_0 b^\dagger b + \lambda_{12} (b^\dagger + b) \sum_{\sigma} (d_{1\sigma}^\dagger d_{2\sigma} + d_{2\sigma}^\dagger d_{1\sigma}) &= \omega_0 b^\dagger b + \lambda_{12} (b^\dagger + b) (n_e - n_o) \\ &= \omega_0 \tilde{b}^\dagger \tilde{b} - \left( \lambda_{12}^2 / \omega_0 \right) (n_e - n_o)^2 \end{aligned}$$

Increasing  $\lambda_{12}$  changes ground state from single-occupied to double-occupied:

$$n_e = 2 \quad n_o = 0$$

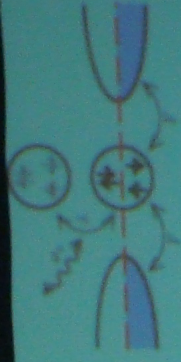
or

$$n_e = 0 \quad n_o = 2$$

Gain in phonon energy should balance Coulomb cost of double occupation around

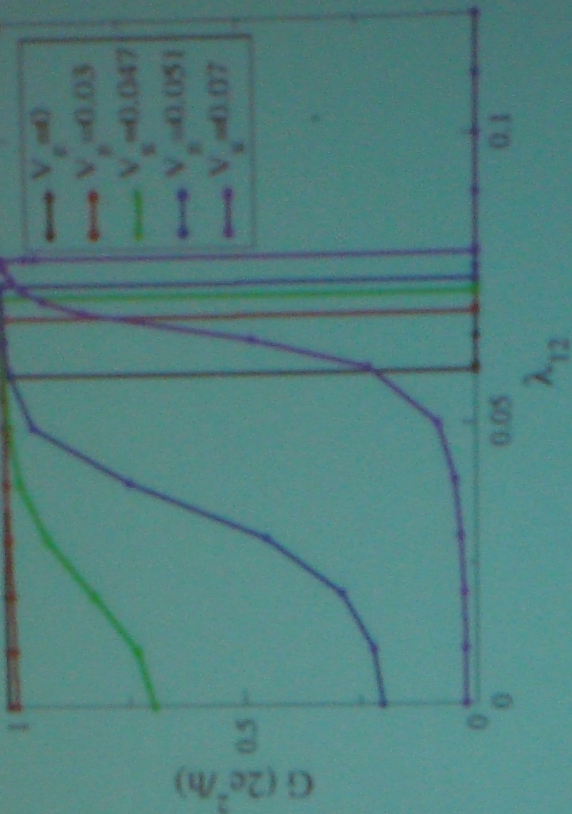
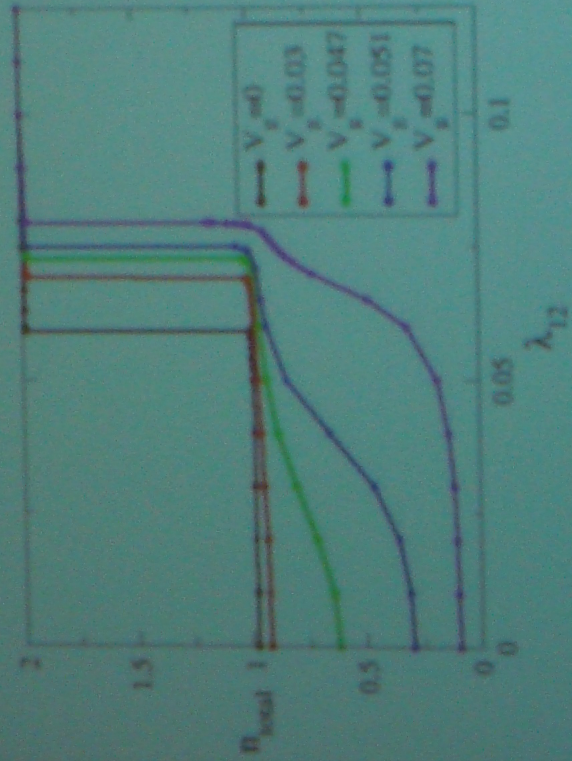
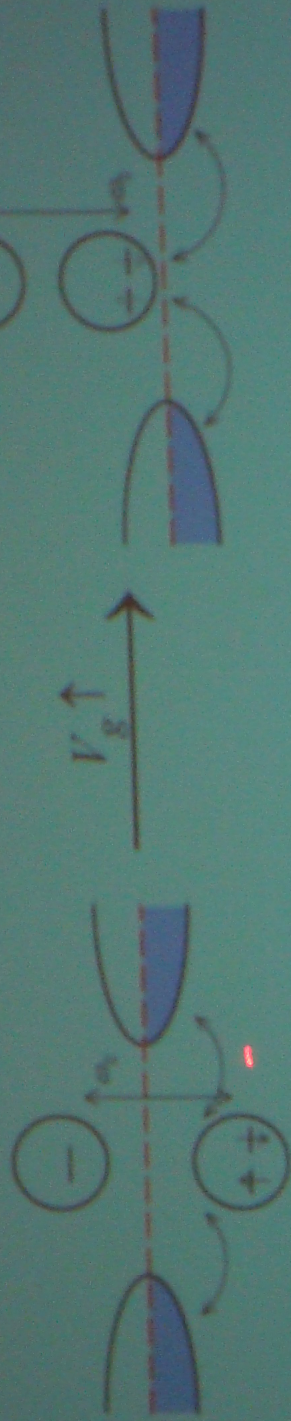
$$4\lambda_{12}^2 / \omega_0 = O(U_j)$$

Does the ground state change continuously or abruptly?





# QPT between Kondo phase and phonon-dominated phase



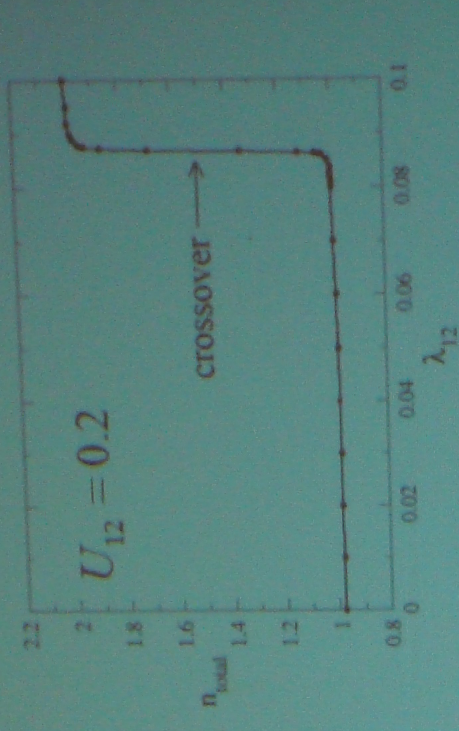
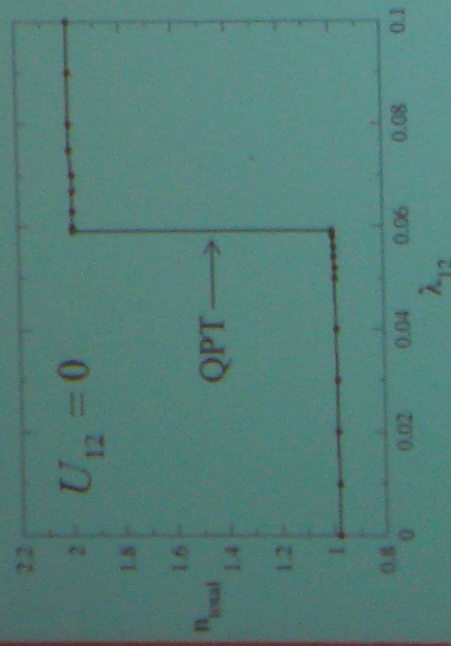
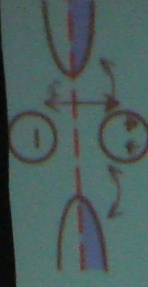
Kondo effect is killed at a first-order QPT

$$G = G_0 \sin^2(\pi n_{total}/2)$$

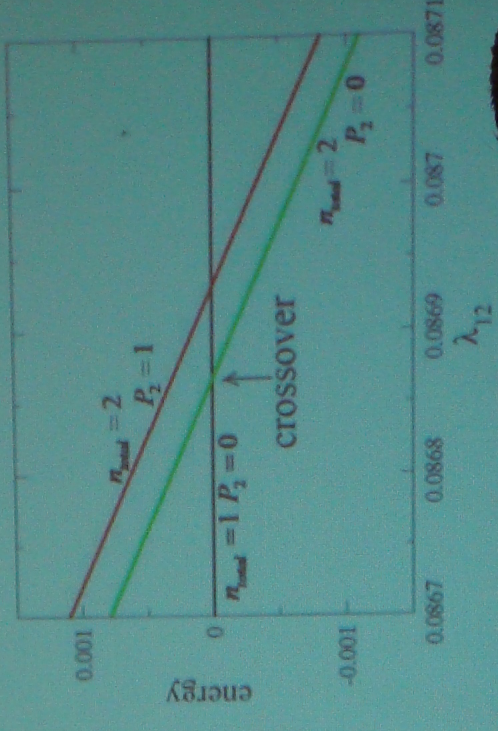
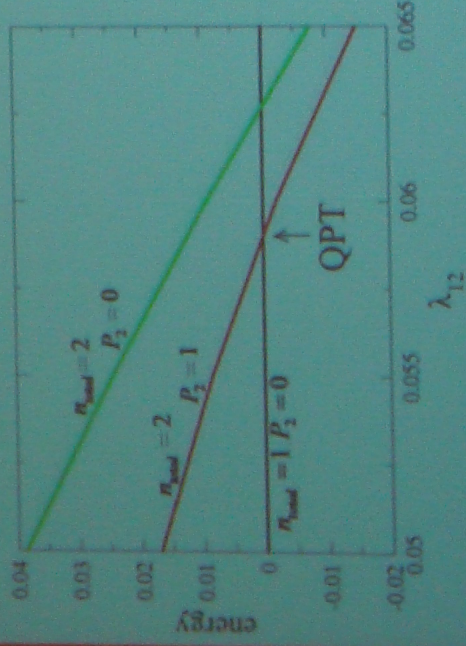


# Why is there a QPT?

For  $\Gamma_2 = 0$ ,  $n_2$  changes only under  $\lambda_{12}(b^\dagger + b) \sum_{\sigma} (d_{1\sigma}^\dagger d_{2\sigma} + d_{2\sigma}^\dagger d_{1\sigma})$   
 $P_2 = (n_2 + n_b) \bmod 2$  is conserved.

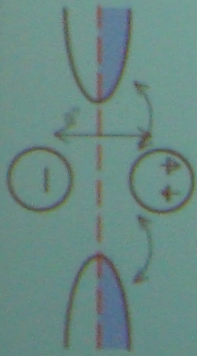


QPT  
vs.  
crossover



Lowest  
energies  
of isolated  
molecule



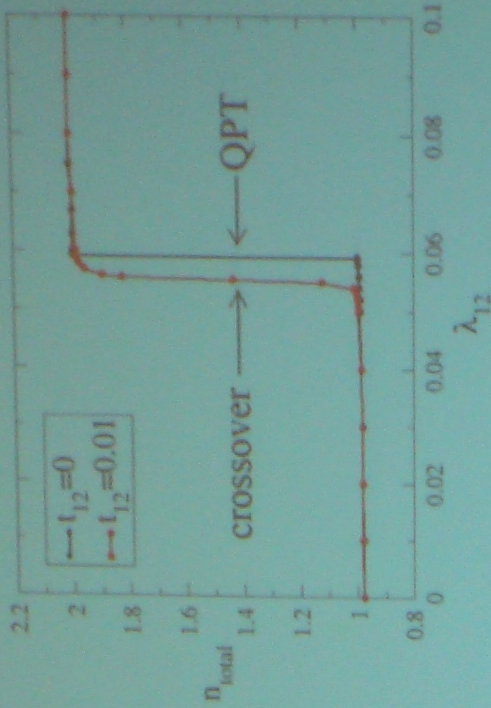
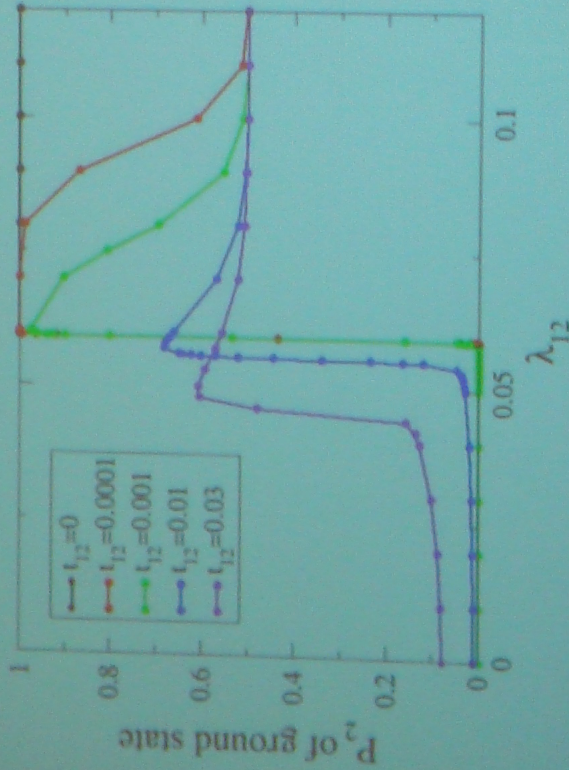


### QPT or crossover?

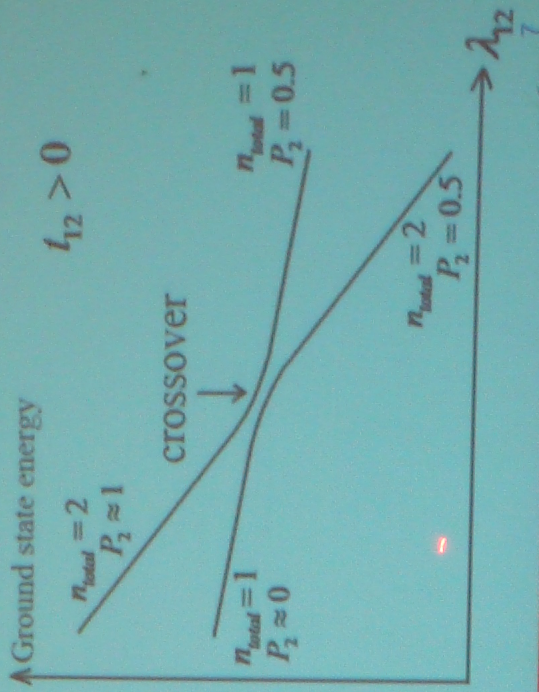
Turning on direct inter-level tunneling

$$t_{12} \sum_{\sigma} (d_{1\sigma}^{\dagger} d_{2\sigma} + d_{2\sigma}^{\dagger} d_{1\sigma})$$

Conservation of  $(n_2 + n_b) \bmod 2 = P_2$  is broken, QPT is replaced by crossover



### Schematic plot of ground state energy

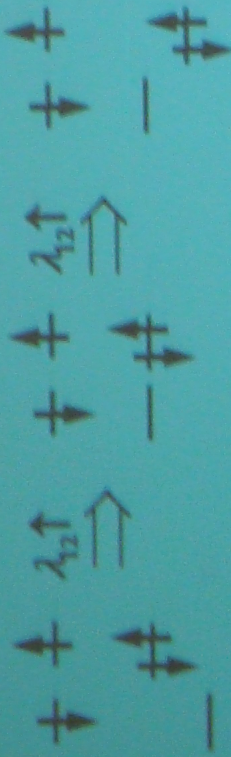




### Interplay of two phonon effects

$$\lambda_{12} \sum_{\sigma} (d_{1\sigma}^{\dagger} d_{2\sigma} + d_{2\sigma}^{\dagger} d_{1\sigma}) (b^{\dagger} + b)$$

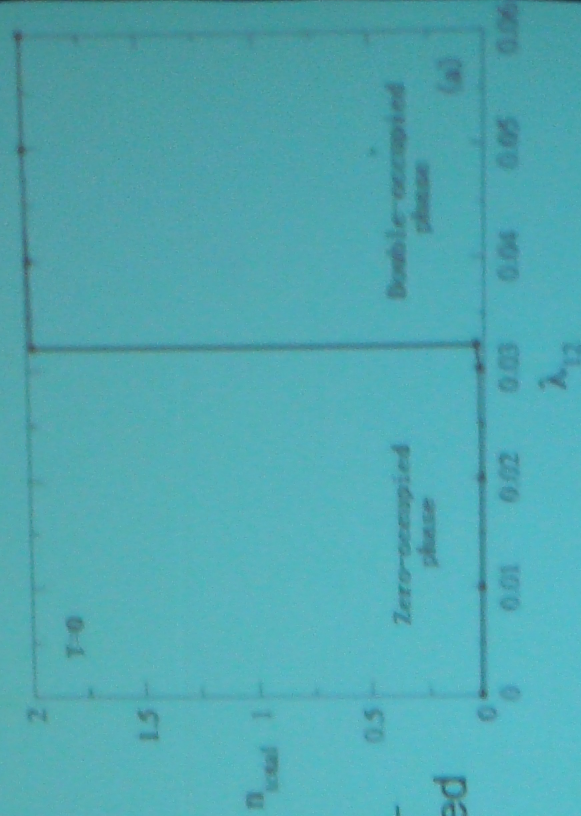
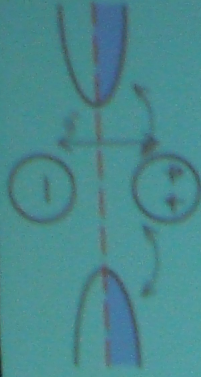
$$-\lambda_1 (n_1 - 1)(b^{\dagger} + b)$$



zero-occupied  $\rightleftharpoons$  degenerate  $\rightleftharpoons$  double-occupied

Phonon-assisted inter-level tunneling

Holstein-type phonon coupling to orbital-1

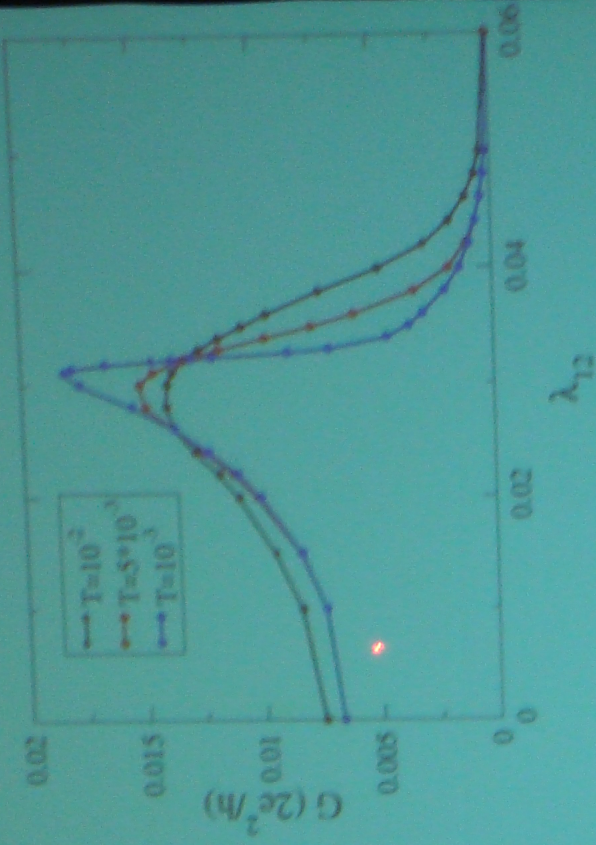


Does ground state change abruptly or continuously?

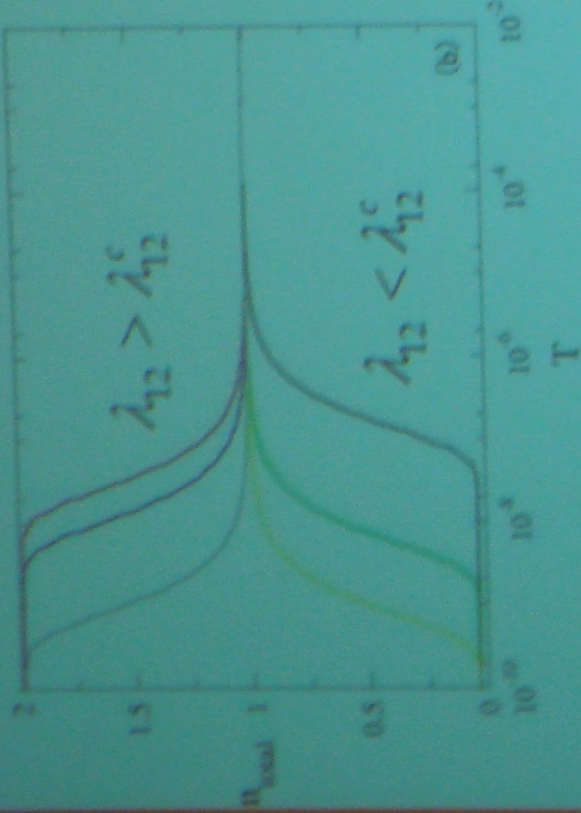
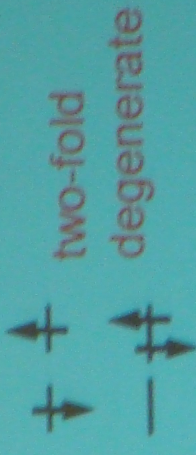


# Interplay of two phonon effects

## Finite temperature behavior



Conductance enhanced around:



Crossover temperature

$$T^* \propto |\lambda_{12} - \lambda_{12}^c|$$

QPT or not?



## Summary

- Have studied phonon-assisted inter-level tunneling in a single-molecule device.
- Competition between electron-electron and electron-phonon interactions can produce a first-order QPT between Kondo phase and phonon-dominated phase.
- The QPT arises from a special symmetry of the problem. Weak symmetry breaking leads to a sharp crossover instead.
- With Holstein phonon coupling, find a QPT or sharp crossover between two different phonon-dominated regimes.